

Measuring the gluon distribution of nuclei: diffractive e+A collisions at eRHIC

Matthew A. C. Lamont
Brookhaven National Lab



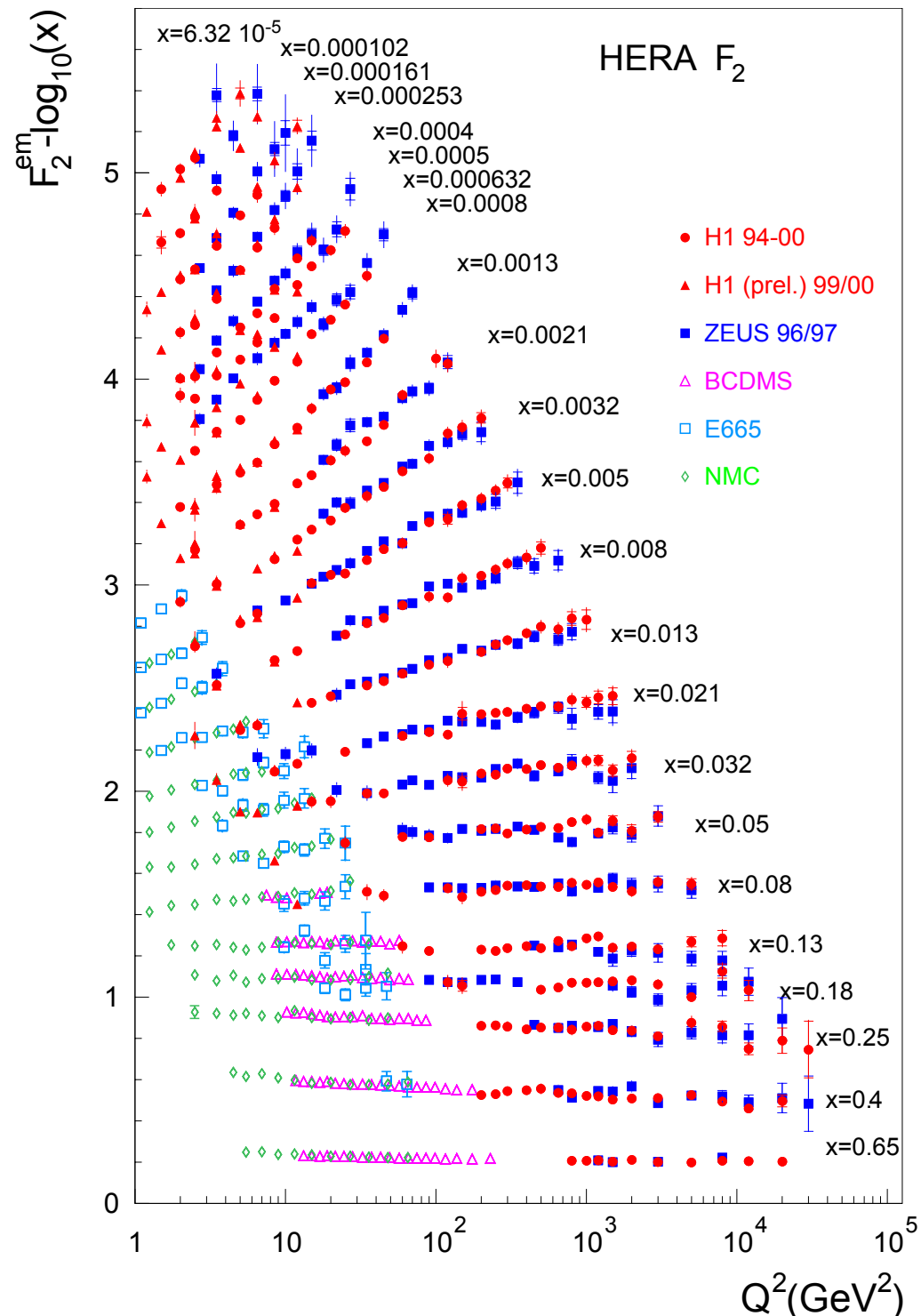


What did we learn from e+p collisions at HERA?

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

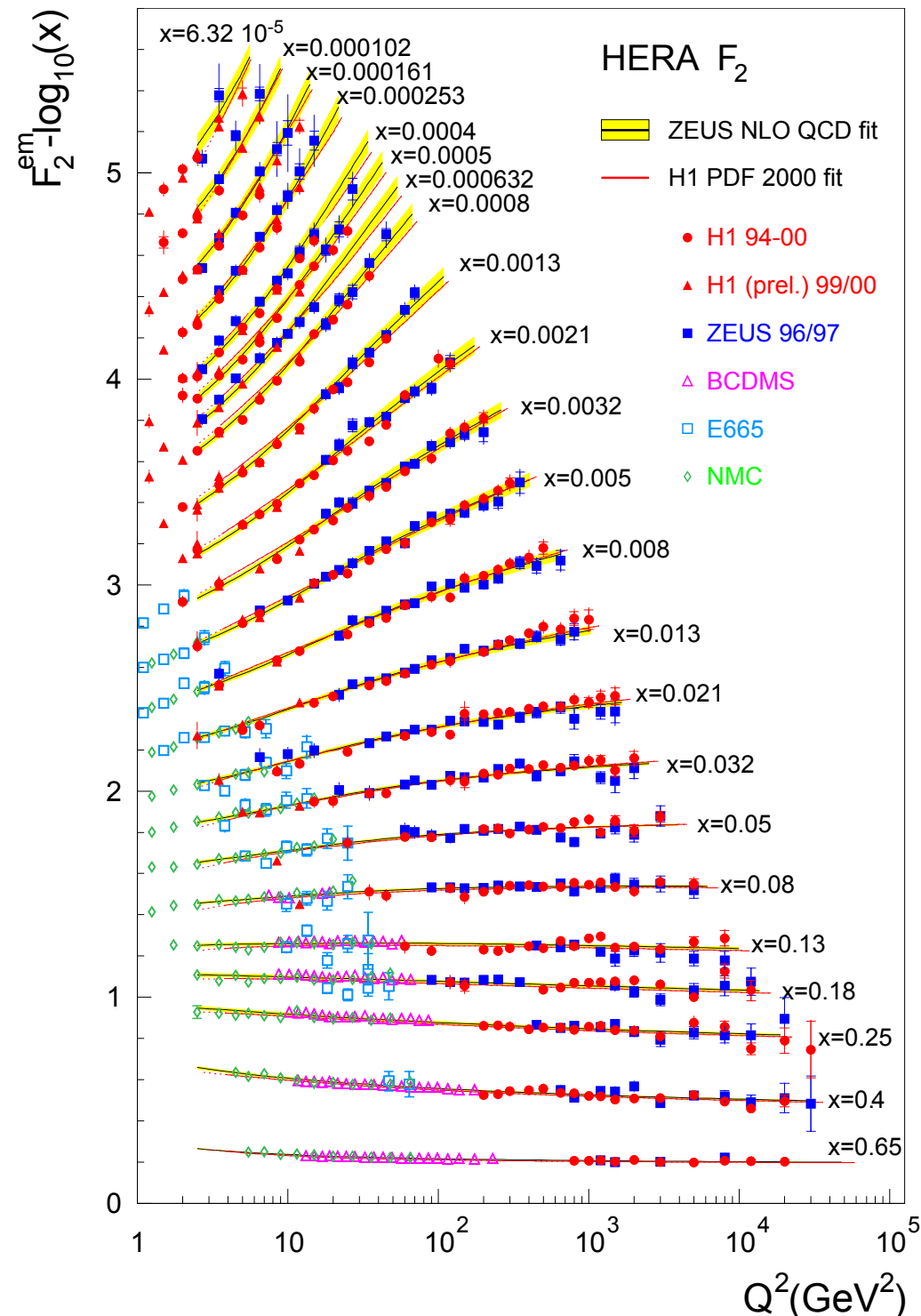
**quark+anti-quark
momentum distributions**

**gluon momentum
distribution**



What did we learn from e+p collisions at HERA?

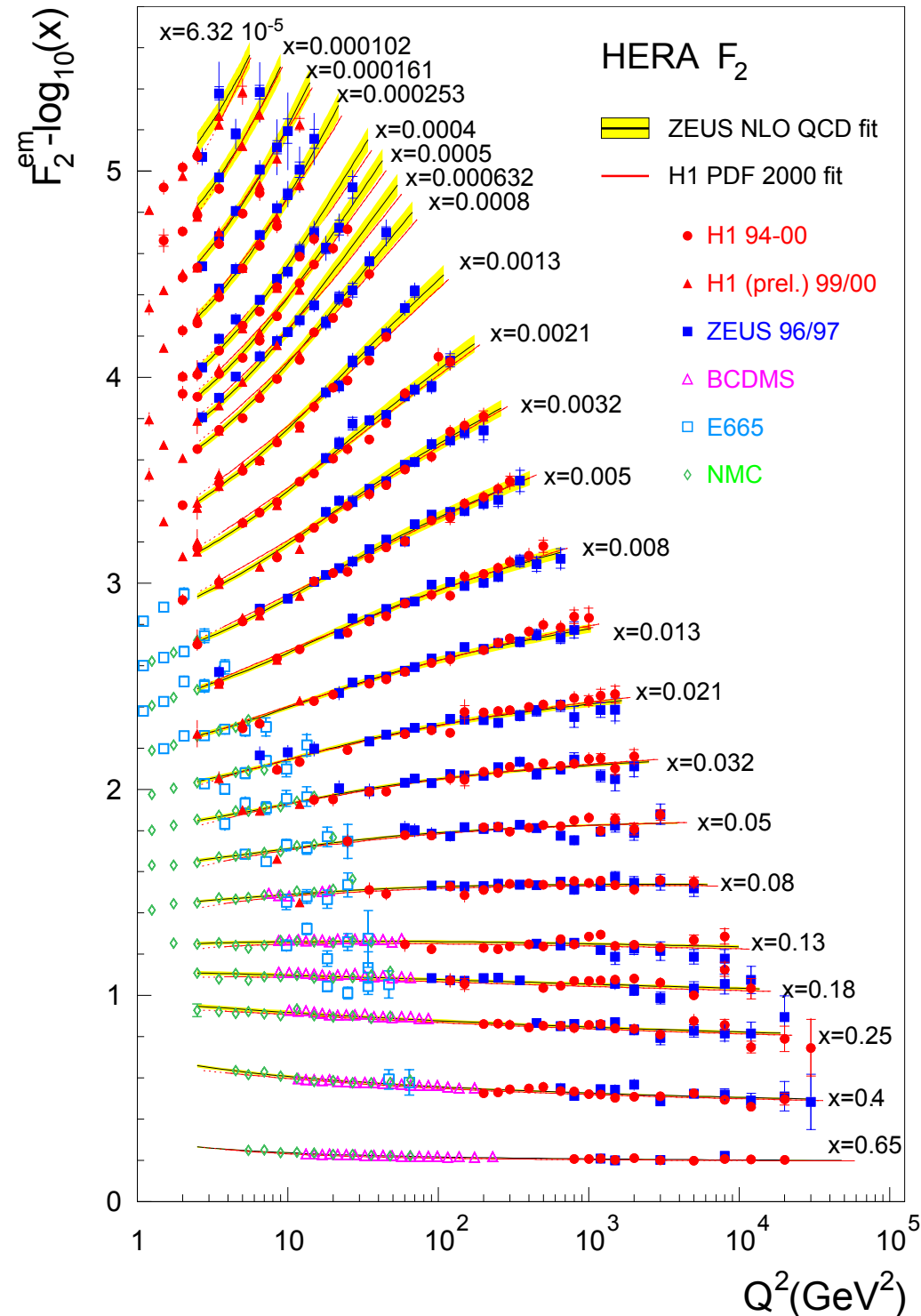
$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$



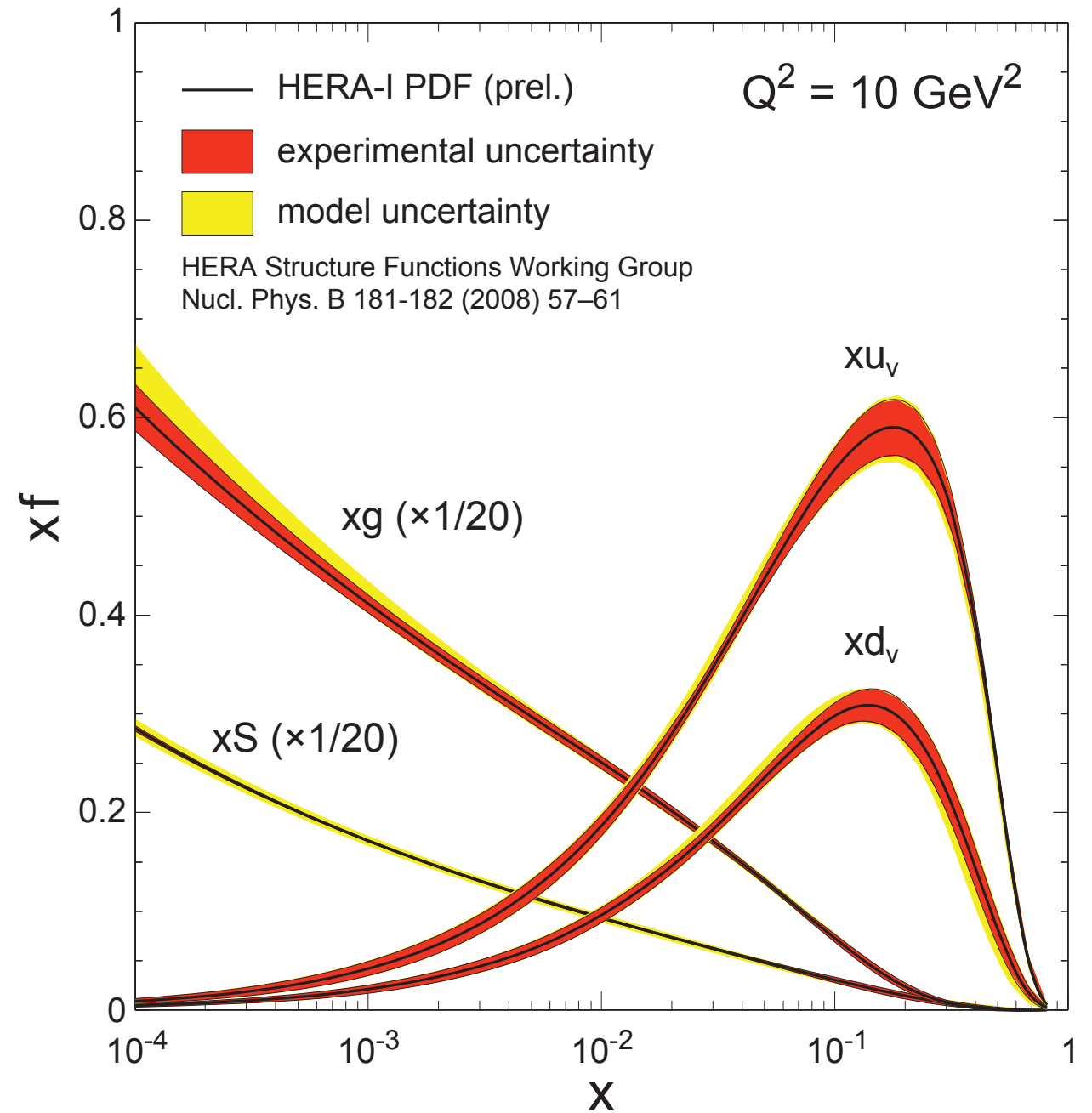
Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP Evolution $\Rightarrow G(x, Q^2)$

What did we learn from e+p collisions at HERA?

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

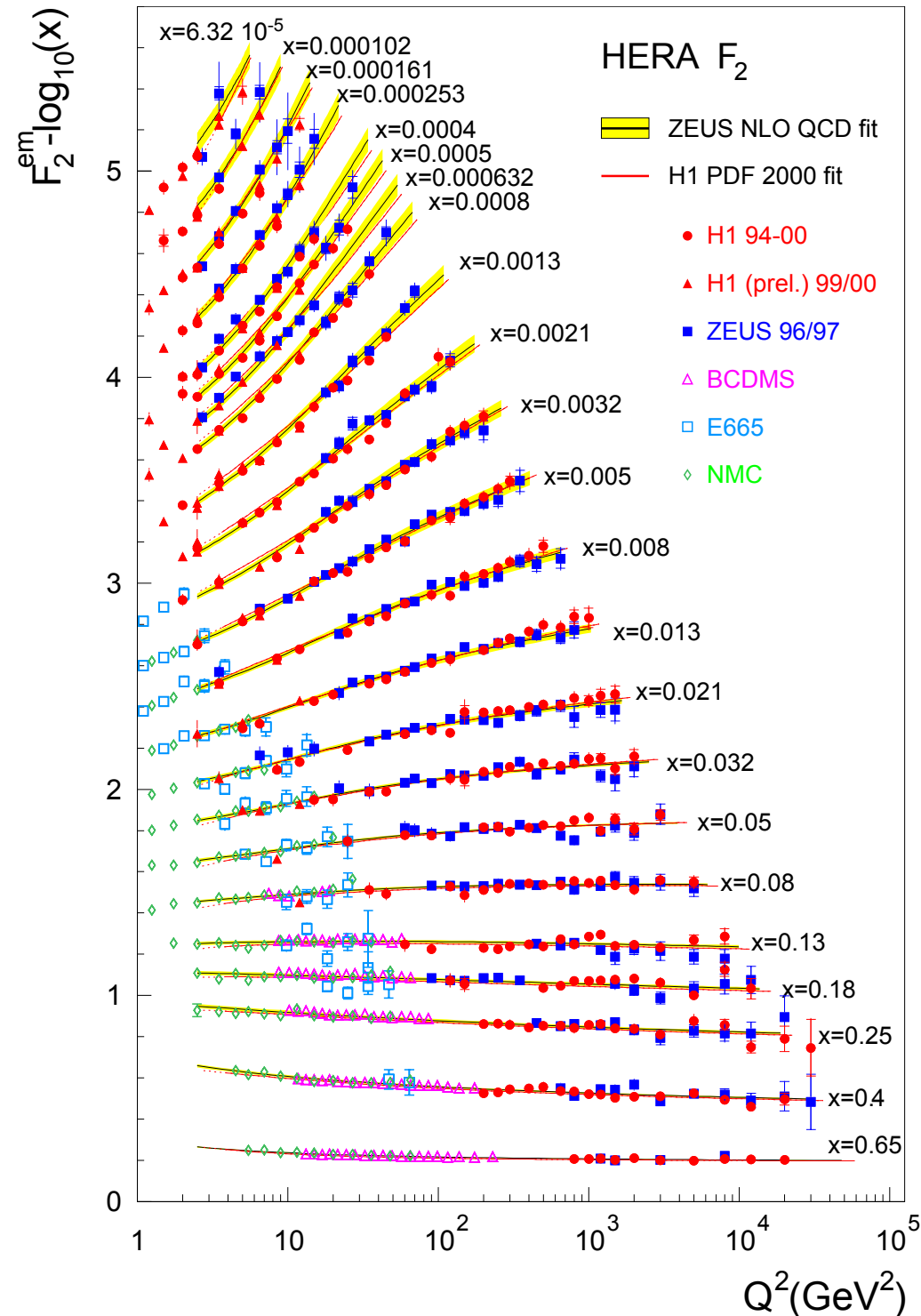


Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP Evolution $\Rightarrow G(x, Q^2)$

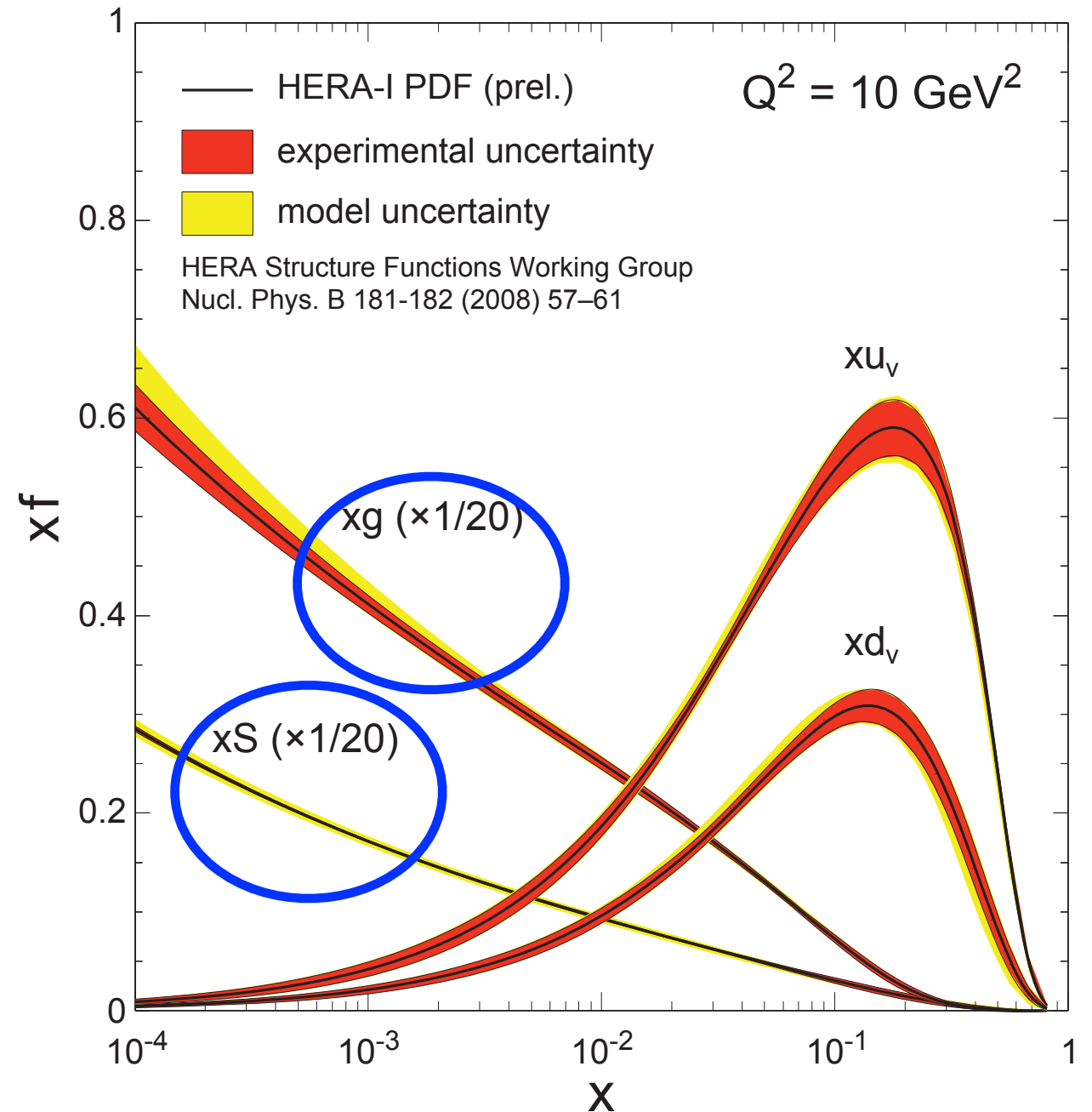


What did we learn from e+p collisions at HERA?

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

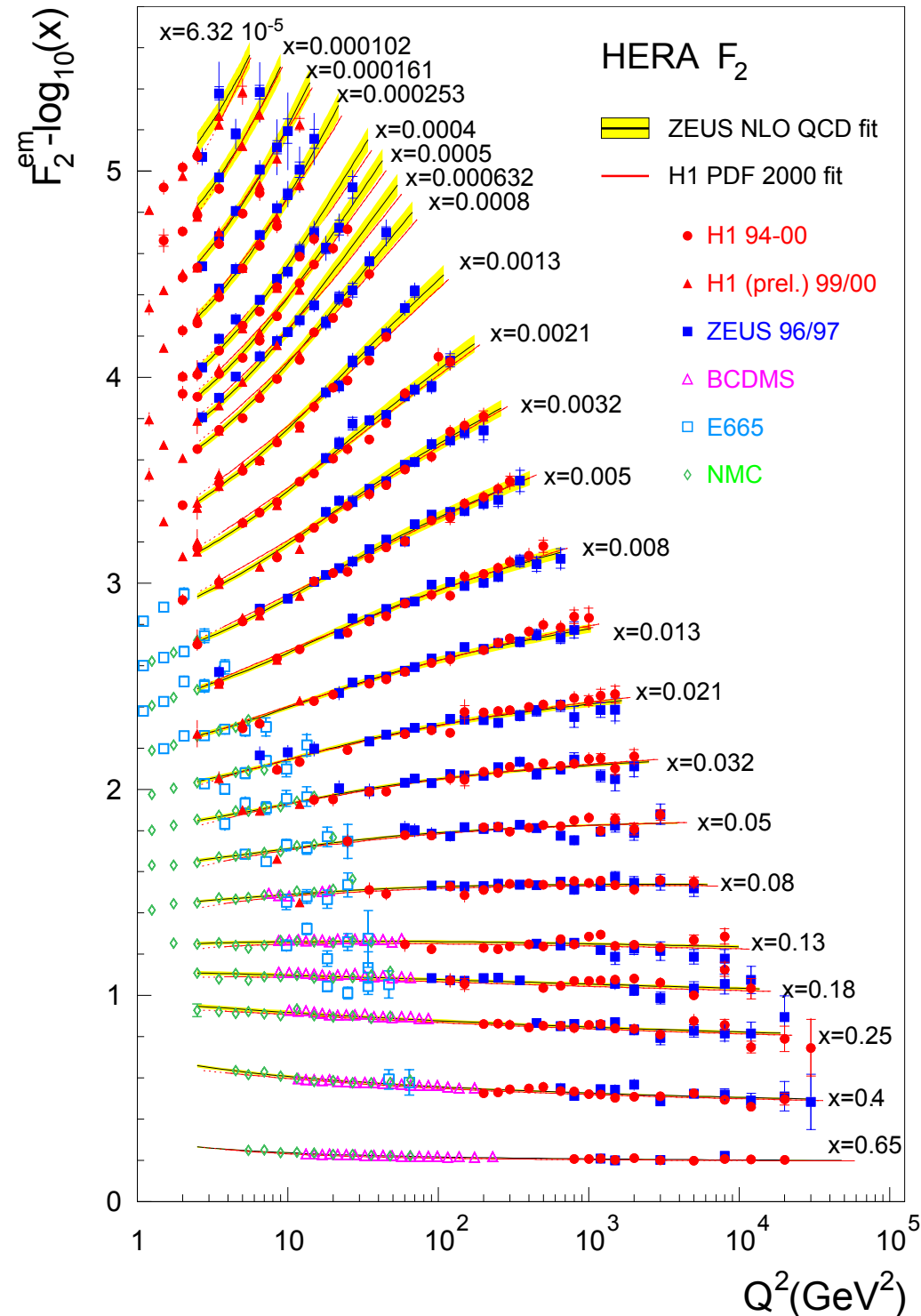


Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP Evolution $\Rightarrow G(x, Q^2)$

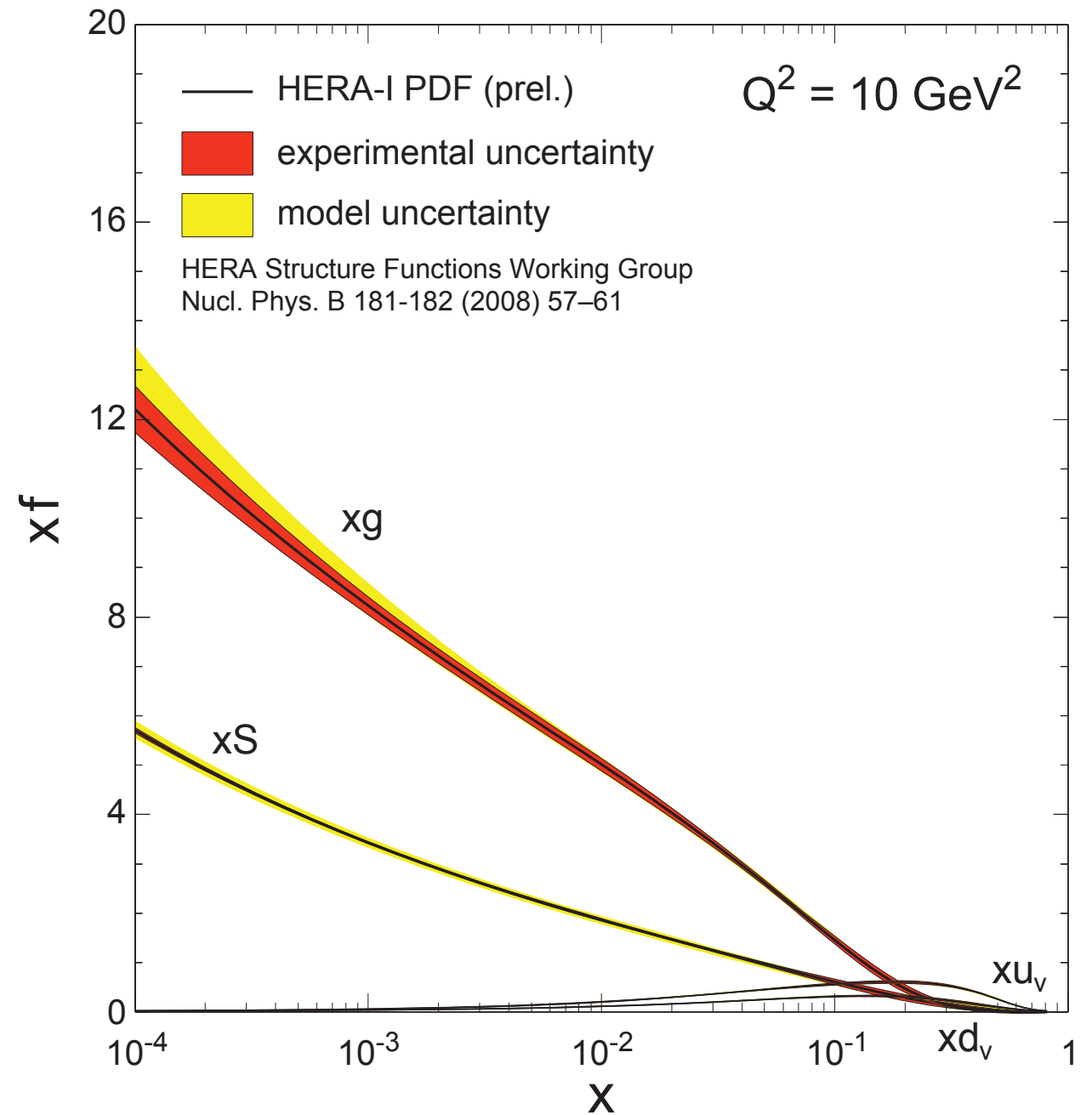


What did we learn from e+p collisions at HERA?

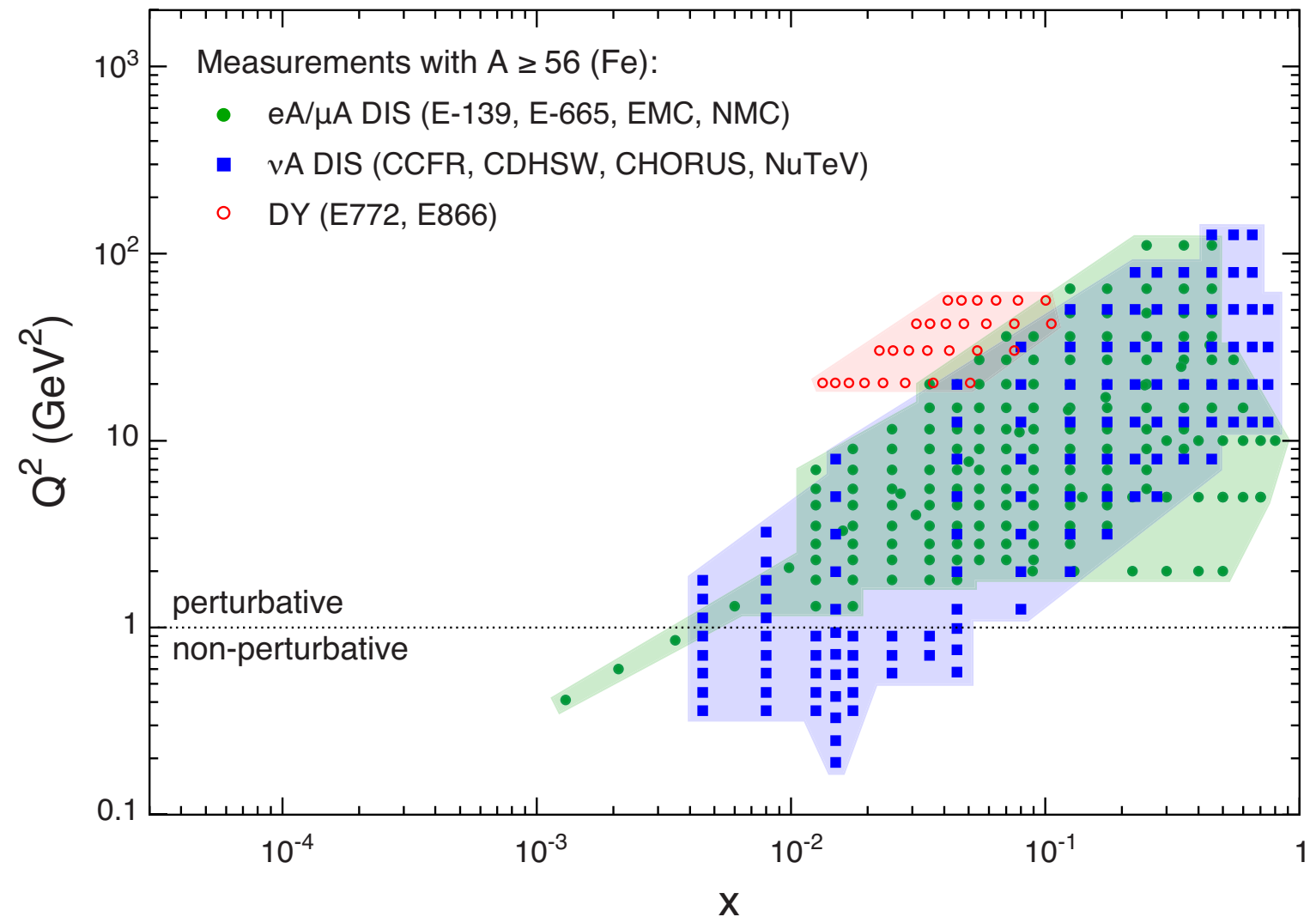
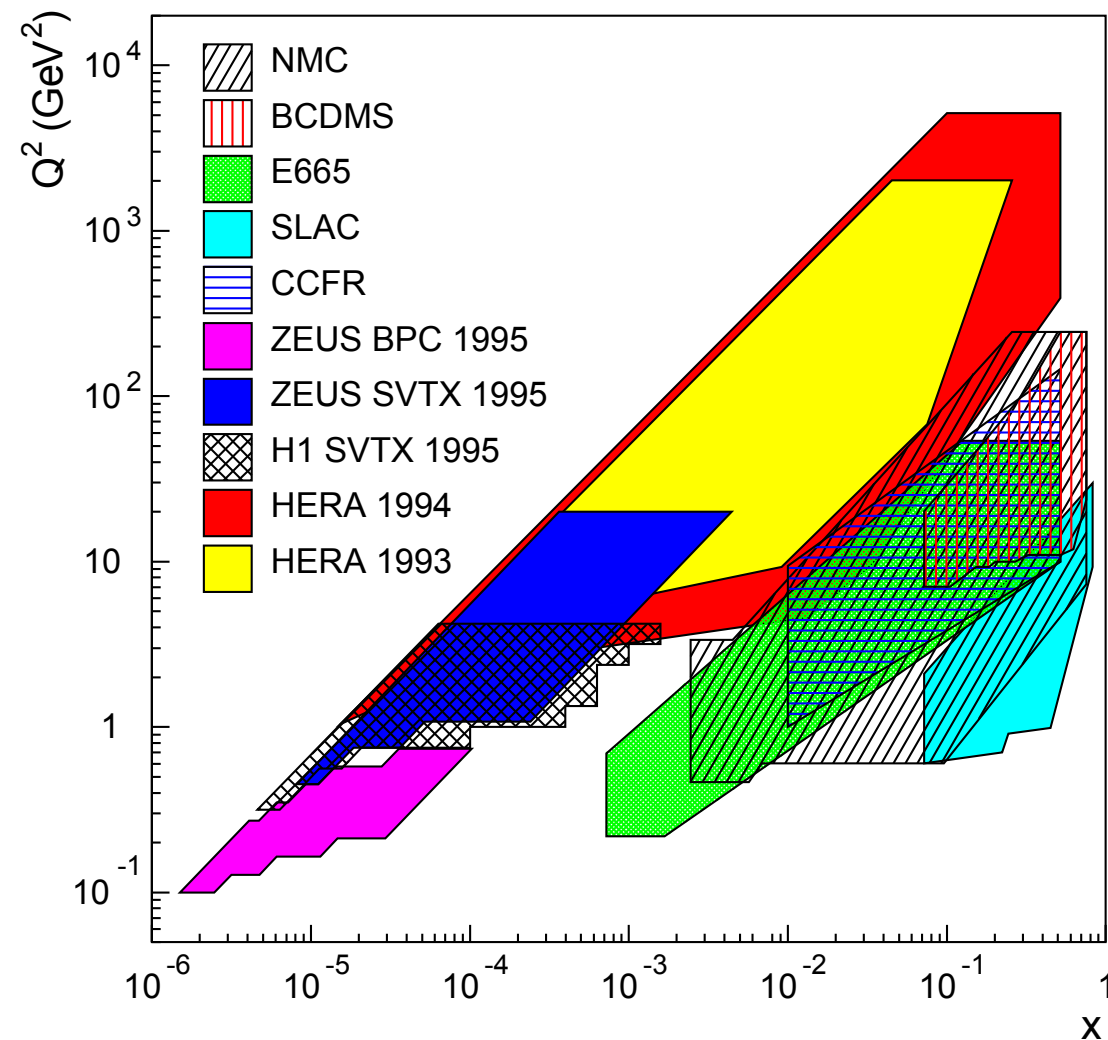
$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$



Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP Evolution $\Rightarrow G(x, Q^2)$



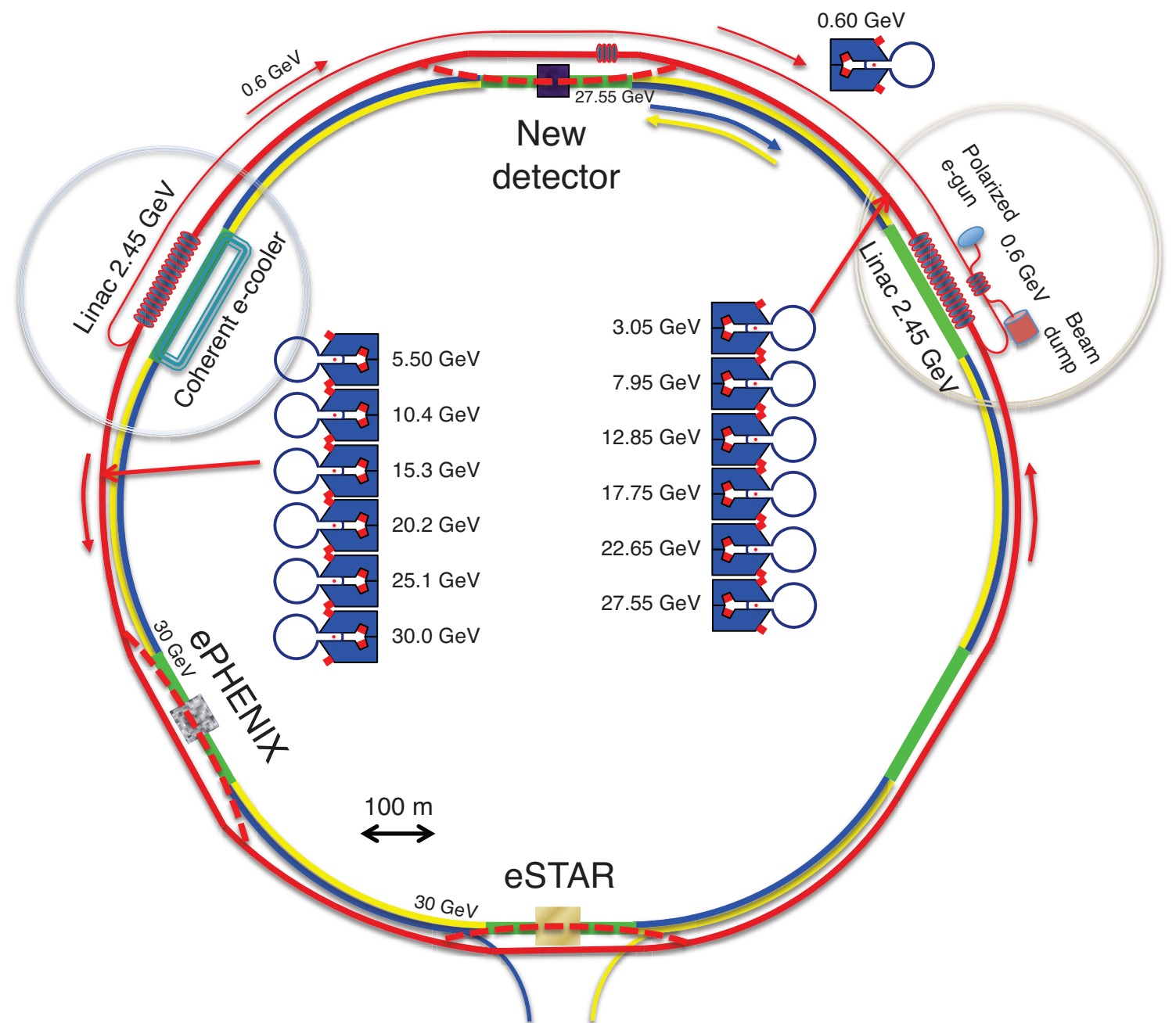
What do we know about the structure of nuclei?



- e+p data covers large part of phase space
 - ➔ low x and large Q^2
- e+A data only a small fraction of this (e+A was a fixed target programme at HERA)
 - ➔ high-medium x and low Q^2

The eRHIC project

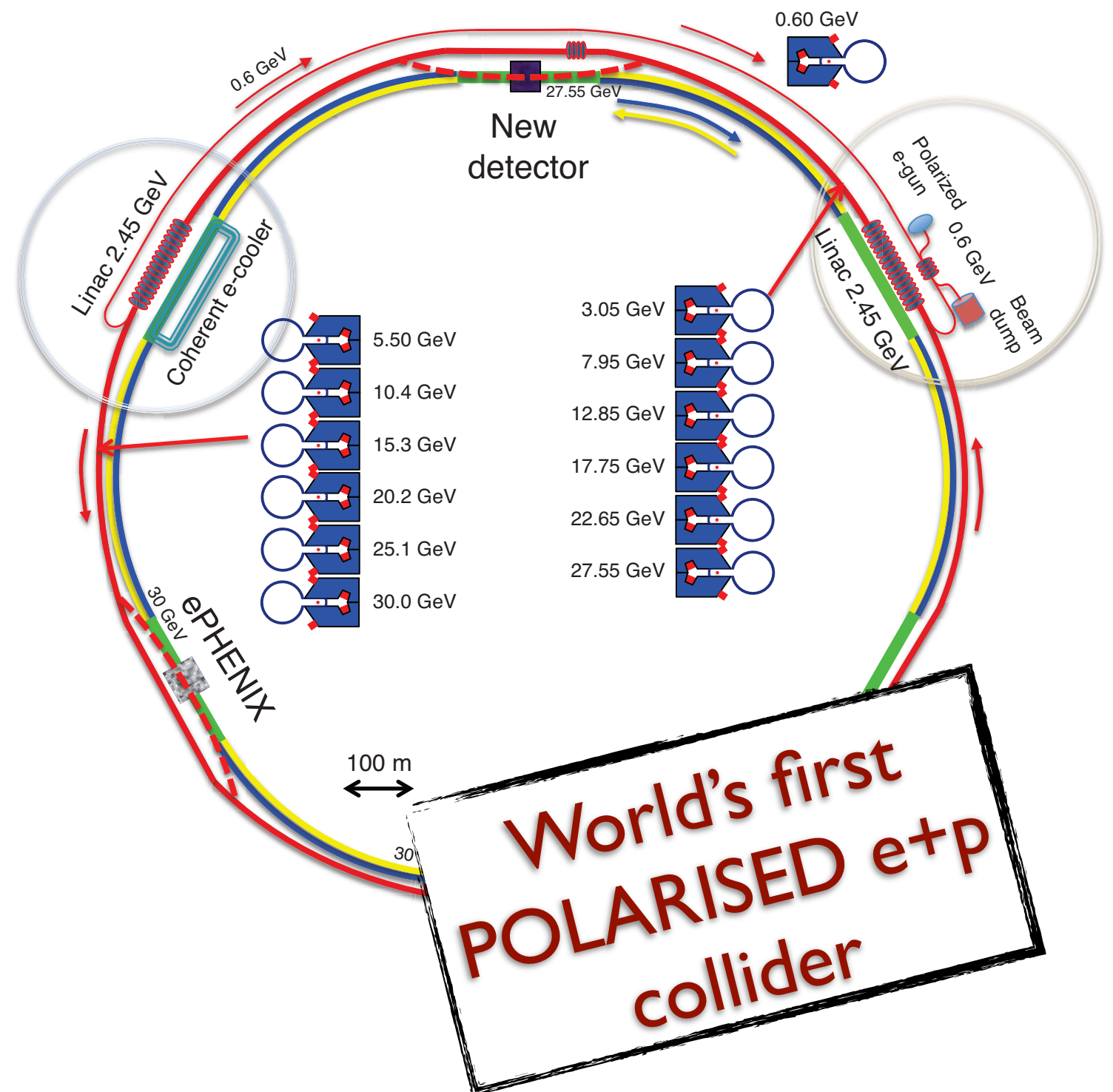
- eRHIC:
 - ➔ Utilises the RHIC ion beams
 - ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
 - ➔ 2-stage approach
 - Stage 1: e^- 5-10 GeV
 - Stage 2: e^- 20-30 GeV
 - ➔ Space for new detector at IP12
 - Possibilities for collisions in current STAR and PHENIX IPs



The eRHIC project

- eRHIC:

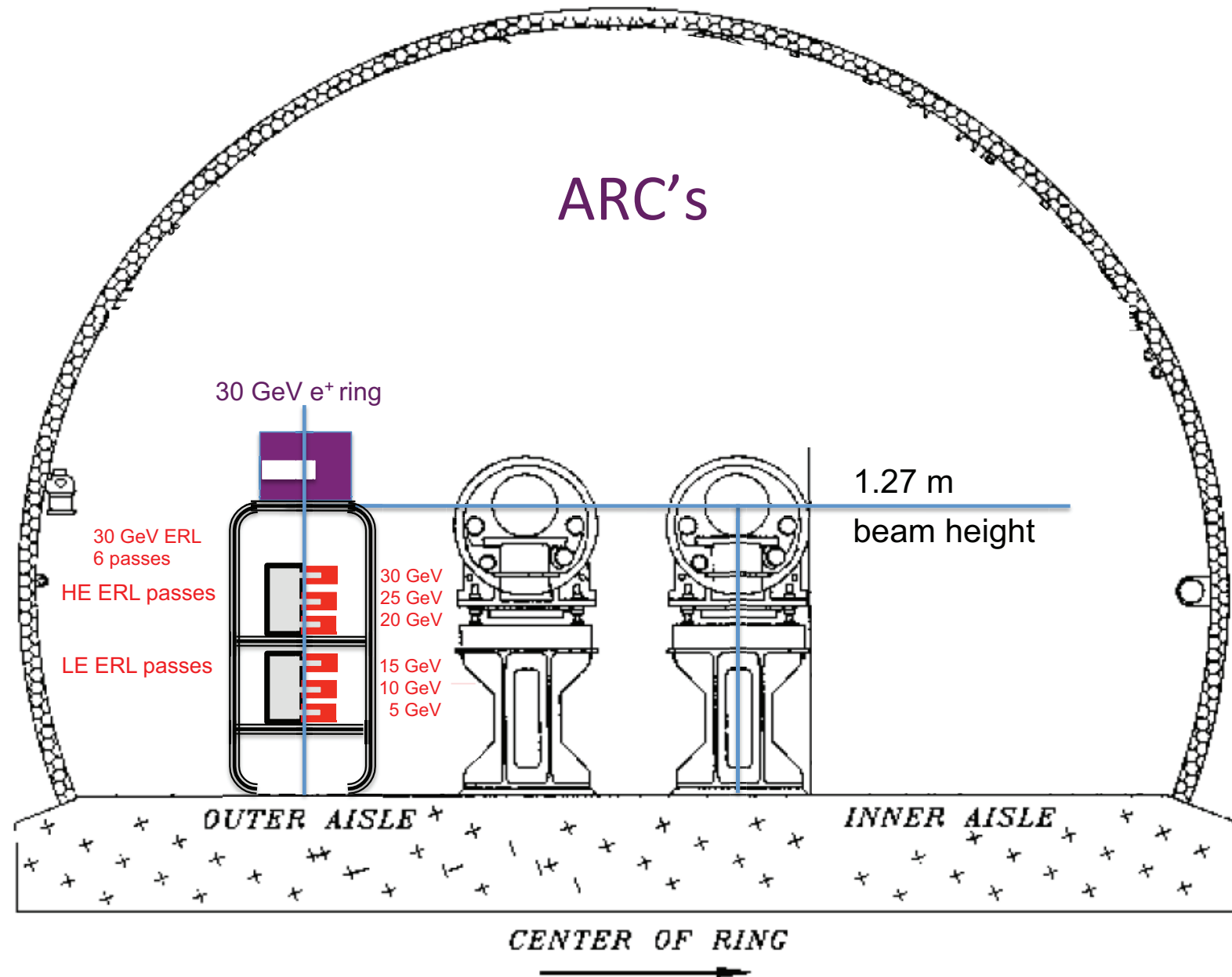
- ➔ Utilises the RHIC ion beams
- ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - ▶ 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
- ➔ 2-stage approach
 - ▶ Stage 1: e^- 5-10 GeV
 - ▶ Stage 2: e^- 20-30 GeV
- ➔ Space for new detector at IP12
 - ▶ Possibilities for collisions in current STAR and PHENIX IPs



The eRHIC project

- eRHIC:

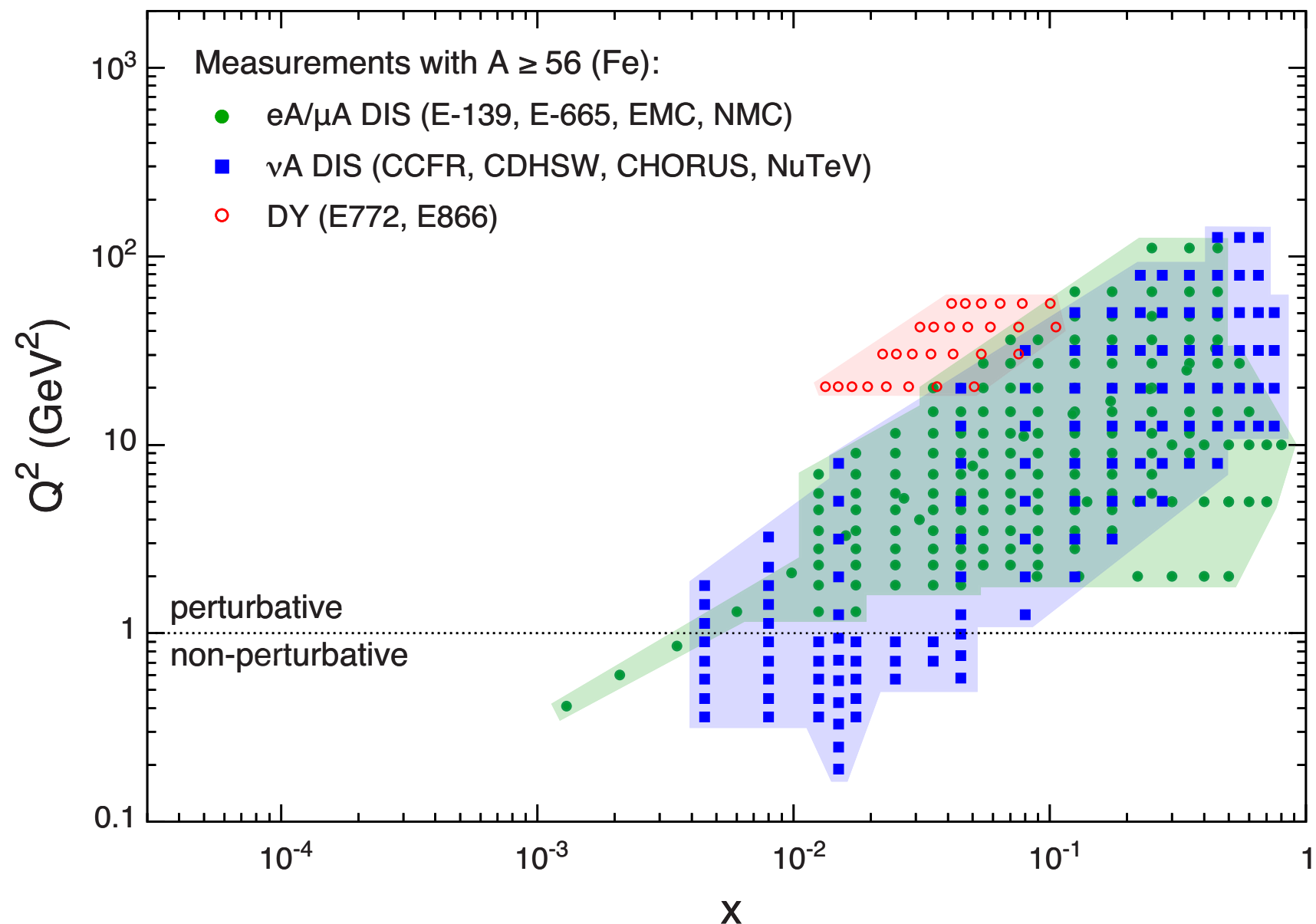
- ➔ Utilises the RHIC ion beams
- ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - ▶ 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
- ➔ 2-stage approach
 - ▶ Stage 1: e^- 5-10 GeV
 - ▶ Stage 2: e^- 20-30 GeV
- ➔ Space for new detector at IP12
 - ▶ Possibilities for collisions in current STAR and PHENIX IPs



The eRHIC project

- eRHIC:

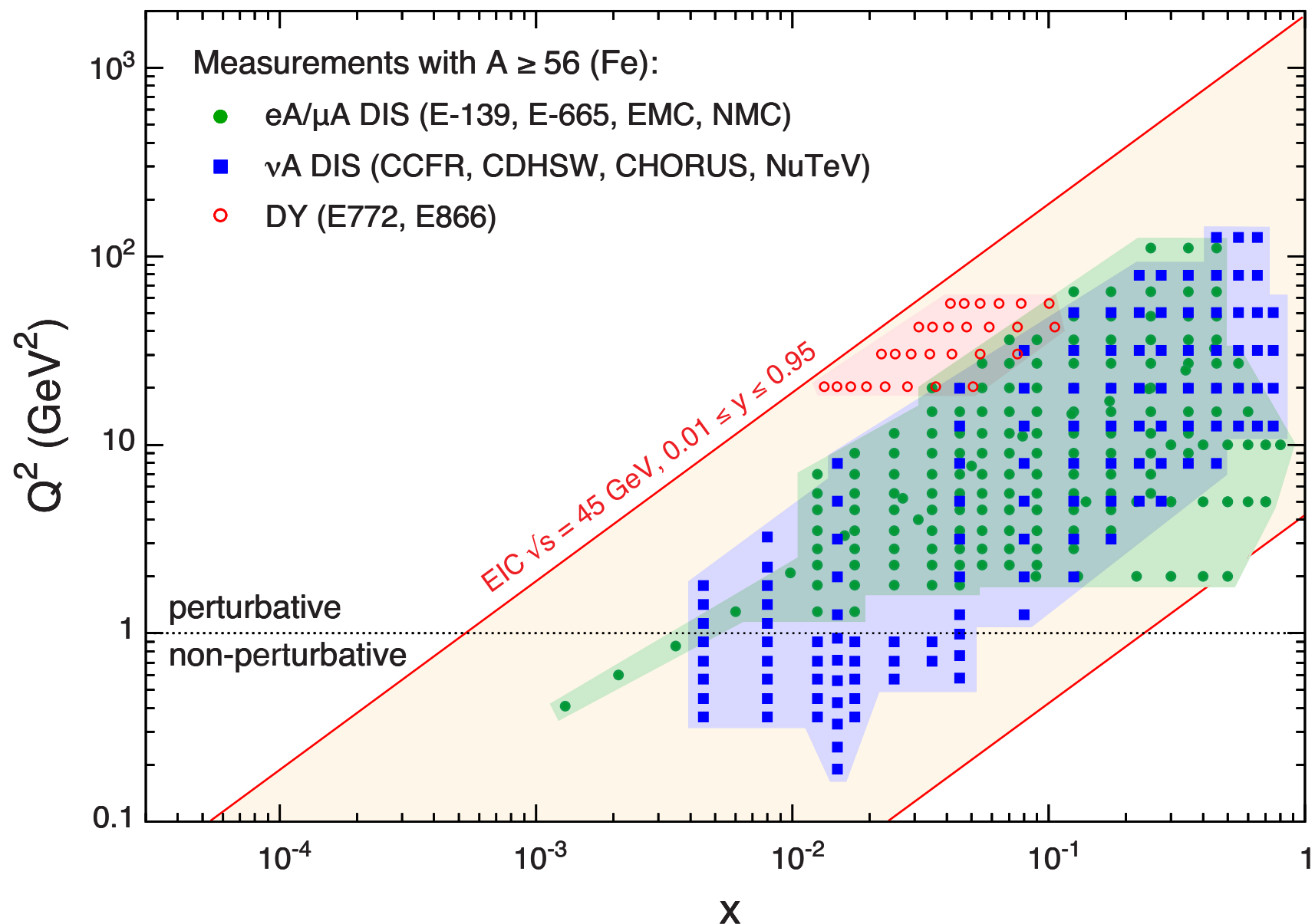
- ➔ Utilises the RHIC ion beams
- ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - ▶ 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
- ➔ Space for new detector at IP12
 - ▶ Possibilities for collisions in current STAR and PHENIX IPs



The eRHIC project

- eRHIC:

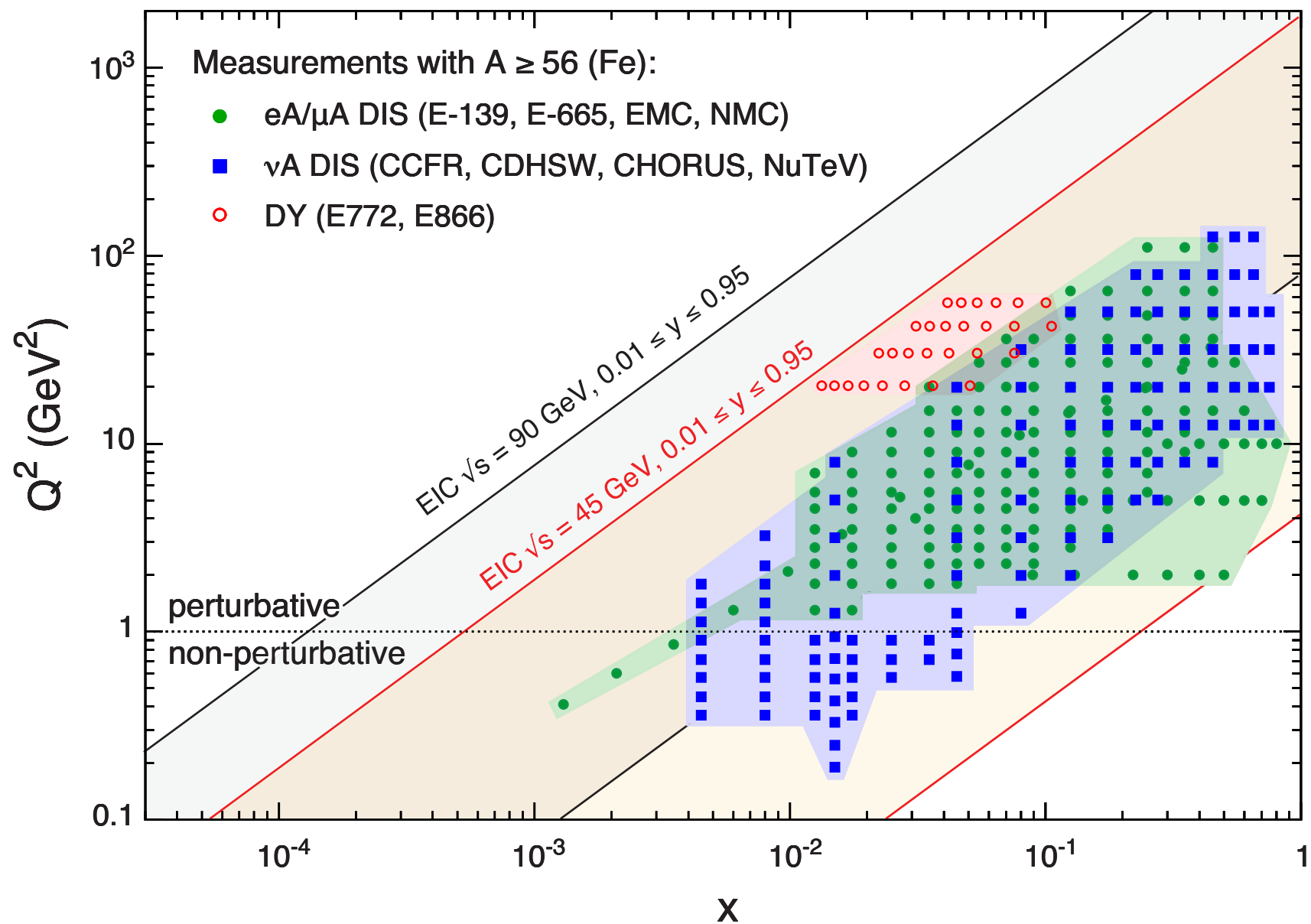
- ➔ Utilises the RHIC ion beams
- ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - ▶ 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
- ➔ Space for new detector at IP12
 - ▶ Possibilities for collisions in current STAR and PHENIX IPs



The eRHIC project

- eRHIC:

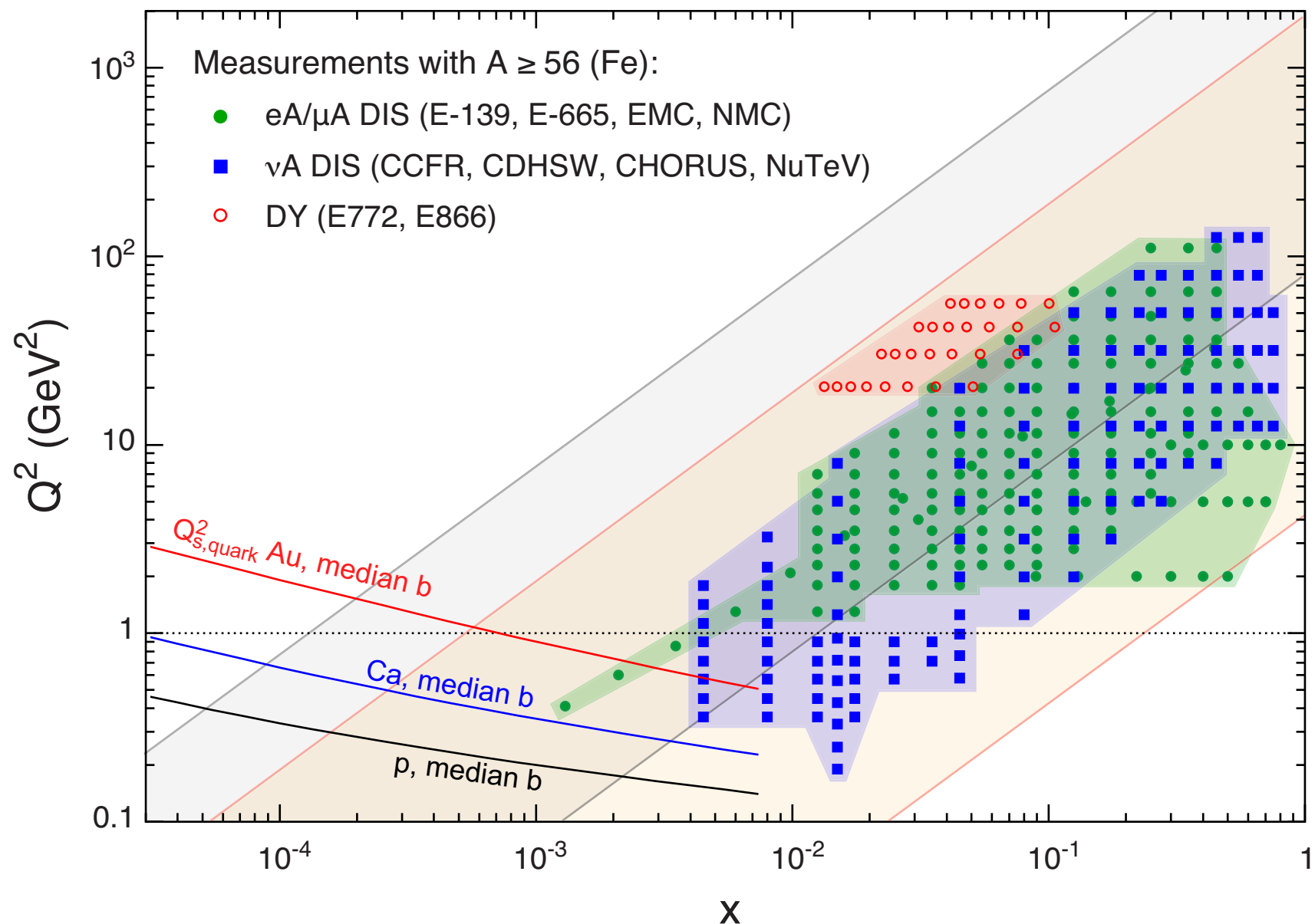
- ➔ Utilises the RHIC ion beams
- ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - ▶ 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
- ➔ Space for new detector at IP12
 - ▶ Possibilities for collisions in current STAR and PHENIX IPs



The eRHIC project

- eRHIC:

- ➔ Utilises the RHIC ion beams
- ➔ Two 2.45 GeV Energy Recovery Linacs (ERLs) accelerate the e^- beam
 - ▶ 6 separate rings accelerate the e^- up to a maximum energy of 30 GeV
- ➔ Space for new detector at IP12
 - ▶ Possibilities for collisions in current STAR and PHENIX IPs





Fundamental questions which arise:

- What is the role of strong gluon fields, parton saturation effects and collective gluon excitations in scattering off nuclei?



Fundamental questions which arise:

- What is the role of strong gluon fields, parton saturation effects and collective gluon excitations in scattering off nuclei?
- Can we experimentally find the evidence of non-linear QCD evolution in high-energy scattering off nuclei?



Fundamental questions which arise:

- What is the role of strong gluon fields, parton saturation effects and collective gluon excitations in scattering off nuclei?
- Can we experimentally find the evidence of non-linear QCD evolution in high-energy scattering off nuclei?
- What is the momentum and spatial distribution of gluons and sea quarks in nuclei?



Fundamental questions which arise:

- What is the role of strong gluon fields, parton saturation effects and collective gluon excitations in scattering off nuclei?
- Can we experimentally find the evidence of non-linear QCD evolution in high-energy scattering off nuclei?
- What is the momentum and spatial distribution of gluons and sea quarks in nuclei?
- Are there strong colour (quark and gluon density) fluctuations inside of a large nucleus? How does the nucleus respond to the propagation of a colour charge through it?



Important Measurements

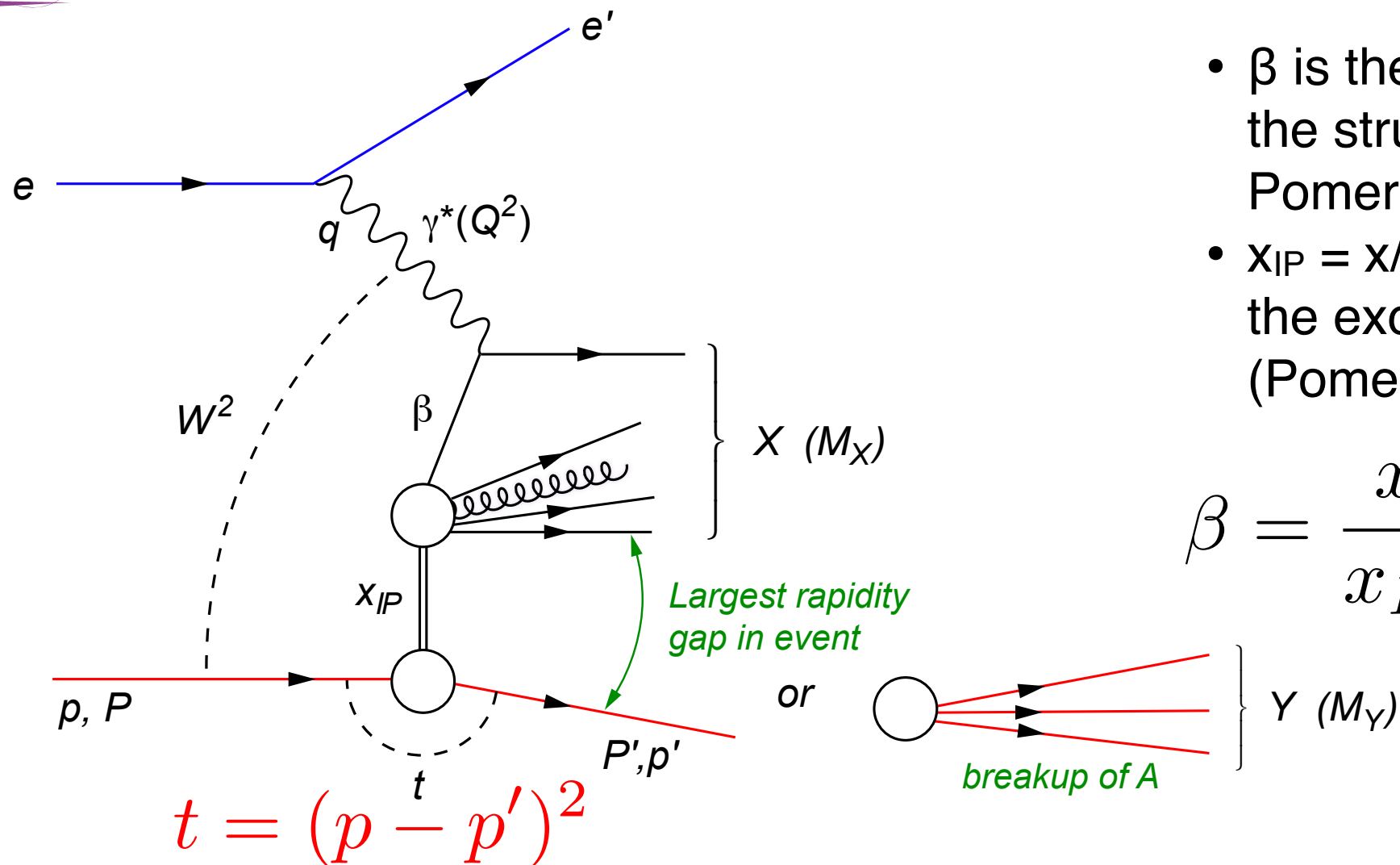
Deliverables	Observables	What we learn	Stage-1	Stage-II
integrated gluon distributions	$F_{2,L}$	nuclear wave function; saturation, Q_s	gluons at $10^{-3} < x < 1$	saturation regime
k_T dependent gluons; gluon correlations	di-hadron correlations	non-linear QCD evolution / universality	onset of saturation	measure Q_s
b-dependent gluons; gluon correlations	DVCS; diffractive vector mesons	interplay between small-x evolution and confinement	moderate x with light, heavy nuclei	smaller x, saturation
transport coefficients in cold matter	large-x SIDIS; jets	parton energy loss, shower evolution; energy loss mechanisms	light flavours and charm; jets	rare probes and bottom; large-x gluons



Important Measurements

Deliverables	Observables	What we learn	Stage-1	Stage-II
integrated gluon distributions	$F_{2,L}$	nuclear wave function; saturation, Q_s	gluons at $10^{-3} < x < 1$	saturation regime
k_T dependent gluons; gluon correlations	di-hadron correlations	non-linear QCD evolution / universality	onset of saturation	measure Q_s
b-dependent gluons; gluon correlations	DVCS; diffractive vector mesons	interplay between small-x evolution and confinement	moderate x with light, heavy nuclei	smaller x, saturation
transport coefficients in cold matter	large-x SIDIS; jets	parton energy loss, shower evolution; energy loss mechanisms	light flavours and charm; jets	rare probes and bottom; large-x gluons

Hard Diffraction



- β is the momentum fraction of the struck parton w.r.t. the Pomeron
- $x_{IP} = x/\beta$: momentum fraction of the exchanged object (Pomeron) w.r.t. the hadron

$$\beta = \frac{x}{x_{IP}} = \frac{Q^2}{Q^2 + M_X^2 - t}$$

• Diffraction in e+p:

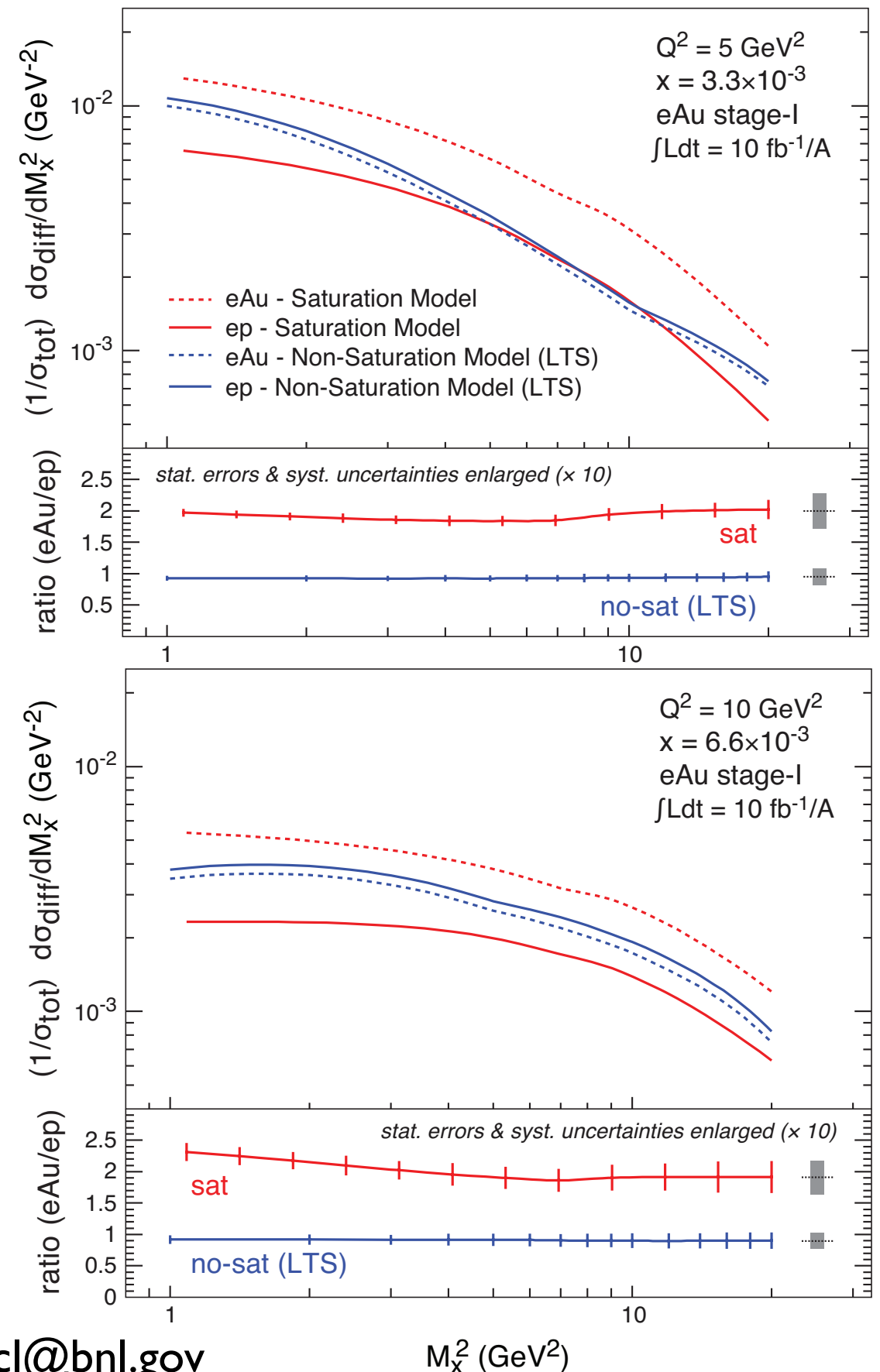
- ➔ HERA: 15% of all events are diffractive

• Diffraction in e+A:

- ➔ Predictions: $\sigma_{\text{diff}}/\sigma_{\text{tot}}$ in e+A ~25-40%
- ➔ Coherent diffraction (nuclei intact)
- ➔ Incoherent diffraction: breakup into nucleons (nucleons intact)

Diffractive cross-sections: Saturation vs Non-Sat

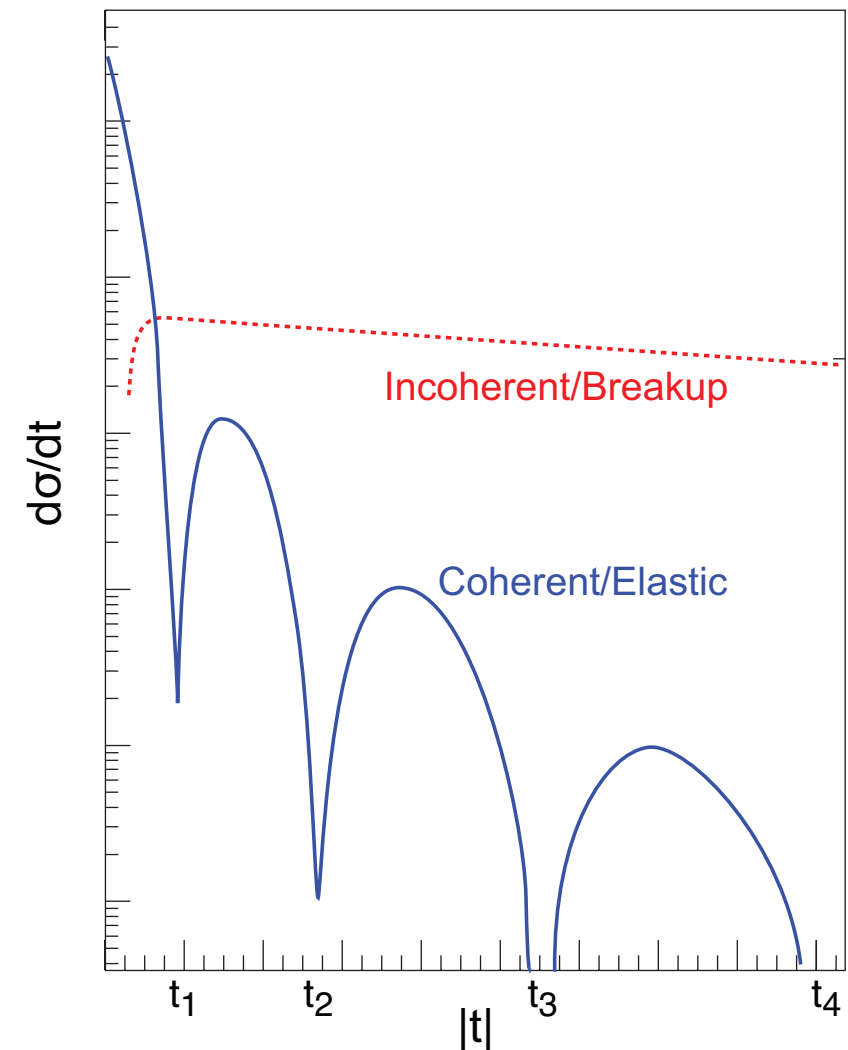
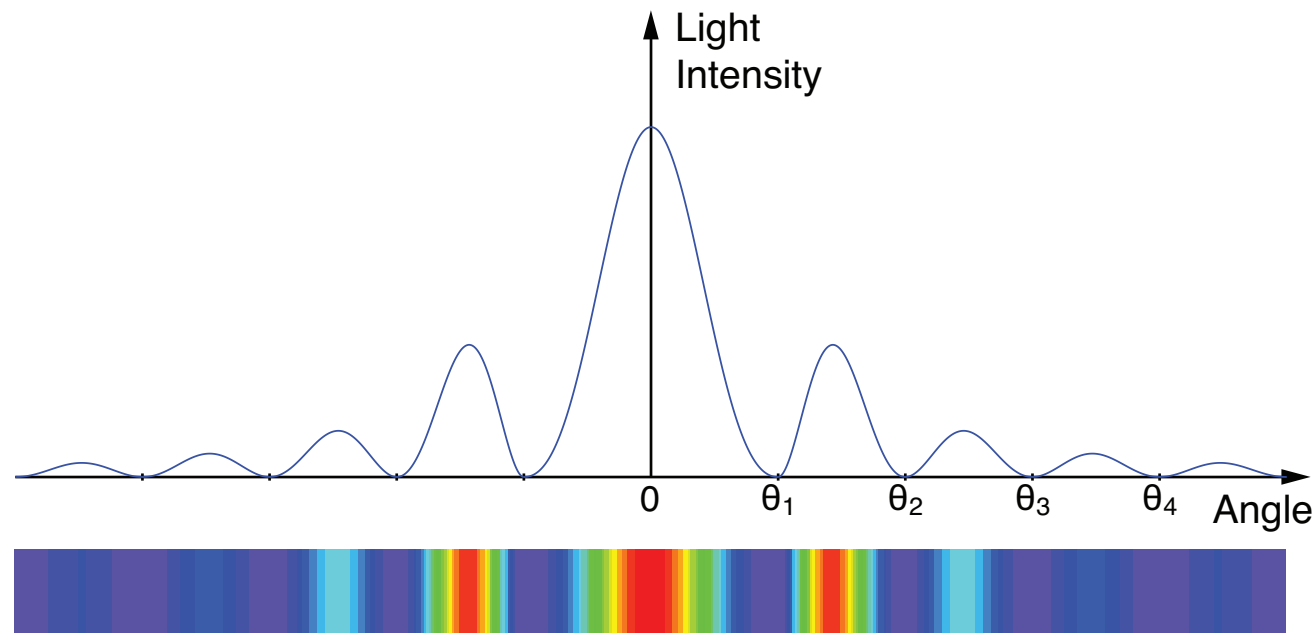
- Ratio of diffractive to total cross-sections between **saturation model (Marquet)** and **Leading-Twist Shadowing (Guzey, Strickman)**.
 - ➔ Very little difference for LTS between e+p and e+Au, independent of Q^2
 - ➔ For saturation model, e+Au $\sim 2 \cdot$ e+p, again independent of Q^2
 - ➔ Simulated error bars (10 fb^{-1}) can easily distinguish between these two scenarios
 - ▶ Note that the errors are scaled on the plot so they are visible!





Exclusive Vector Meson Production in $e+A$

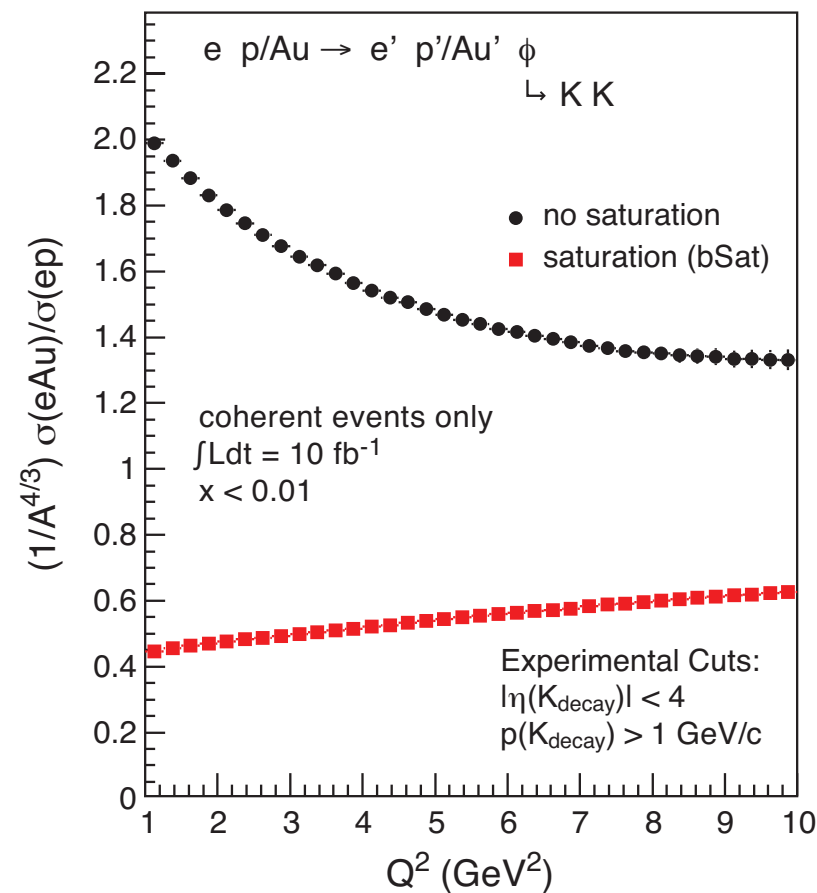
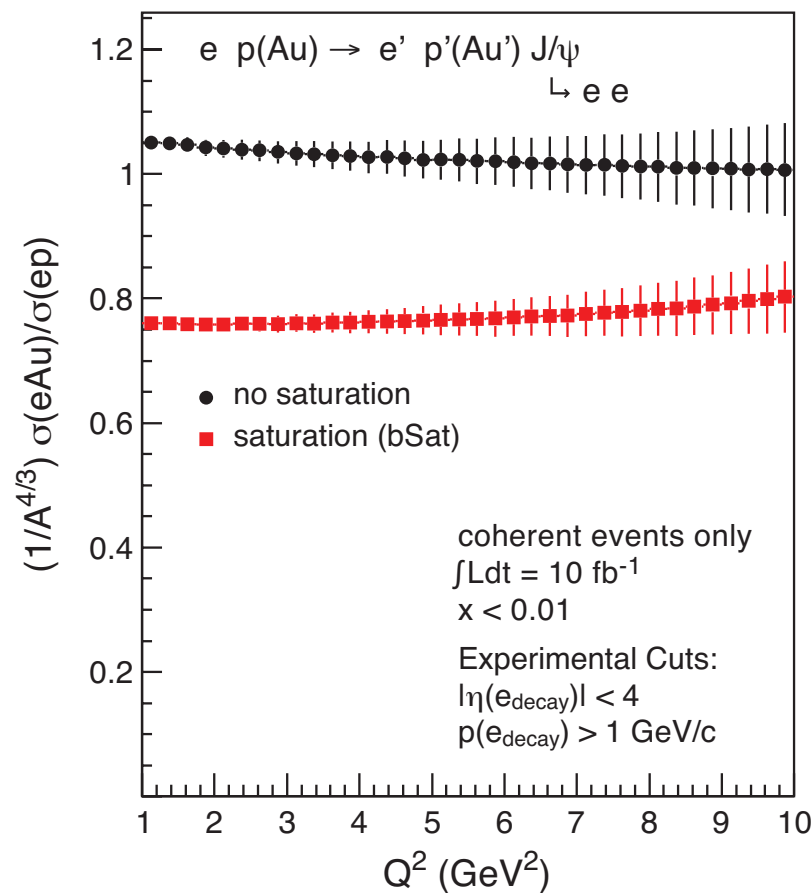
$$e+A \rightarrow e+J/\psi+A'$$



- Low- t : coherent diffraction dominates - **gluon density**
- High- t : incoherent diffraction dominates - **gluon correlations**
- Just like in optics - the positions of the diffractive minima are related to the size of the obstacle

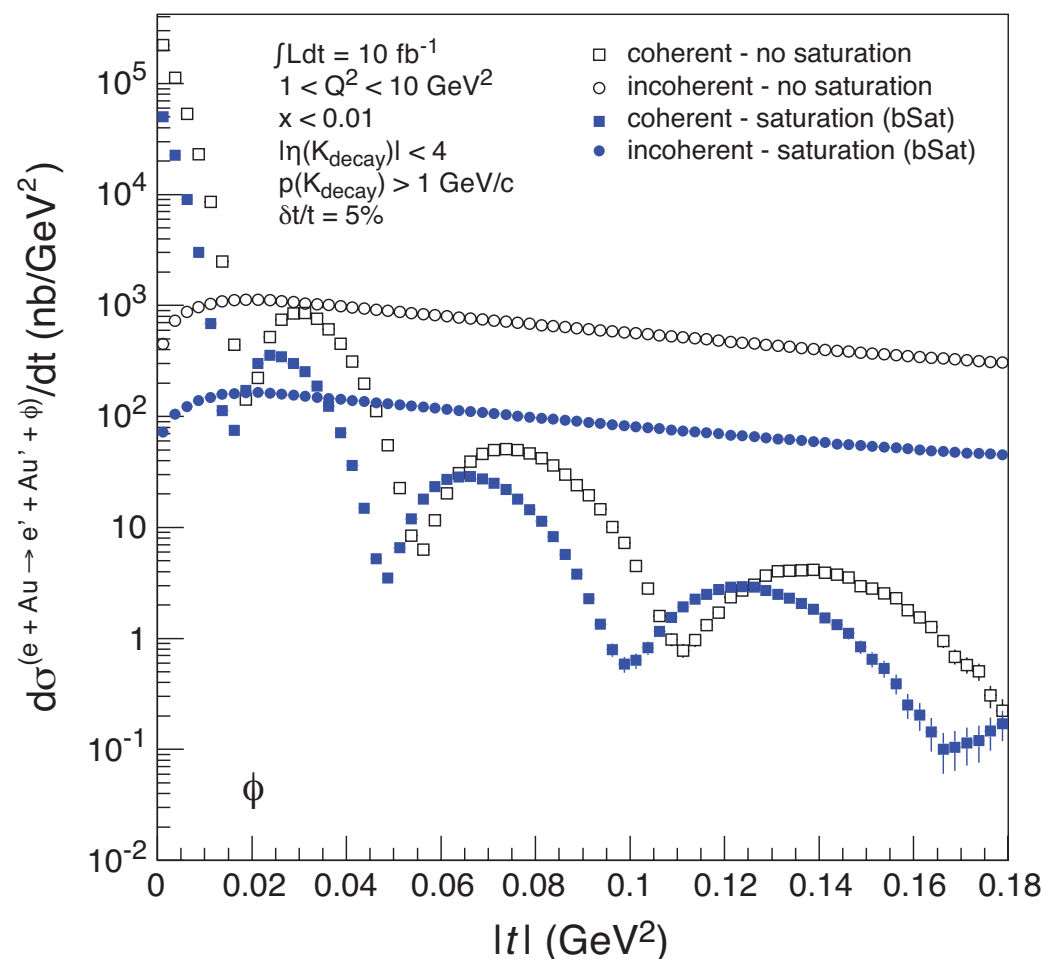
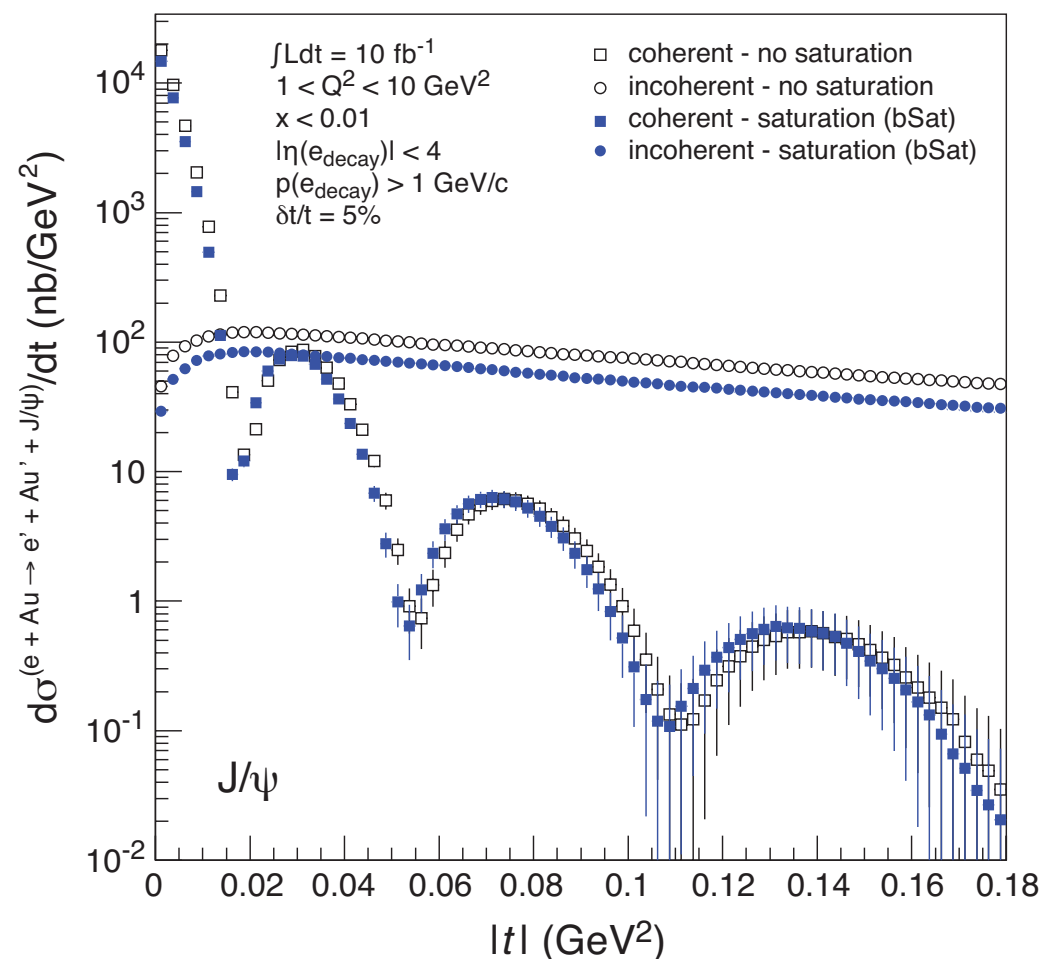
$$\Rightarrow \theta_i \sim 1/(kR)$$

Exclusive Vector Meson Production in e+A



- Diffraction with final state VM - plots from Sartre event generator
 - ➔ Clean - only one new final state particle generated
 - ➔ Unambiguously identified via the presence of a rapidity gap
 - ➔ J/ψ less sensitive to saturation effects than ϕ
 - expected as ϕ has larger wave function

Exclusive Vector Meson Production in e+A



- Low-t: coherent diffraction dominates - **gluon density**
- High-t: incoherent diffraction dominates - **gluon correlations**
- ➔ Need good breakup detection efficiency to discriminate between the two scenarios
 - unlike protons, forward spectrometer won't work for heavy ions
 - **measure emitted neutrons in a ZDC**
 - rapidity gap with absence of break-up fragments sufficient to identify coherent events

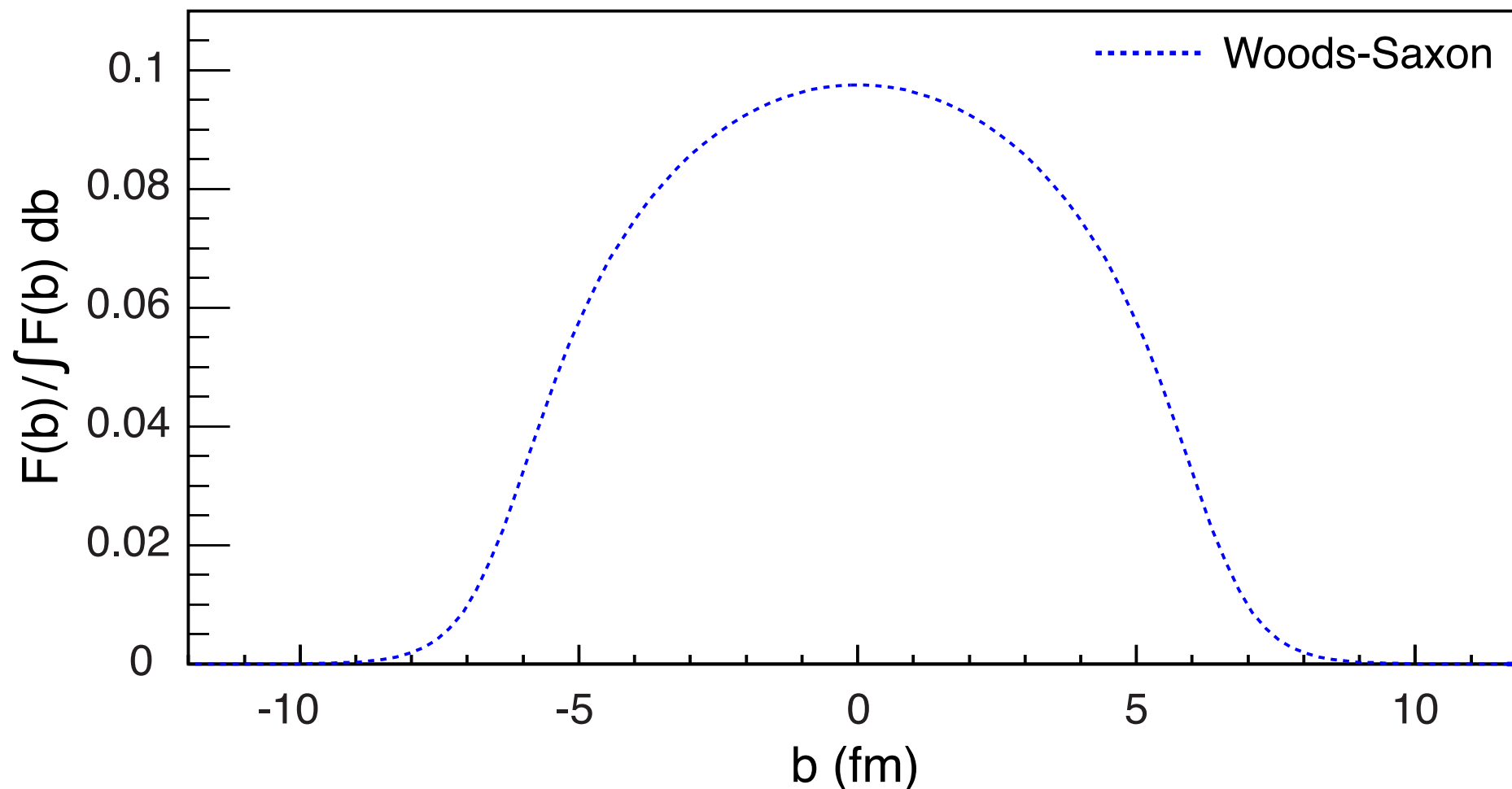
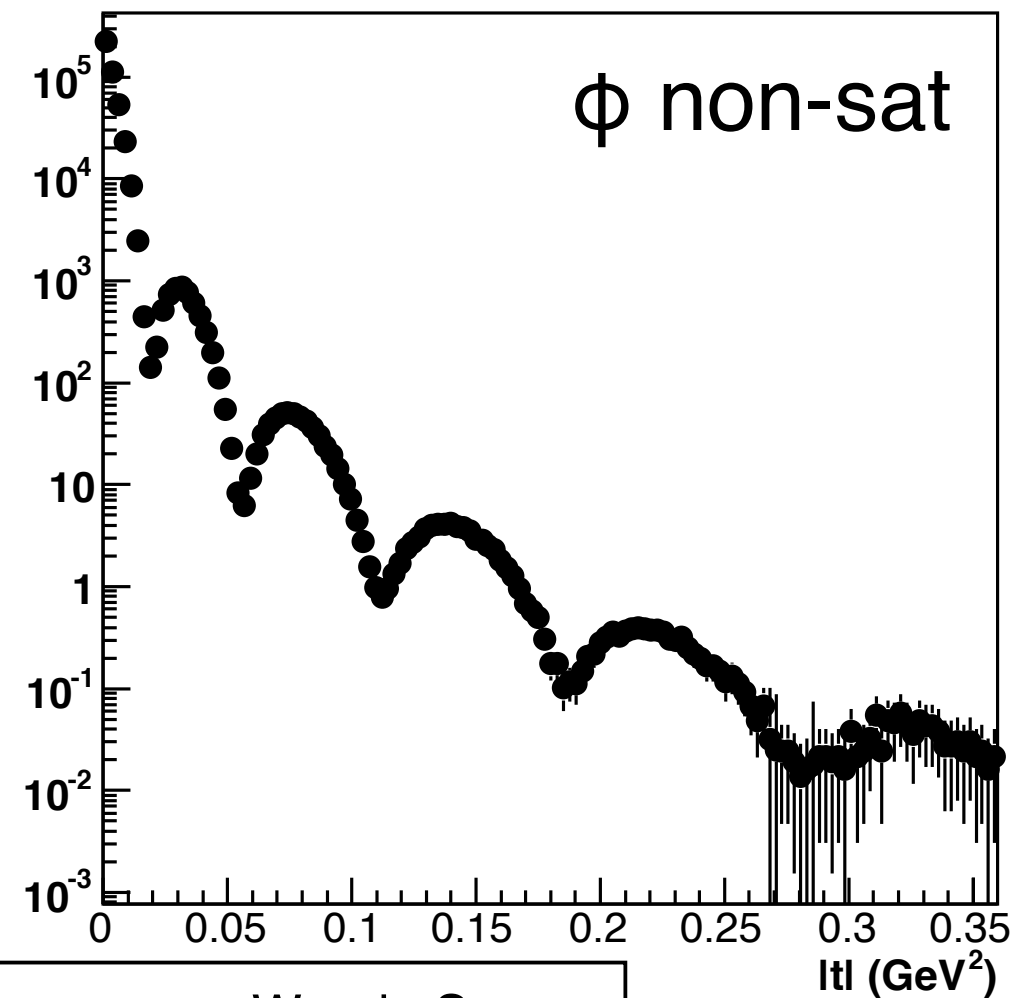


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)

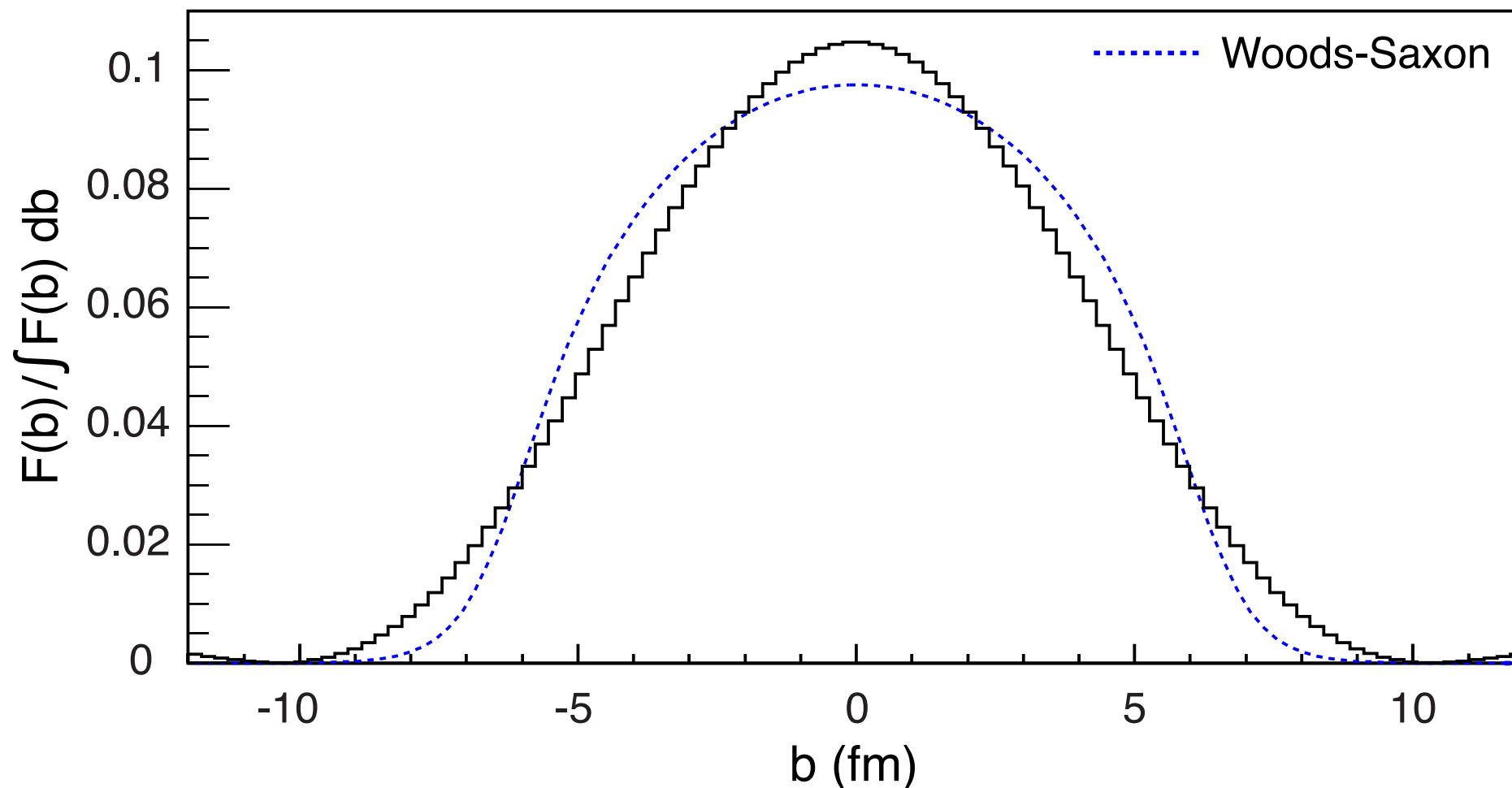
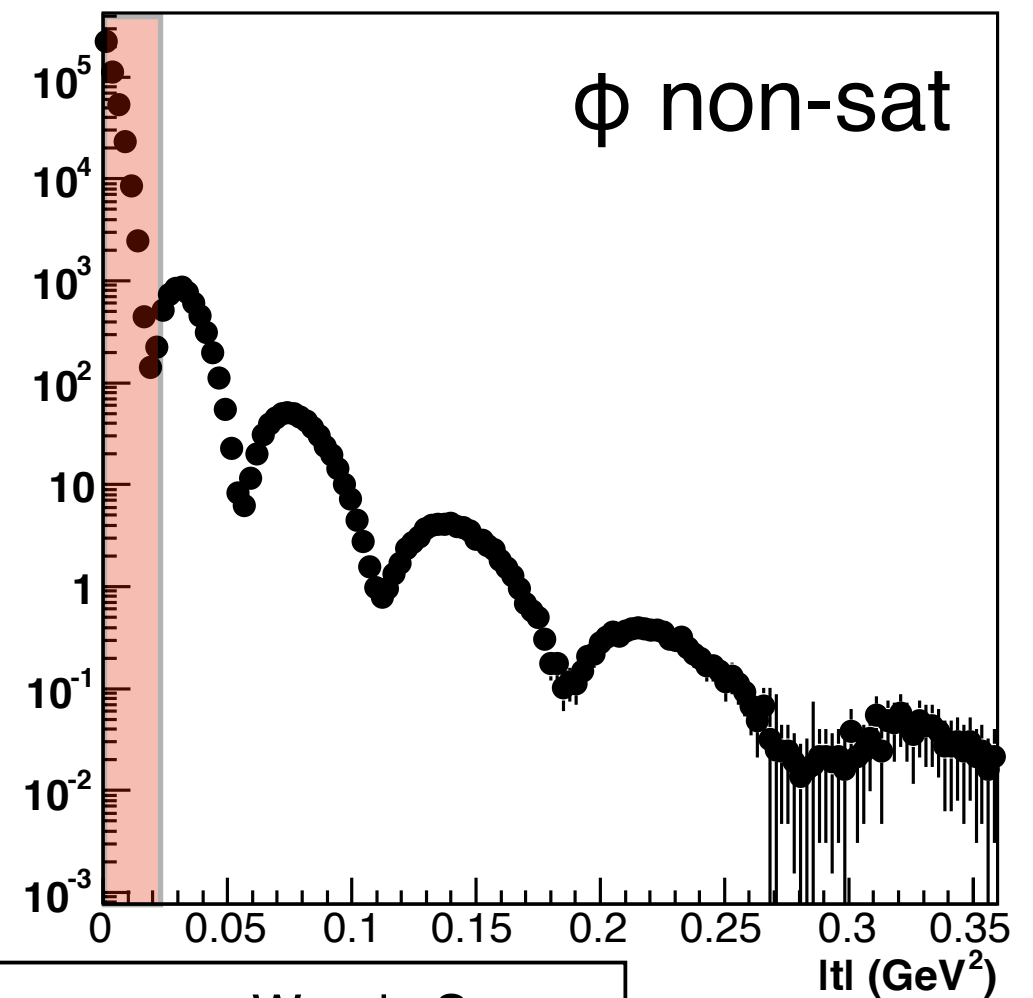


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)

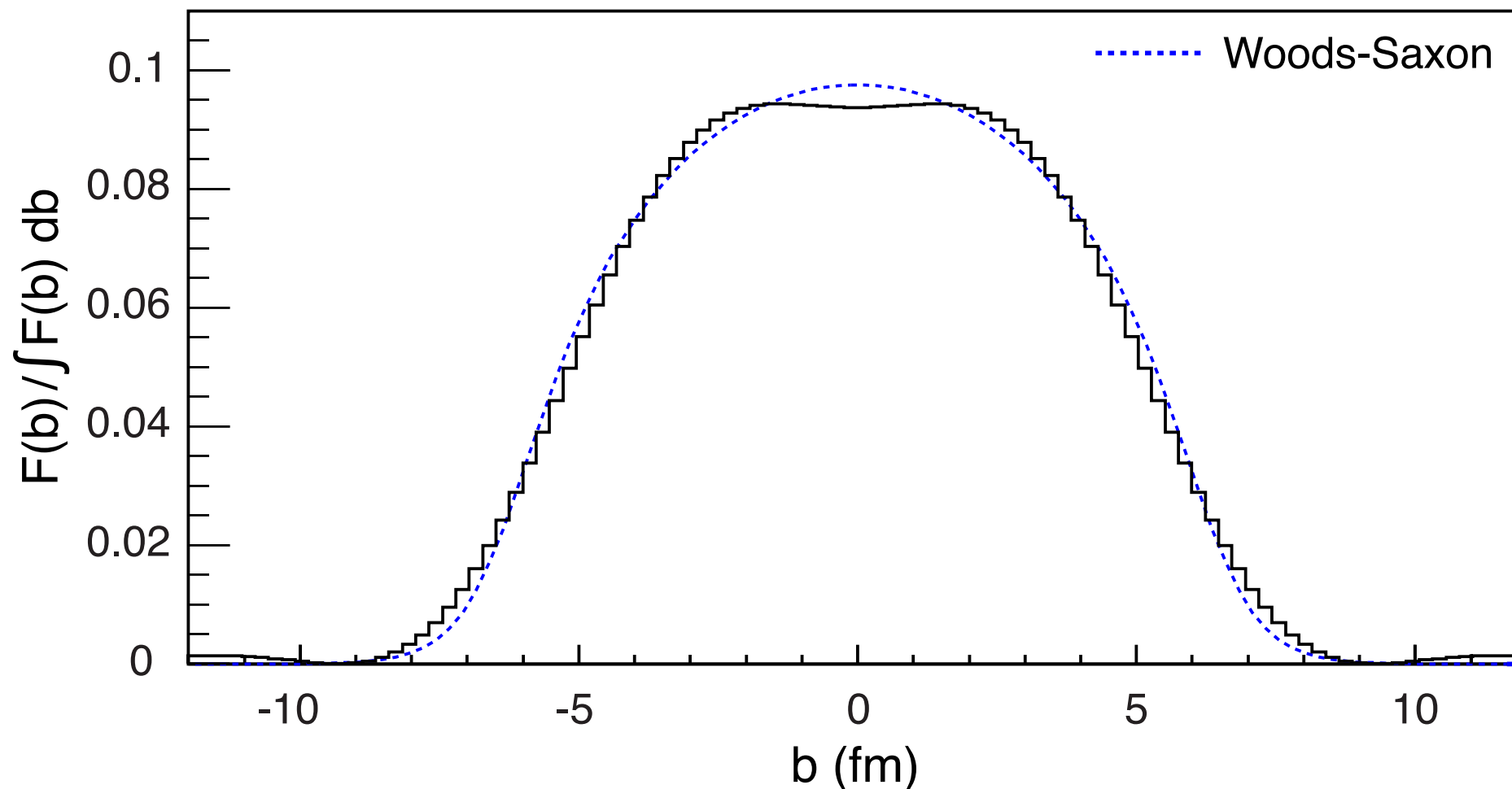
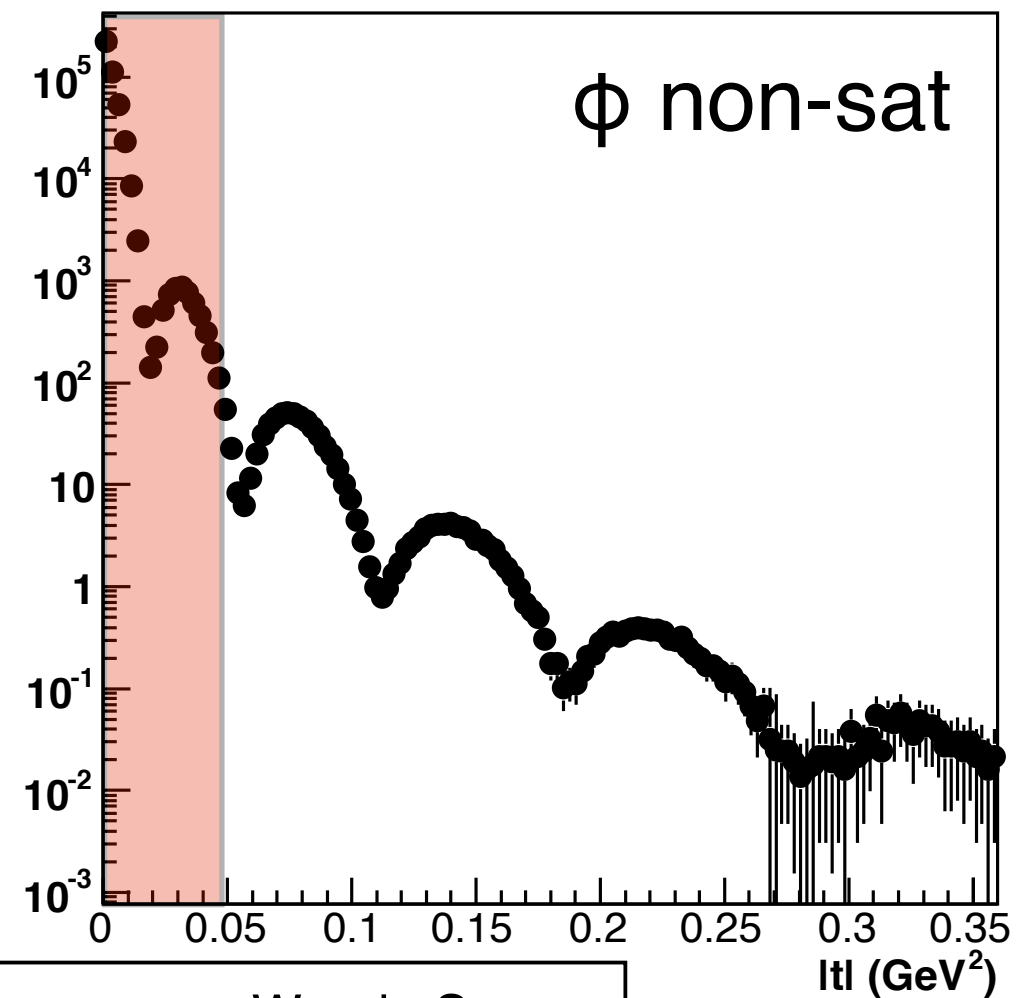


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)

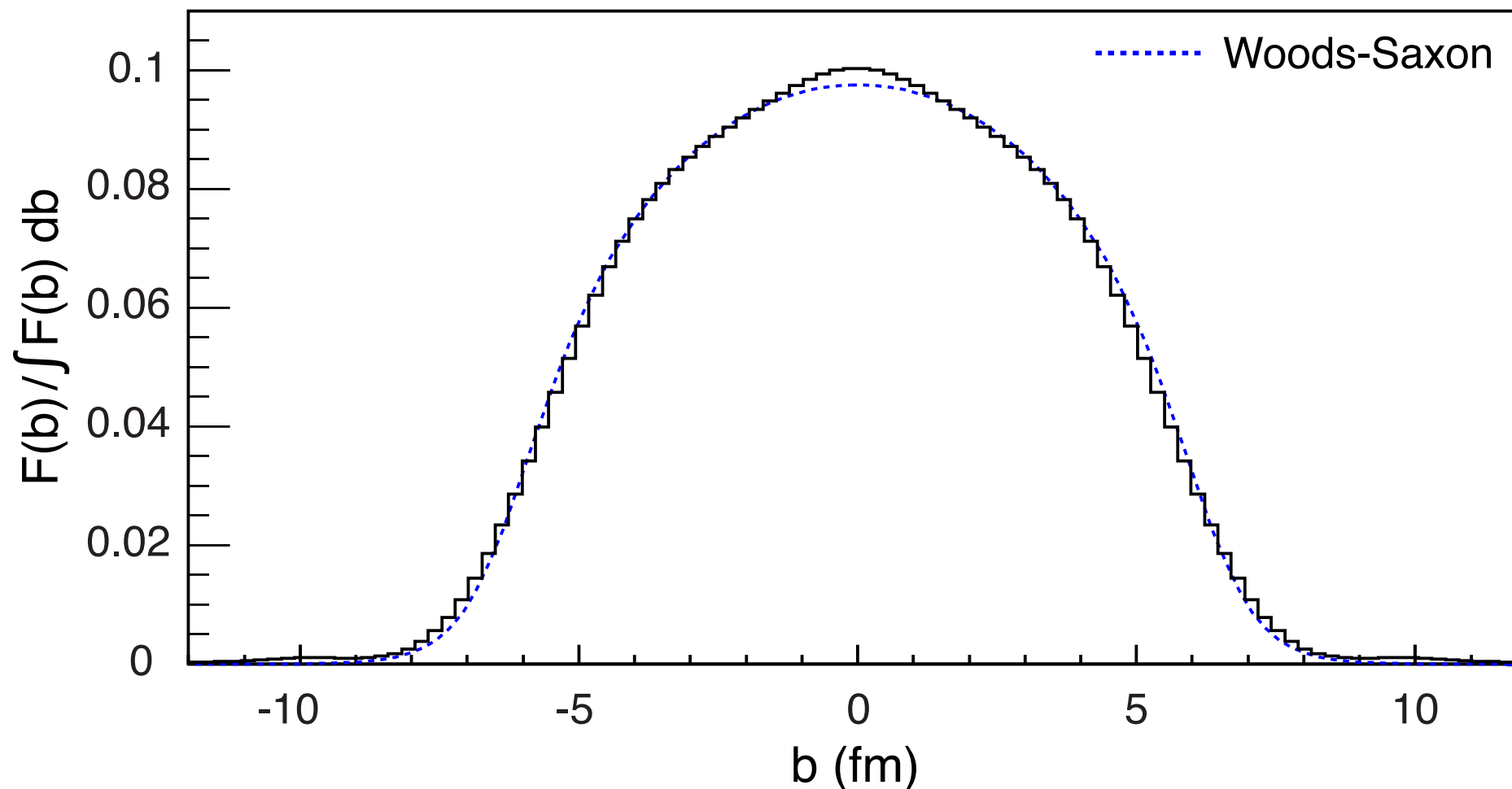
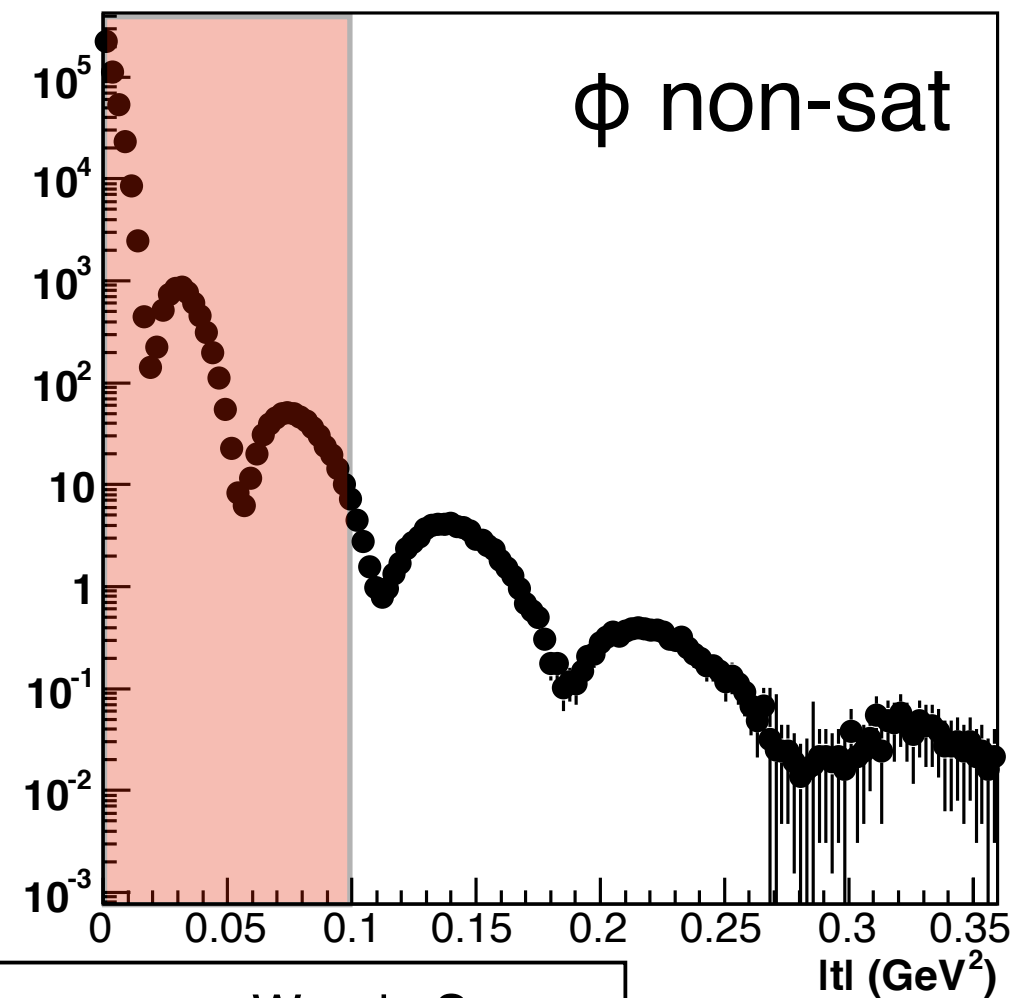


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)

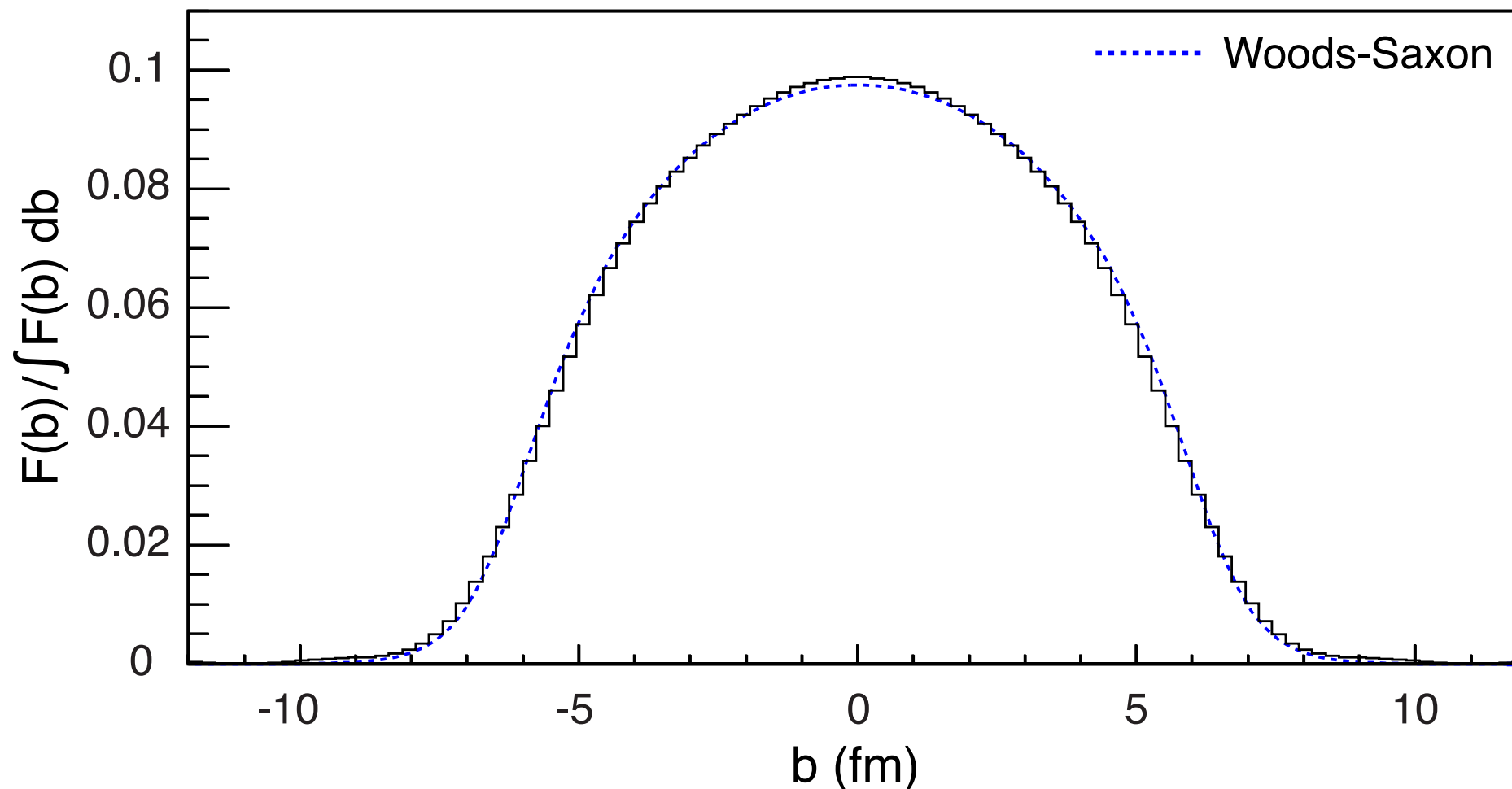
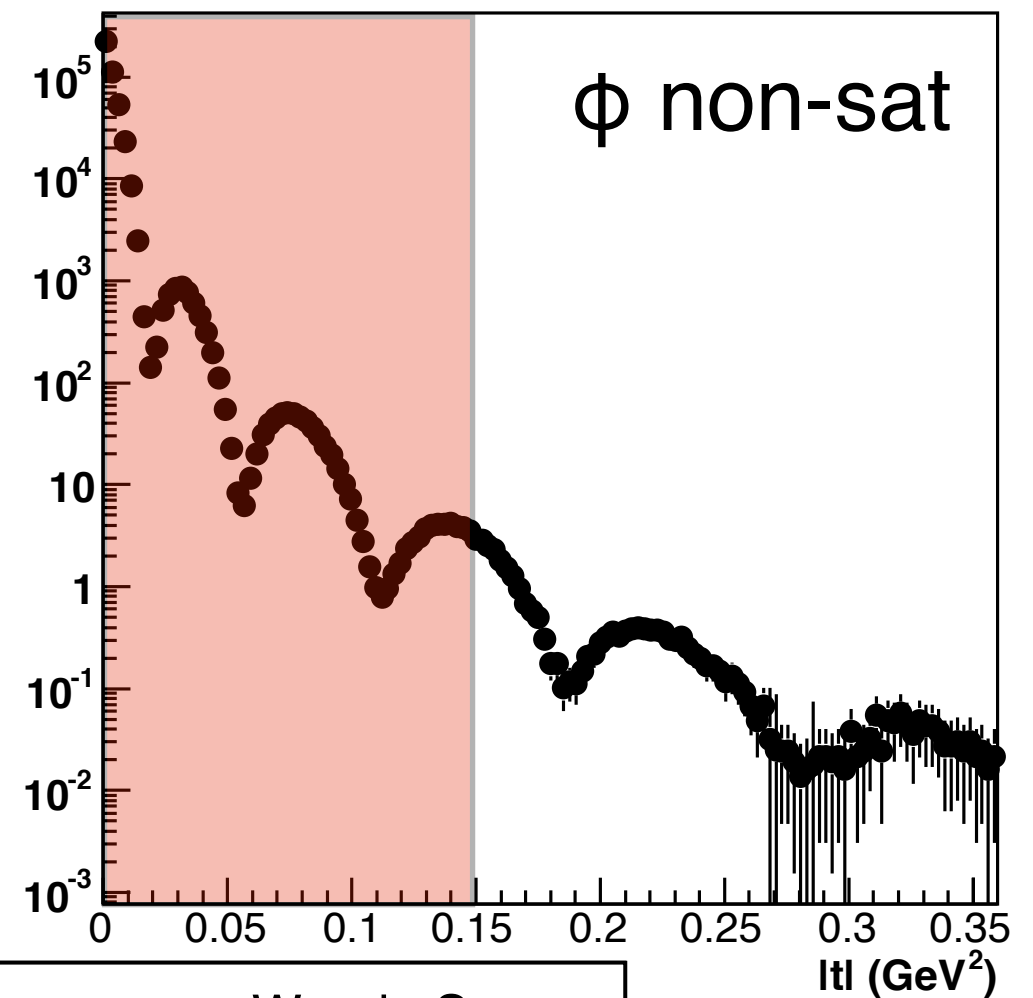


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)

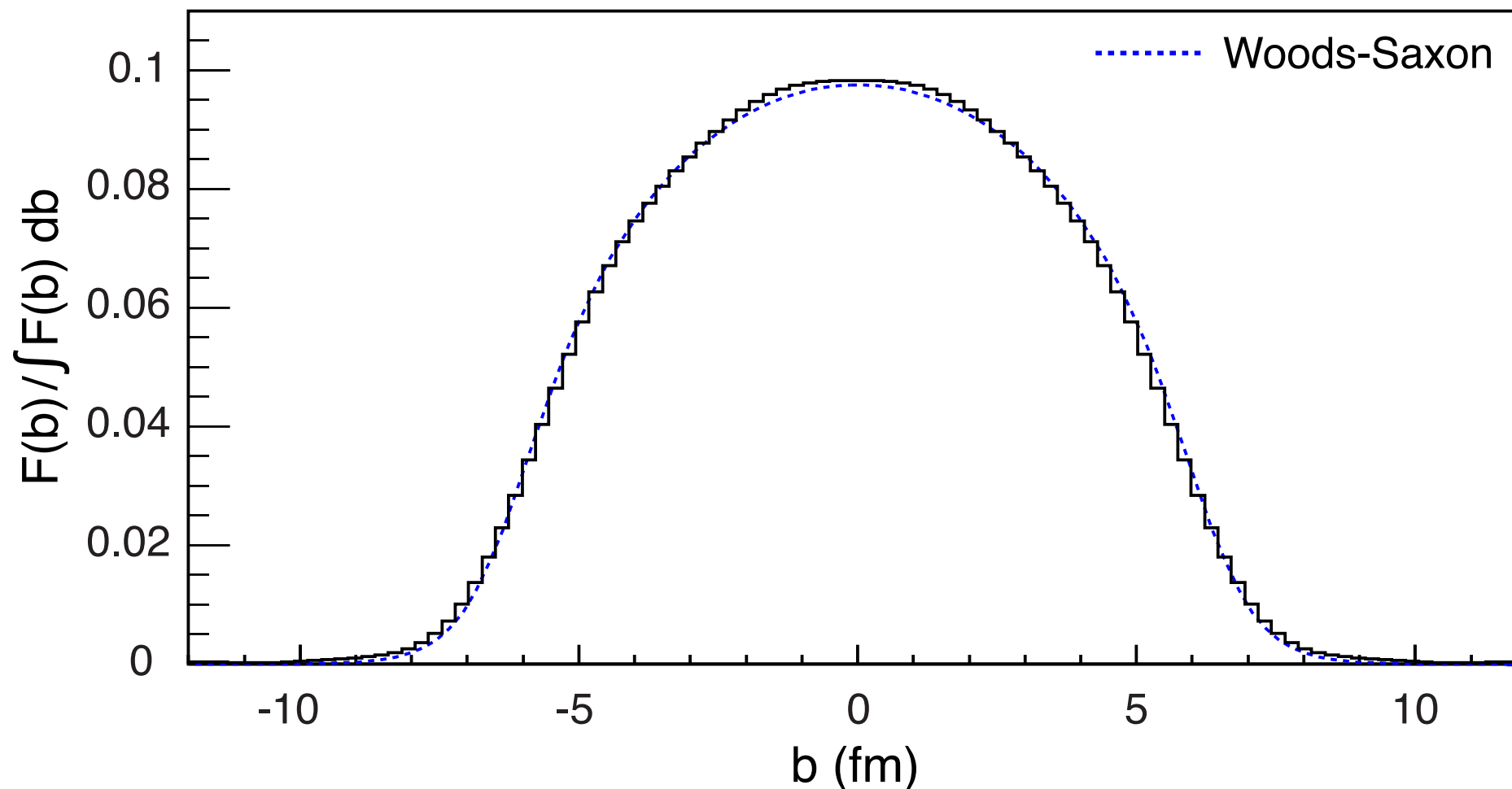
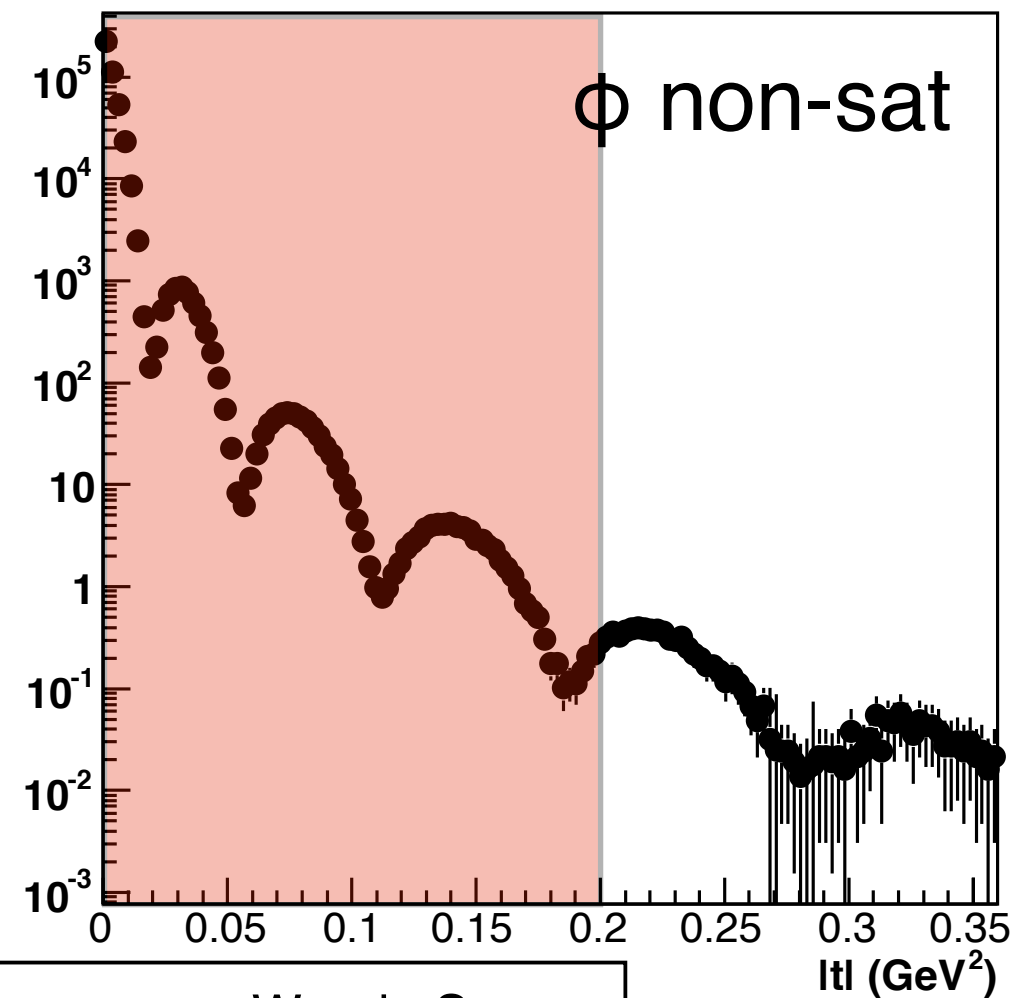


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)

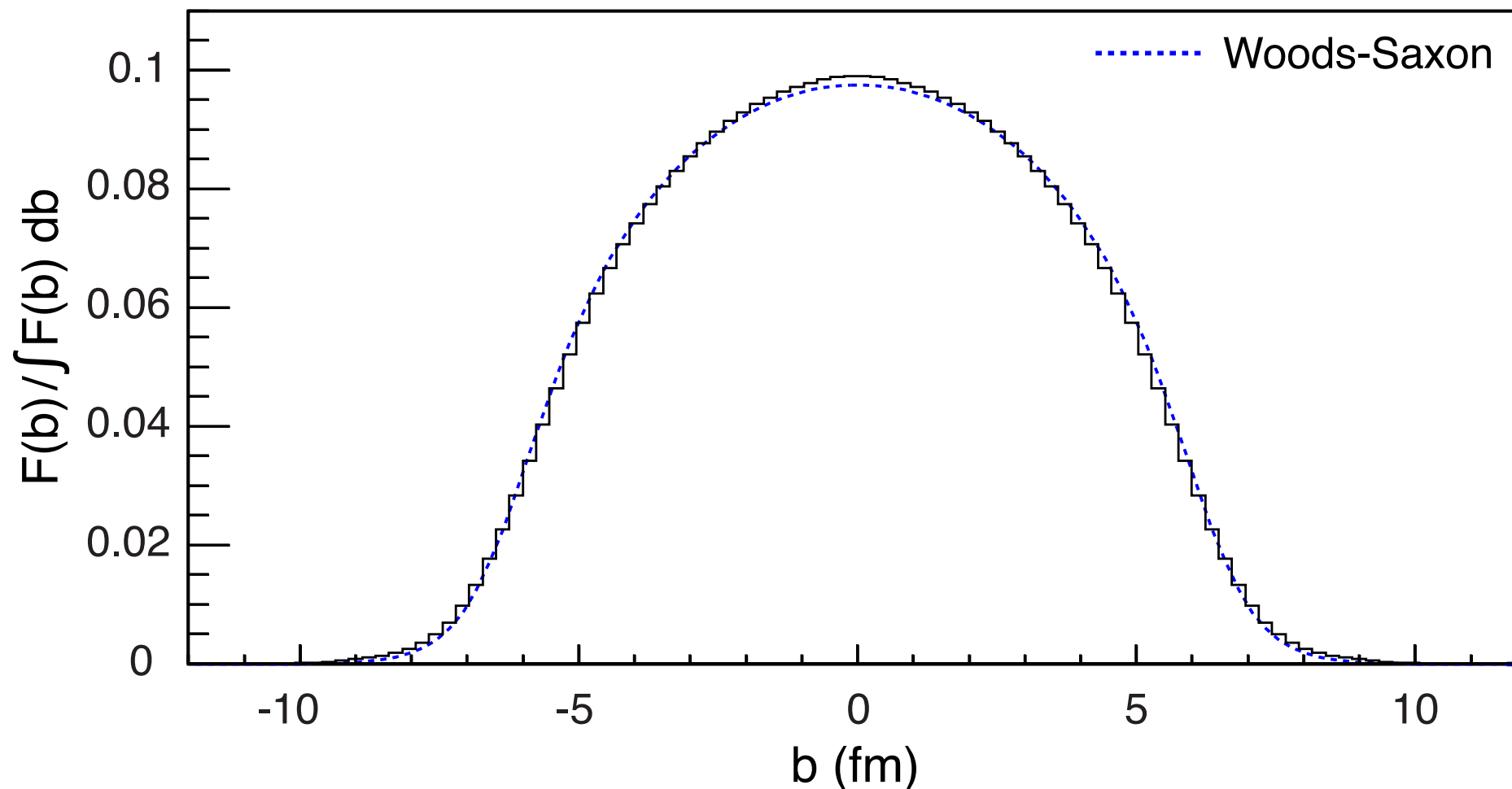
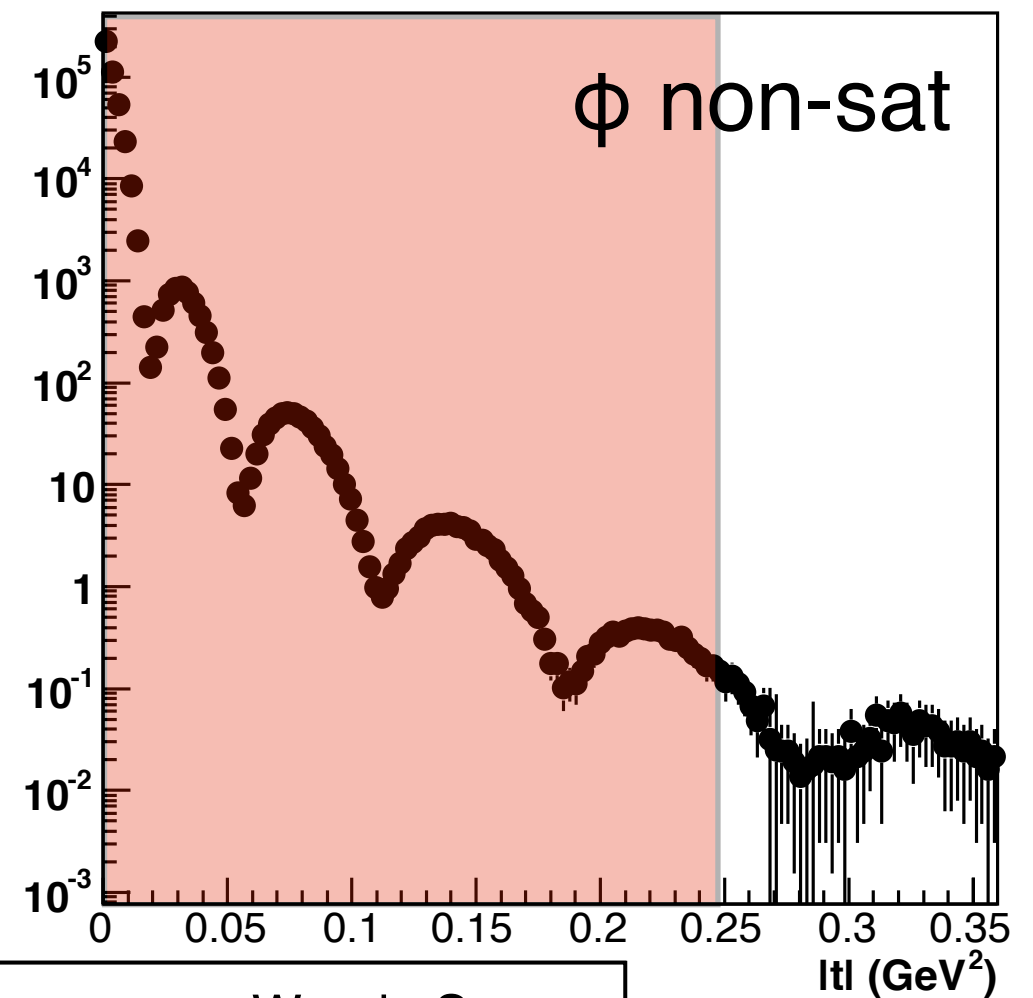


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)



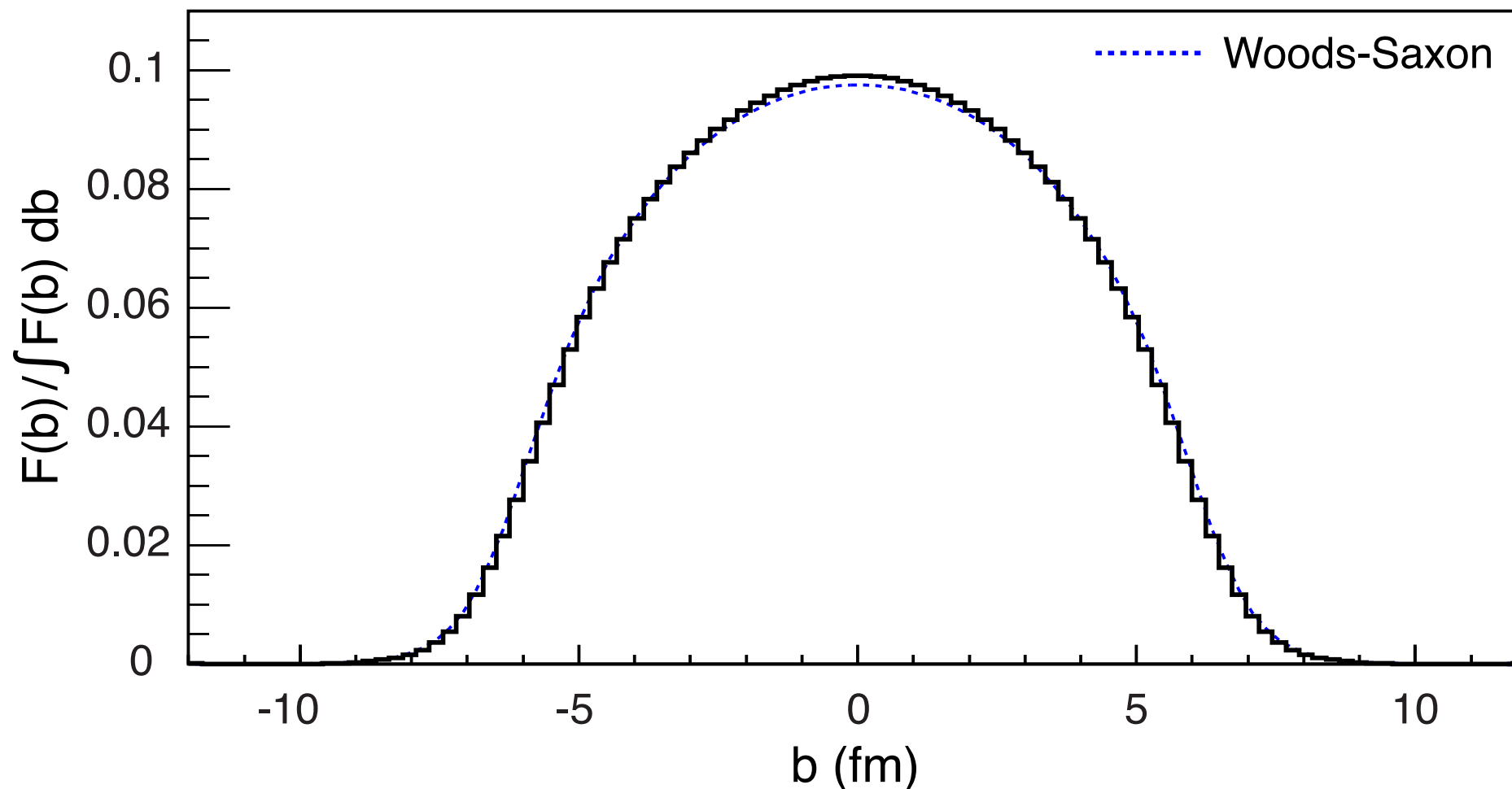
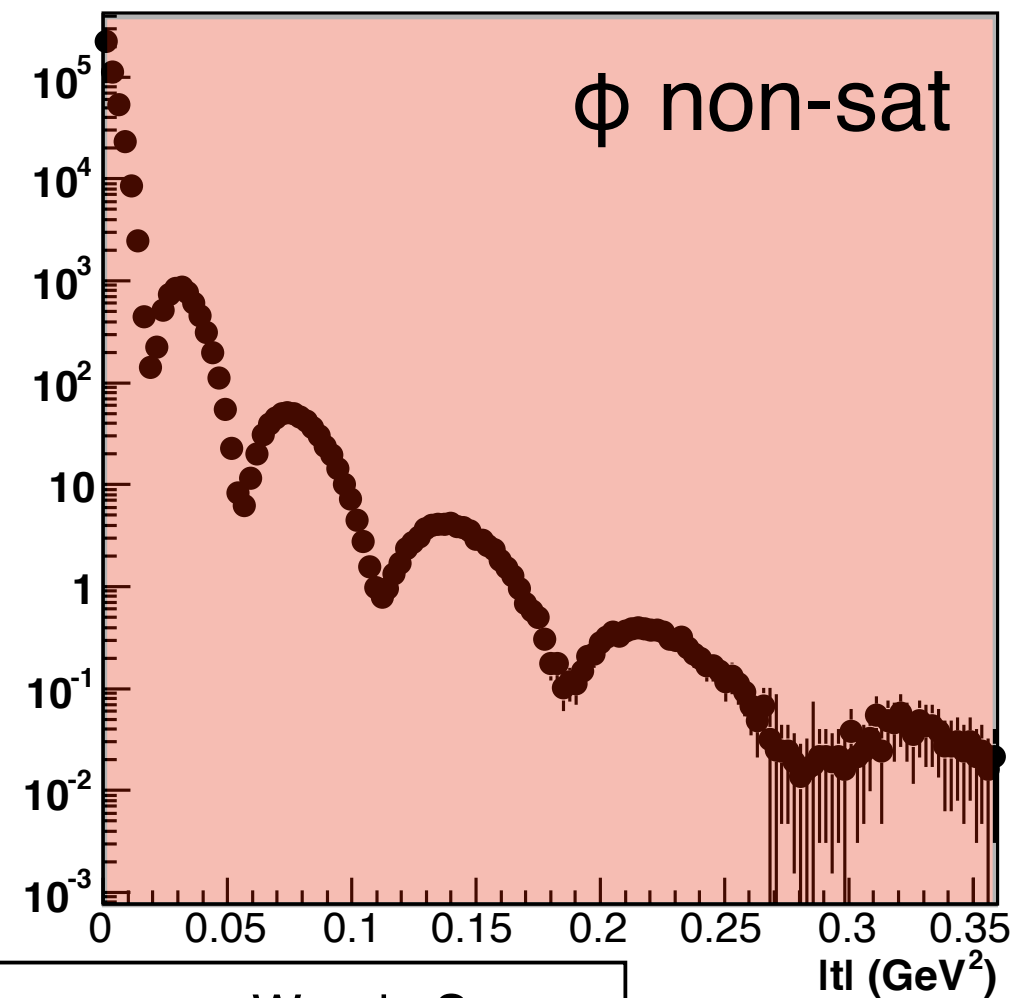


Finding the source...

- Take the $d\sigma/dt$ distribution and perform a Fourier Transform to extract the b -distribution of the gluons

$$F(b) \sim \frac{1}{2\pi} \int_0^\infty d\Delta \Delta J_0(\Delta b) \sqrt{\frac{d\sigma}{dt}}$$

$t = \Delta^2/(1-x) \approx \Delta^2$ (for small x)



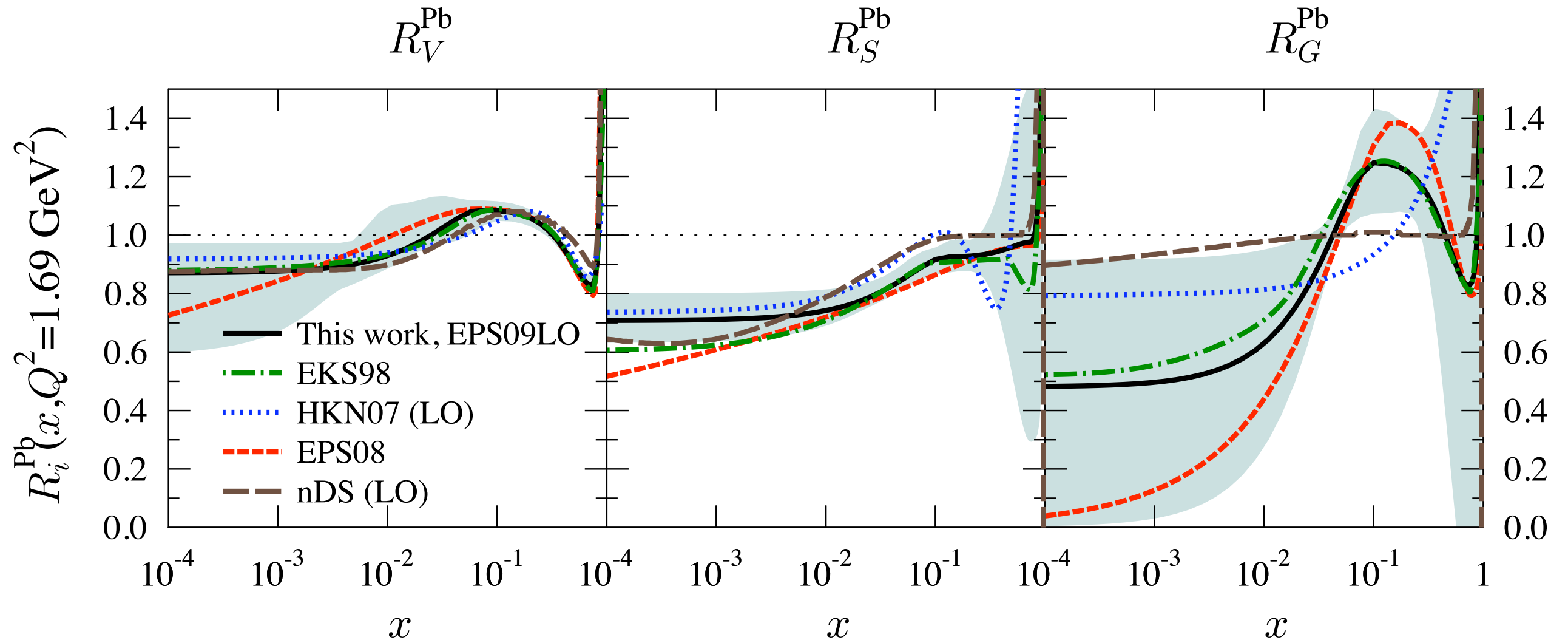


Summary/Conclusions

- The e+A physics programme at an EIC will give us an unprecedented opportunity to study gluons in nuclei at small-x
 - ➔ We can measure the properties of gluons where saturation is the dominant governing phenomena
 - ➔ Understanding the role of gluons in nuclei is also crucial to a quantitative understanding of RHIC (and LHC) heavy-ion results
- Diffractive collisions will give us a good handle on the gluon distribution
 - ➔ The ratio of the cross-sections of diffractive collisions in e+A/e+p itself will help differentiate between linear and non-linear effects
 - ➔ A Fourier Transform of the coherent $d\sigma/dt$ distribution of diffractive vector mesons can allow us to extract the b-dependent initial gluon distribution

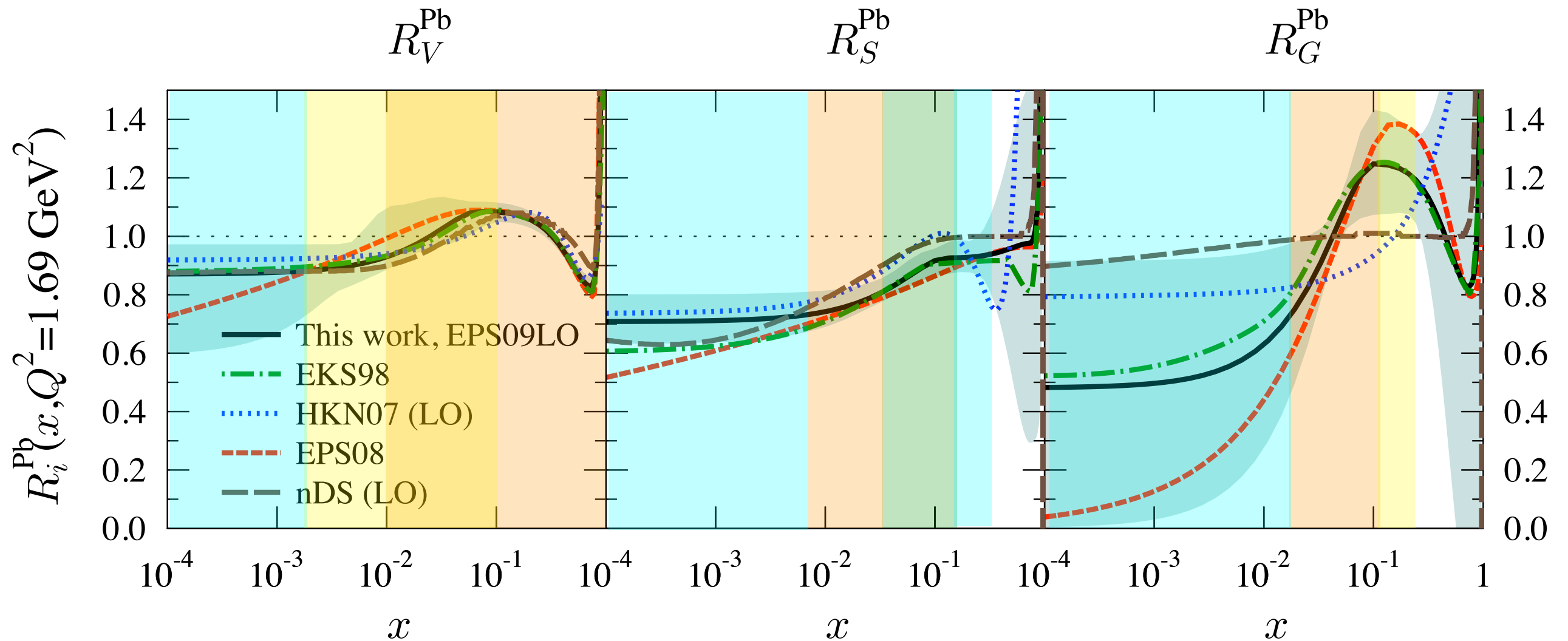
BACKUP

What do we know about the structure of nuclei?



The distribution of valence and sea quarks are relatively well known in nuclei - theories agree well

What do we know about the structure of nuclei?



Constrained by DIS

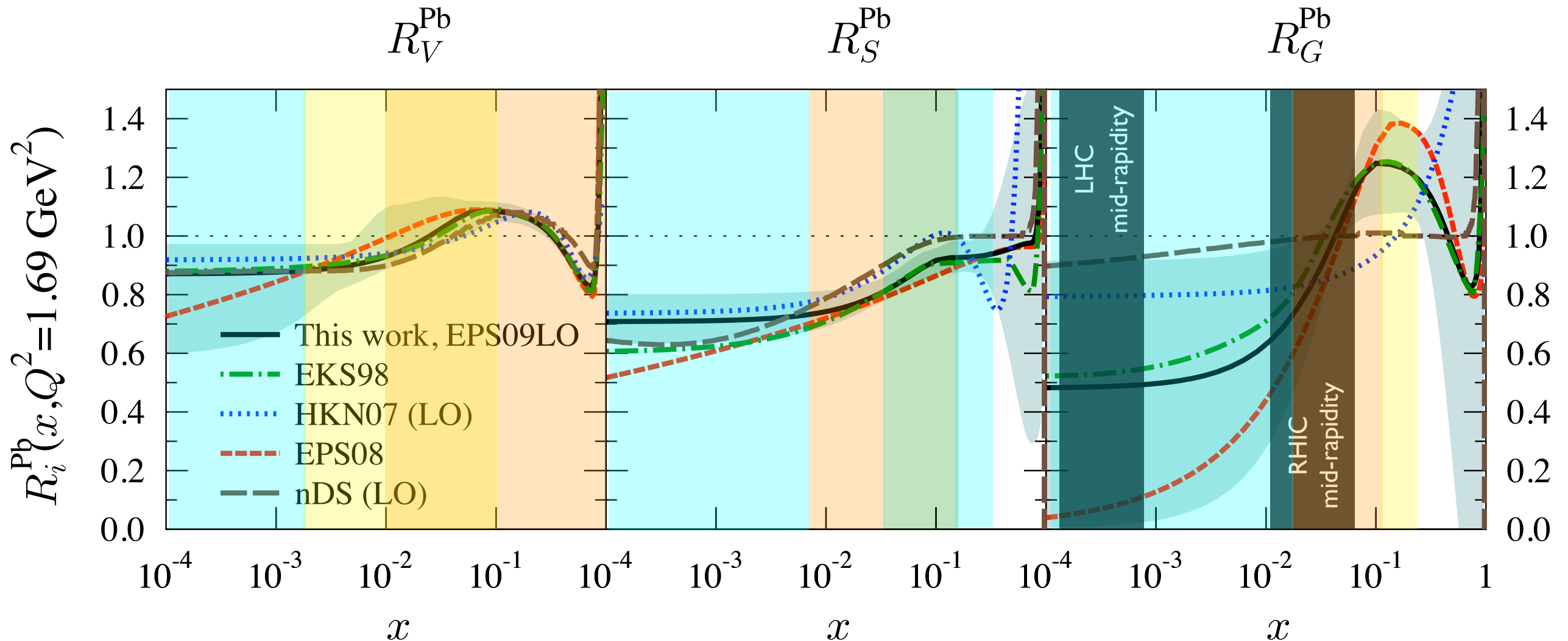
Constrained by DY

Constrained by sum rules

Assumptions

The distribution of valence and sea quarks are relatively well known in nuclei - theories agree well

What do we know about the structure of nuclei?



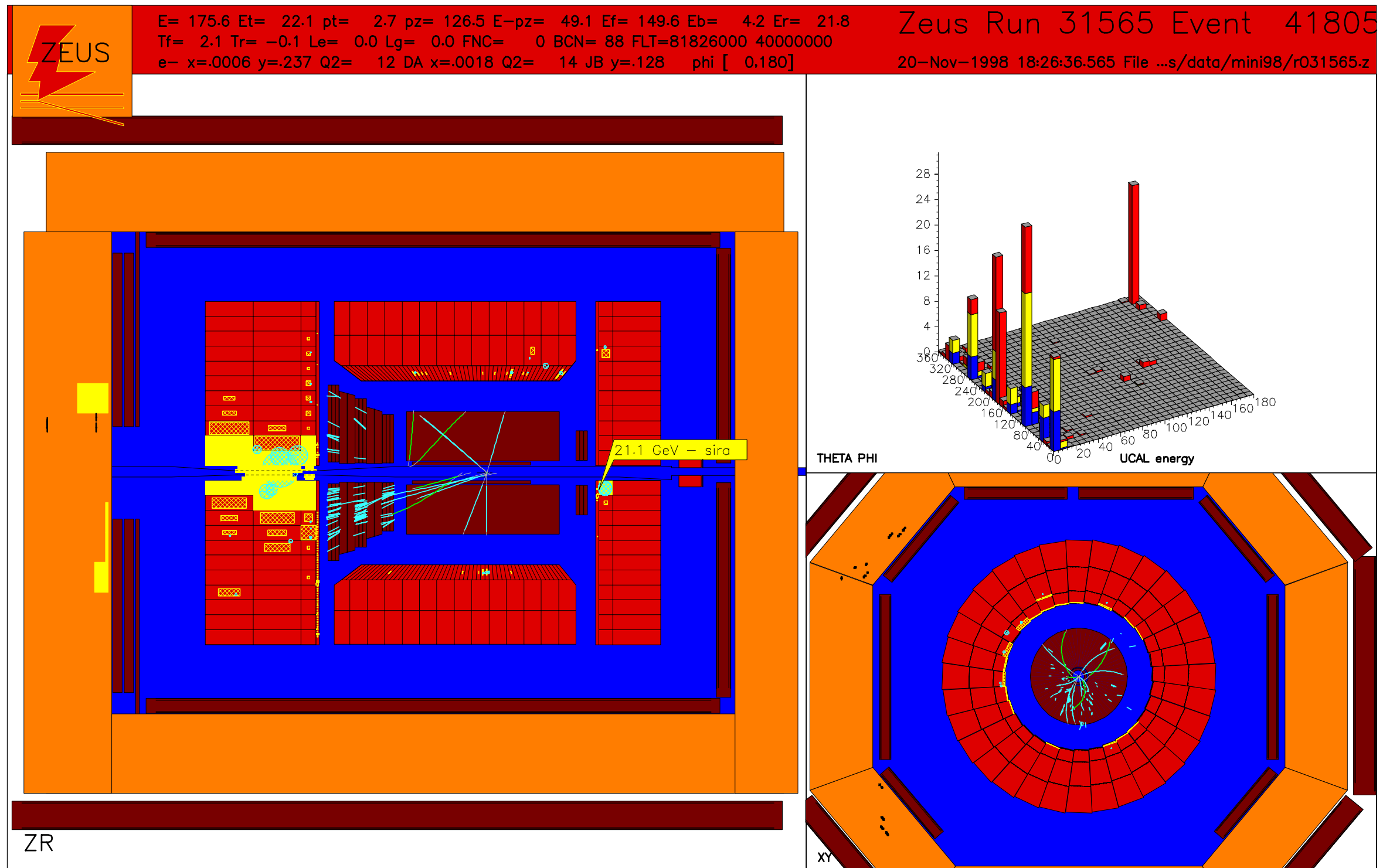
- Constrained by DIS
- Constrained by DY
- Constrained by sum rules
- Assumptions

The distribution of valence and sea quarks are relatively well known in nuclei - theories agree well

Large discrepancies exist in the gluon distributions from models for mid-rapidity LHC and forward RHIC rapidities !!

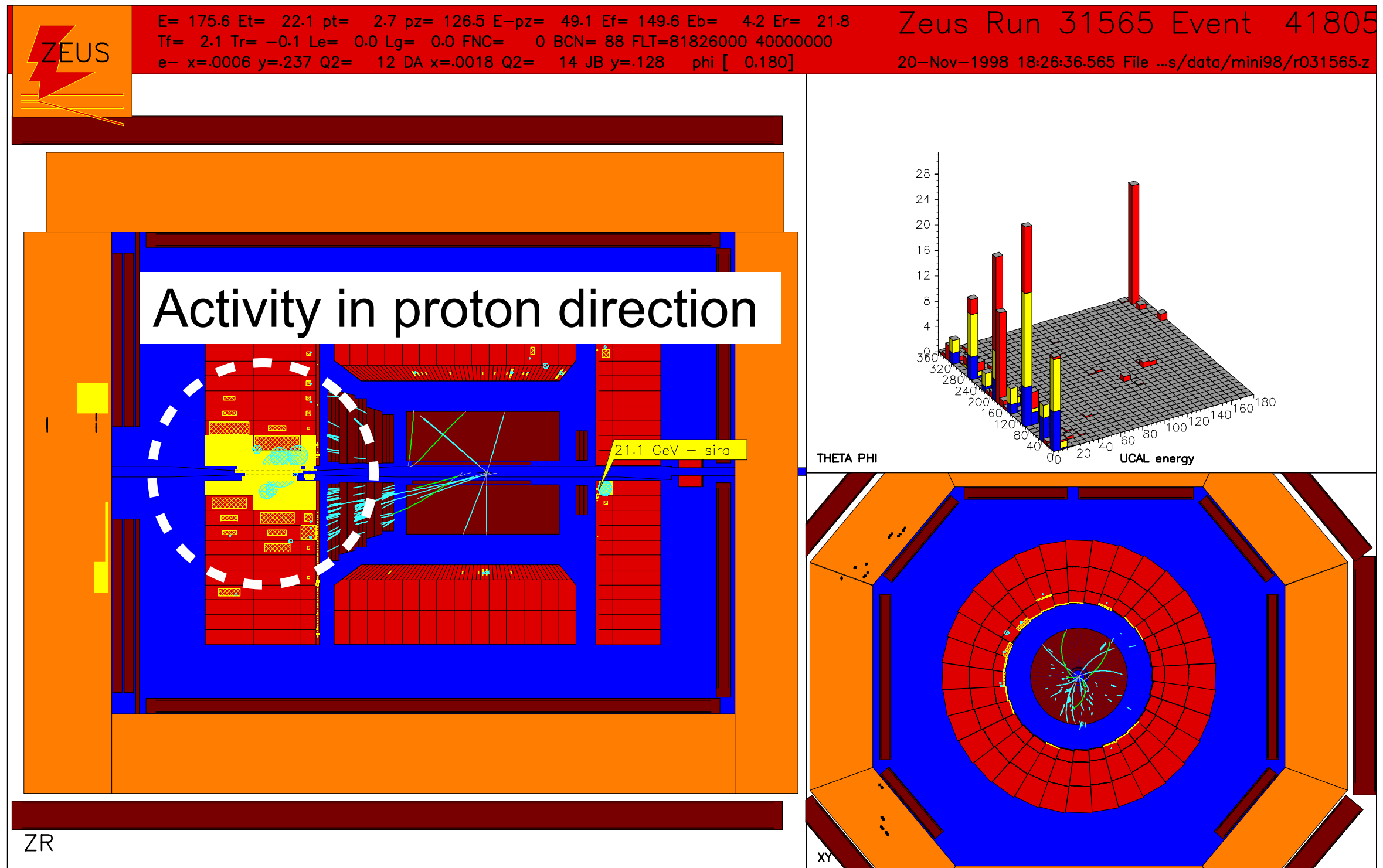
Visualising Diffractive events

A DIS event (experimental view)

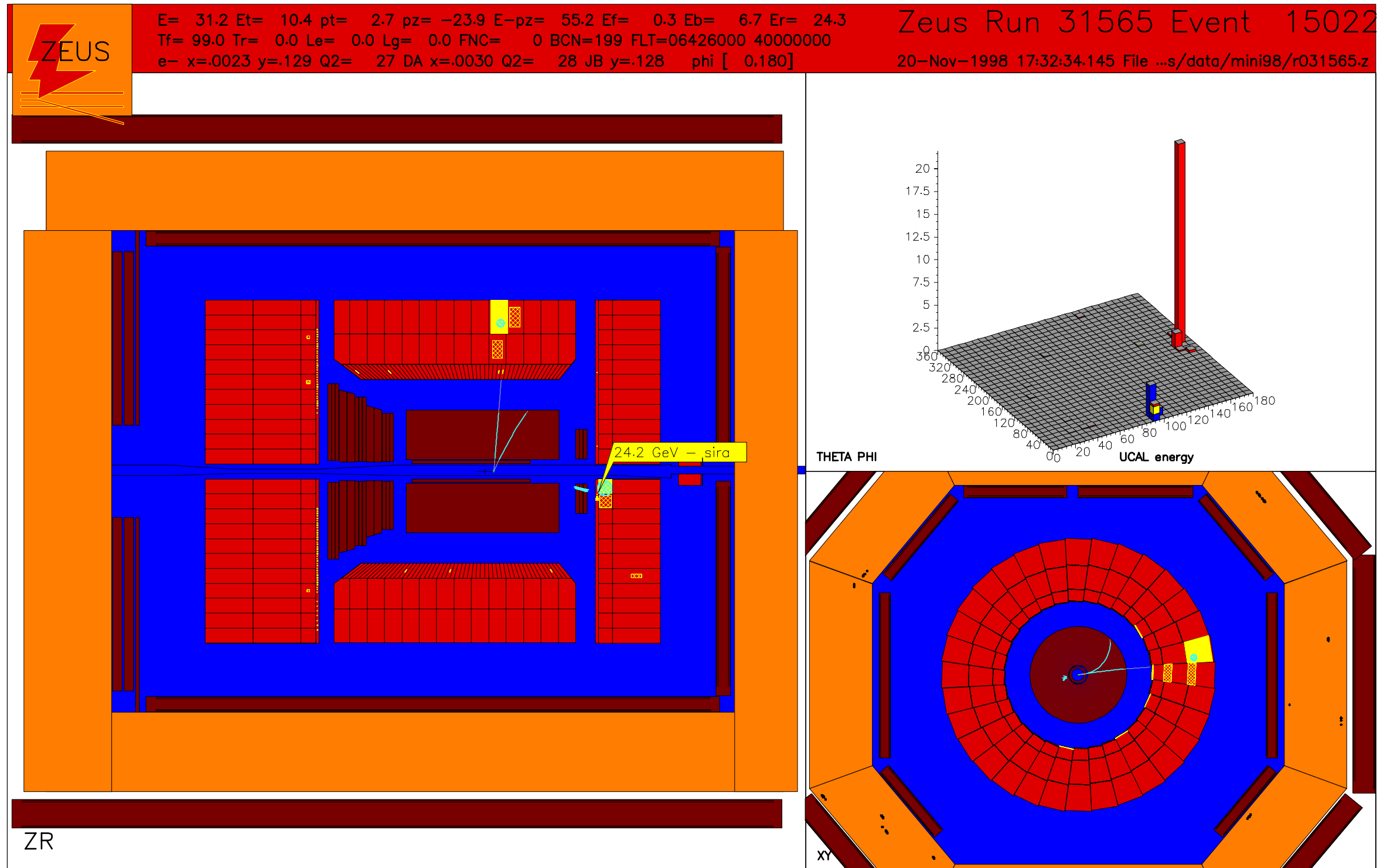


Visualising Diffractive events

A DIS event (experimental view)

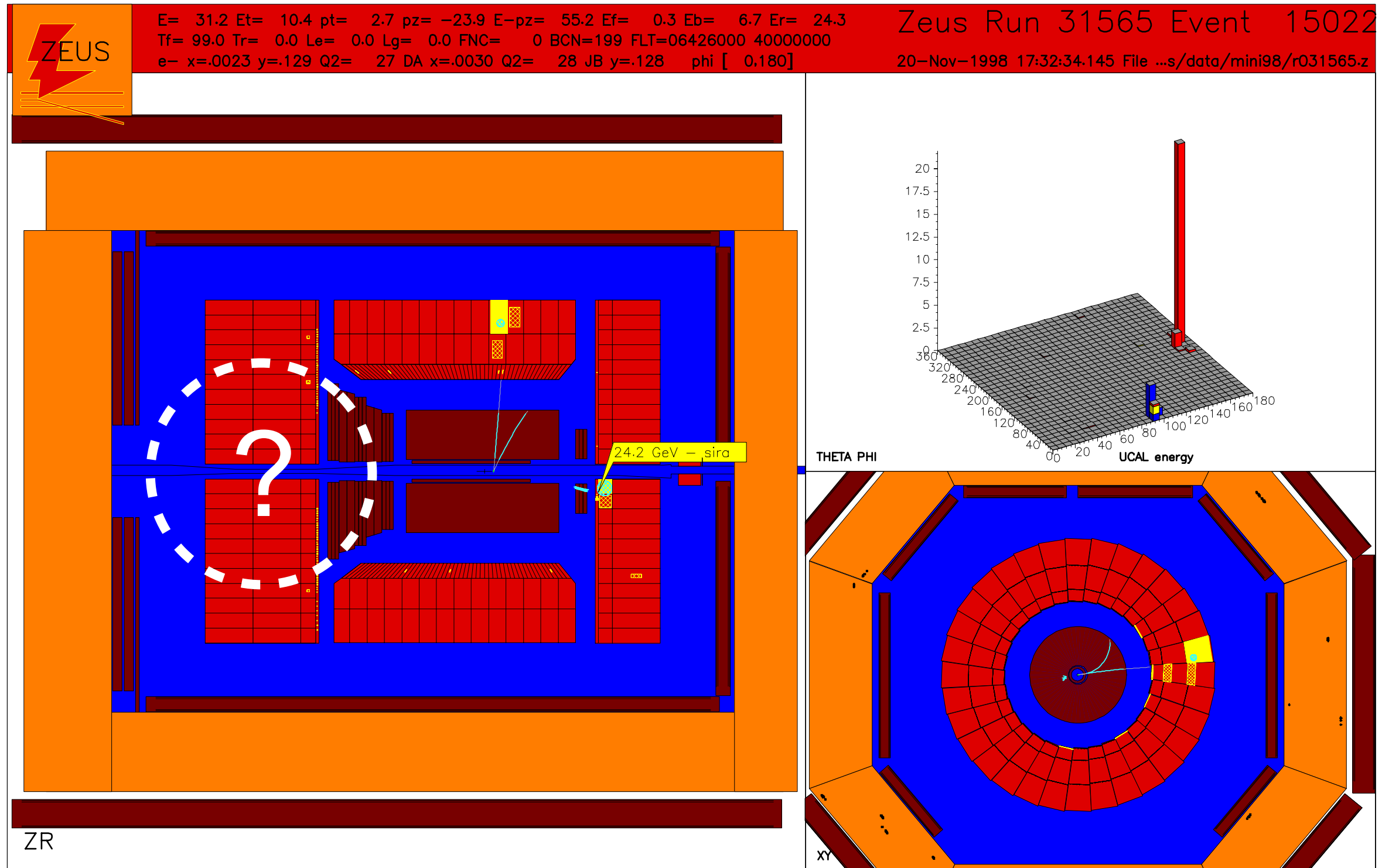


Visualising Diffractive events



Visualising Diffractive events

A diffractive event (experimental view)





Exclusive Vector Meson Production in $e+A$

- Many event generators exist for $e+p$ collisions
 - ➔ Pythia (v6), LEPTO, PEPSI, RAPGAP....
- Dearth of event generators for $e+A$ collisions
 - ➔ DPMJET-III
- Work at BNL (T. Toll, T. Ullrich) to write an $e+A$ generator (SARTRE)
 - ➔ Comparison of saturation vs non-saturation scenarios
 - ➔ First case study is that of exclusive diffractive J/ψ production



b-dependent gluons from DVCS and DVMP

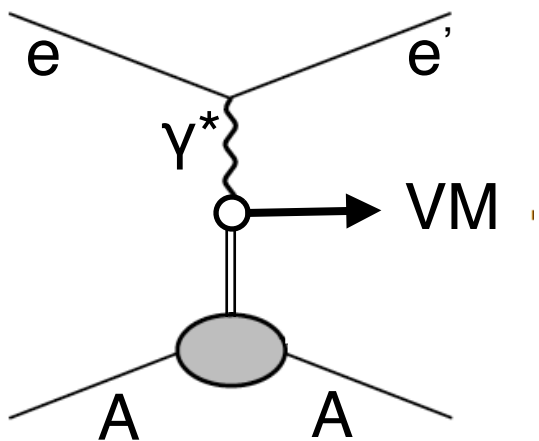
- Transverse position distribution of gluons can be determined from Deeply Virtual Compton Scattering (DVCS: $e+A \rightarrow e+\gamma+A$) and Diffractive Vector Meson Production (DVMP: $e+A \rightarrow e+VM+A$)
 - ➔ Proportional to the square of the gluon distribution!!
- Coherent diffraction (intact nuclear target)
 - ➔ transverse distribution of gluon density
- Incoherent diffraction (dissociated nuclear target)
 - ➔ transverse gluon correlations in addition



b-dependent gluons from DVCS and DVMP

- Transverse position distribution of gluons can be determined from Deeply Virtual Compton Scattering (DVCS: $e+A \rightarrow e+\gamma+A$) and Diffractive Vector Meson Production (DVMP: $e+A \rightarrow e+VM+A$)
 - ➔ Proportional to the square of the gluon distribution!!
- Coherent diffraction (intact nuclear target)
 - ➔ transverse distribution of gluon density
- Incoherent diffraction (dissociated nuclear target)
 - ➔ transverse gluon correlations in addition

DVMP

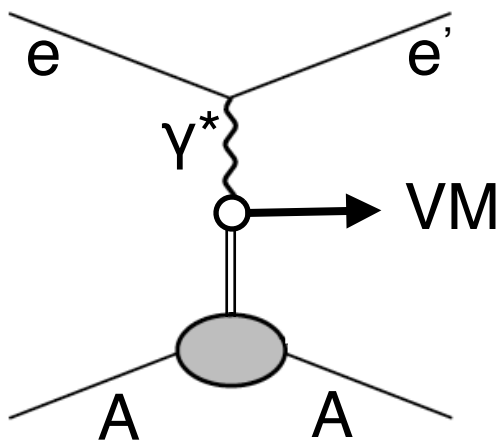




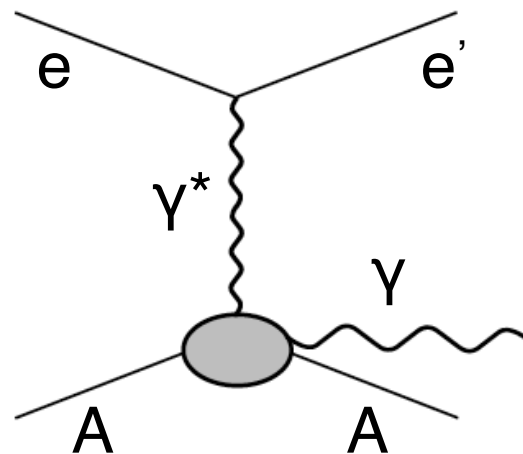
b-dependent gluons from DVCS and DVMP

- Transverse position distribution of gluons can be determined from Deeply Virtual Compton Scattering (DVCS: $e+A \rightarrow e+\gamma+A$) and Diffractive Vector Meson Production (DVMP: $e+A \rightarrow e+VM+A$)
 - ➔ Proportional to the square of the gluon distribution!!
- Coherent diffraction (intact nuclear target)
 - ➔ transverse distribution of gluon density
- Incoherent diffraction (dissociated nuclear target)
 - ➔ transverse gluon correlations in addition

DVMP



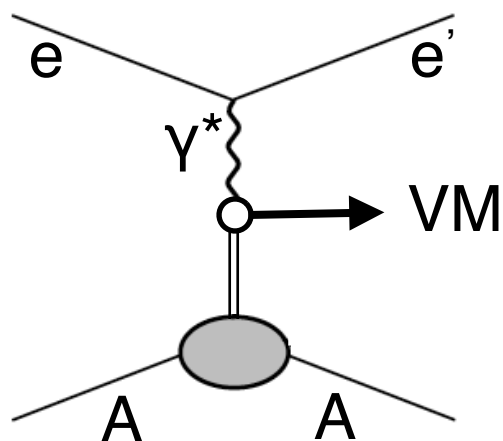
DVCS



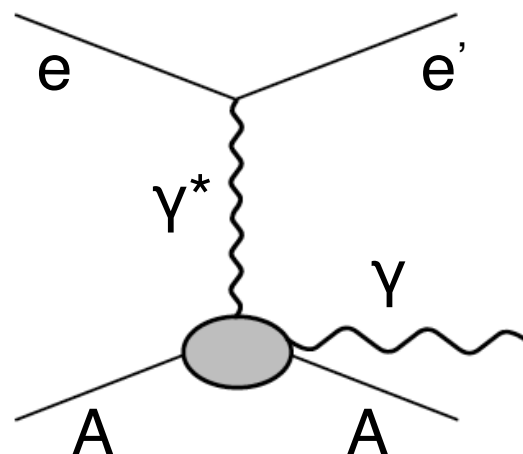
b-dependent gluons from DVCS and DVMP

- Transverse position distribution of gluons can be determined from **D**eeply **V**irtual **C**ompton **S**cattering (**DVCS**: $e+A \rightarrow e+\gamma+A$) and **D**iffractive **V**ector **M**eson **P**roduction (**DVMP**: $e+A \rightarrow e+VM+A$)
 - ➔ Proportional to the square of the gluon distribution!!
- Coherent diffraction (intact nuclear target)
 - ➔ transverse distribution of gluon density
- Incoherent diffraction (dissociated nuclear target)
 - ➔ transverse gluon correlations in addition

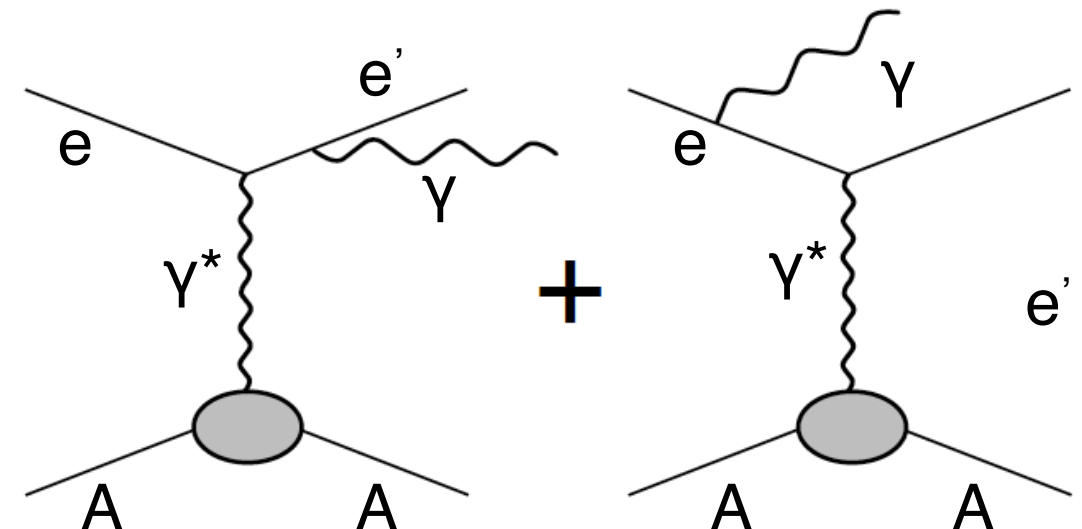
DVMP



DVCS

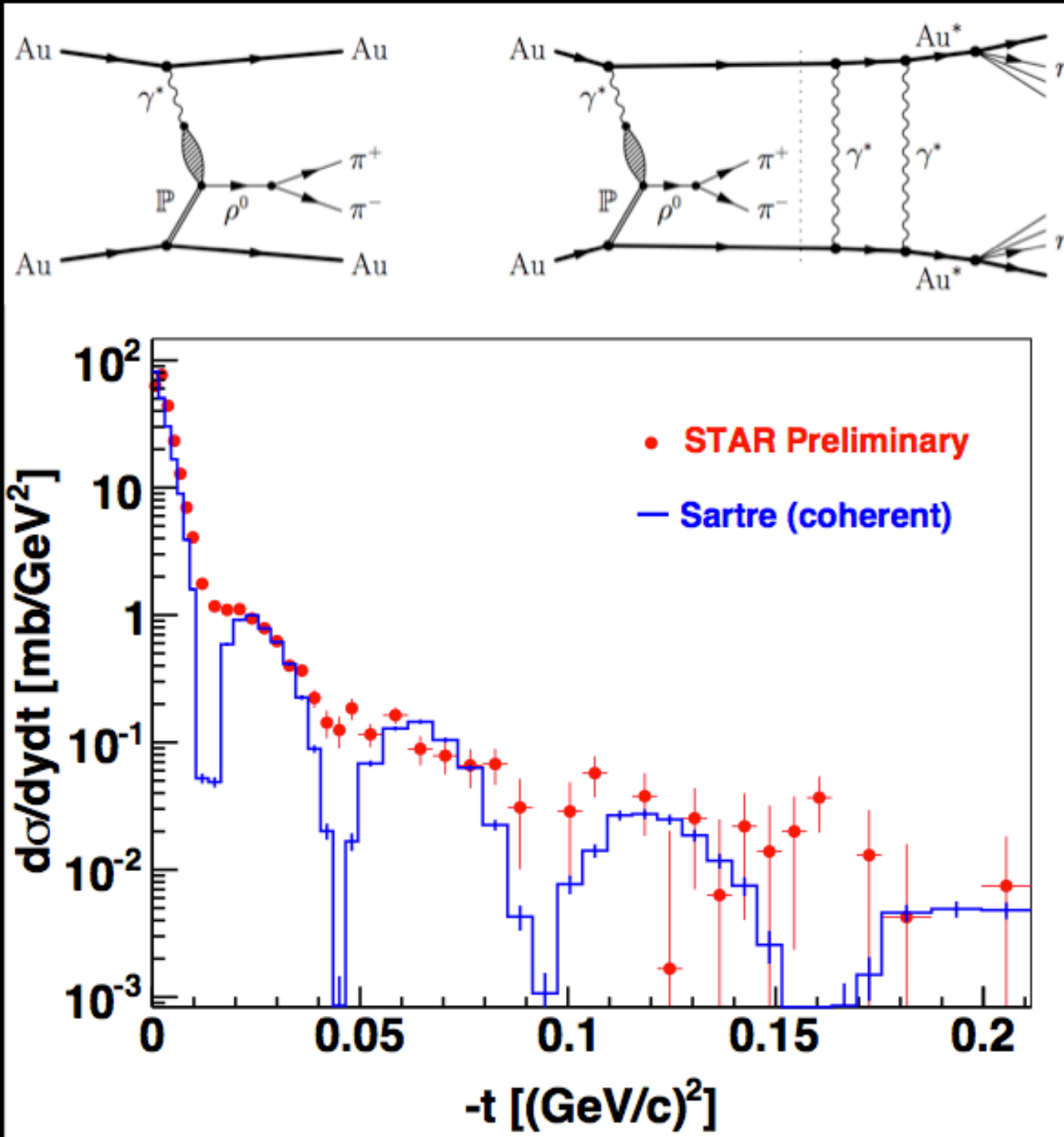


BH



DVCS and Bethe-Heitler interference terms become difficult to distinguish experimentally

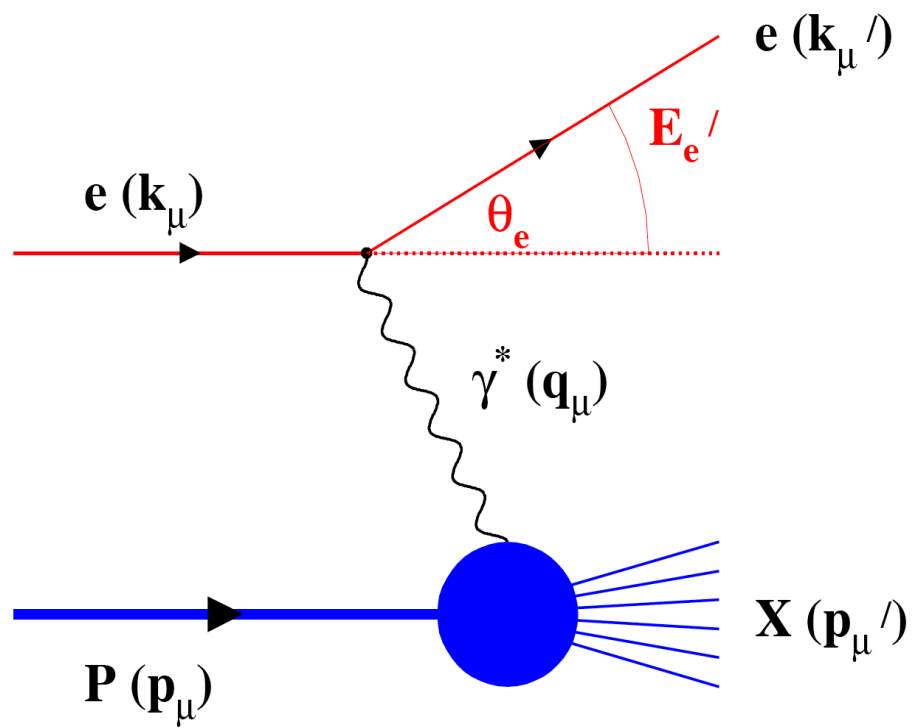
Coherent Diffraction ($\gamma^* + \text{IP}$) in UPC at RHIC



- Coherent diffractive ρ production in Au + Au at $\sqrt{s_{NN}} = 200$ GeV
- Data: STAR/RHIC Ultra-peripheral AuAu Collision
- Simulation: Sartre

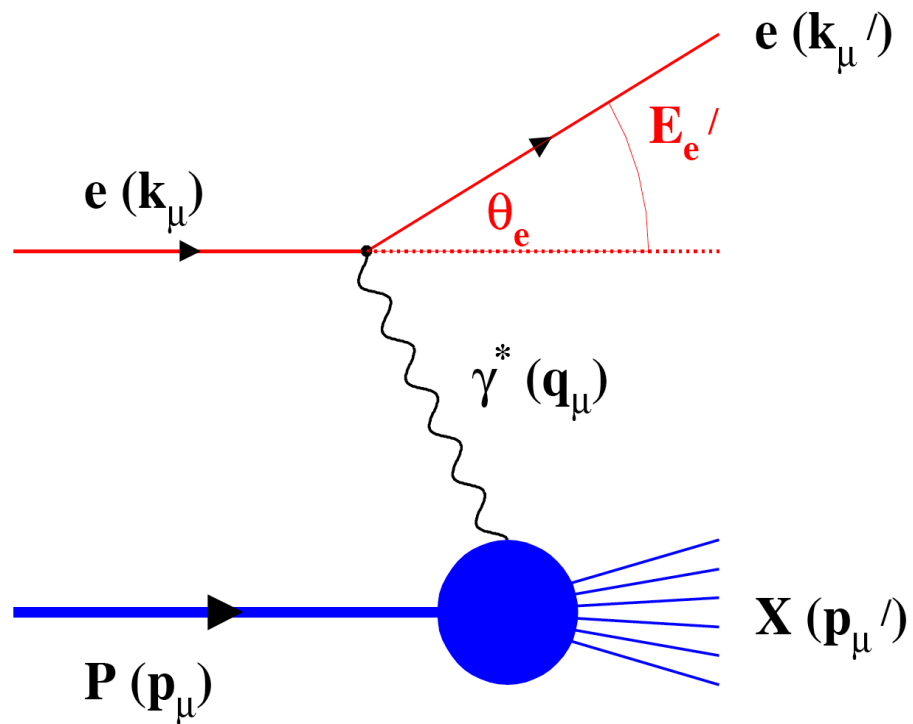
DIS Kinematics

$$e(k) + p(p) \rightarrow e(k') + X(p_X)$$



DIS Kinematics

$$e(k) + p(p) \rightarrow e(k') + X(p_X)$$



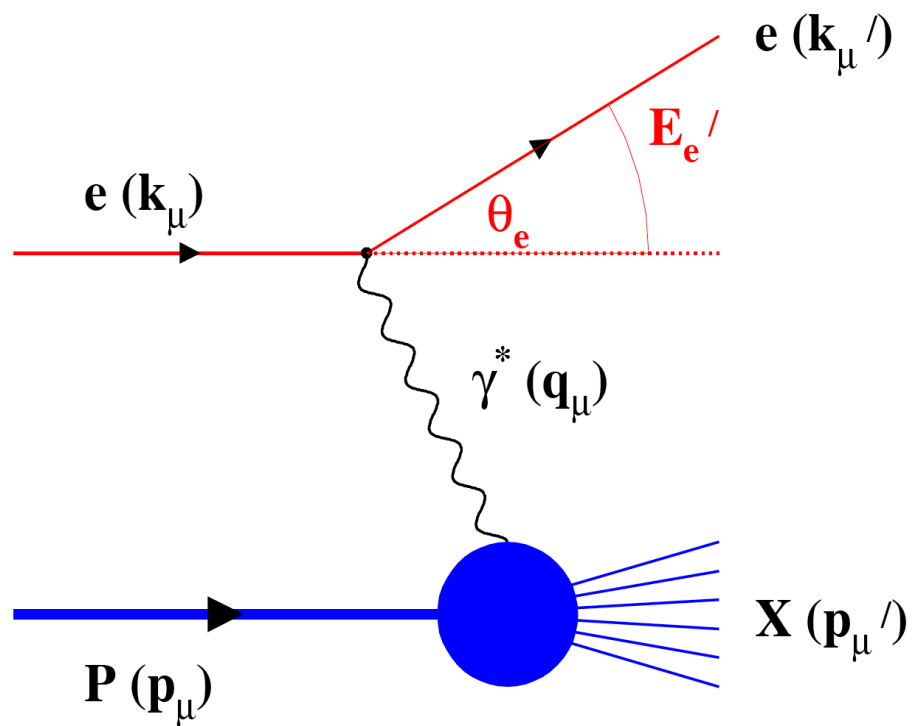
$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 4E_e E'_e \sin^2\left(\frac{\theta_e}{2}\right)$$

Measure of
resolution
power or
"Virtuality"

DIS Kinematics

$$e(k) + p(p) \rightarrow e(k') + X(p_X)$$



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 4E_e E'_e \sin^2\left(\frac{\theta'_e}{2}\right)$$

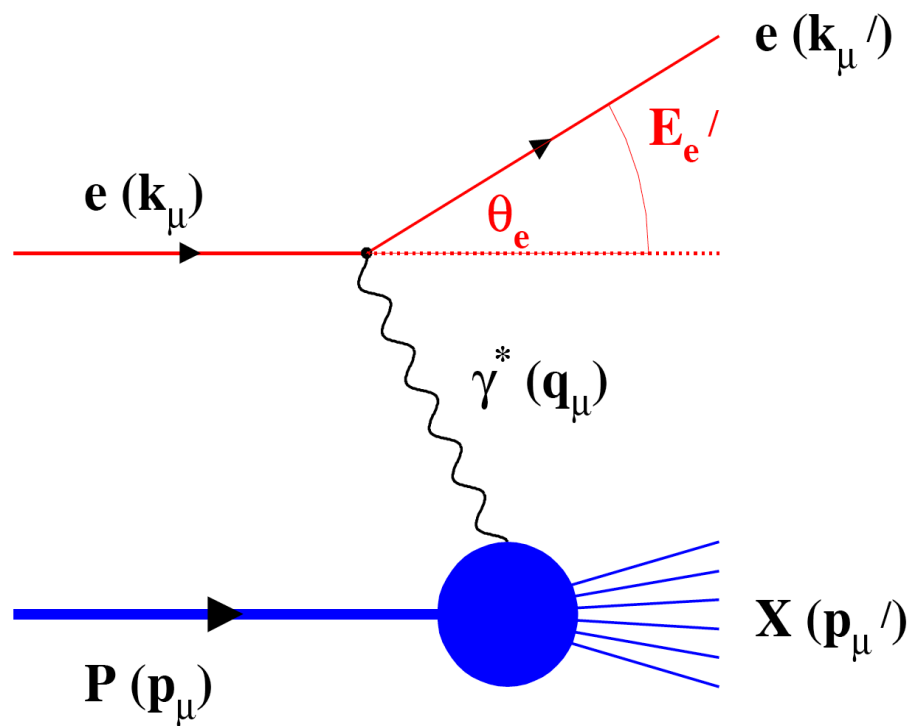
Measure of resolution power or "Virtuality"

$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of inelasticity

DIS Kinematics

$$e(k) + p(p) \rightarrow e(k') + X(p_X)$$



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 4E_e E_e' \sin^2\left(\frac{\theta_e}{2}\right)$$

Measure of resolution power or "Virtuality"

$$y = \frac{pq}{pk} = 1 - \frac{E_e'}{E_e} \cos^2\left(\frac{\theta_e}{2}\right)$$

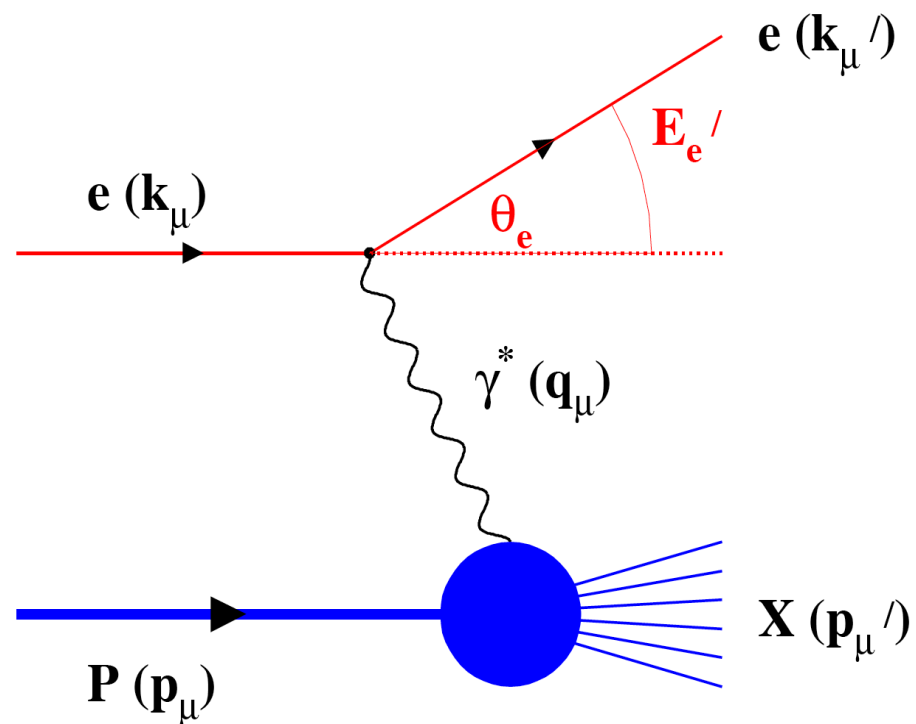
Measure of inelasticity

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

DIS Kinematics

$$e(k) + p(p) \rightarrow e(k') + X(p_X)$$



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

Measure of
resolution
power or
"Virtuality"

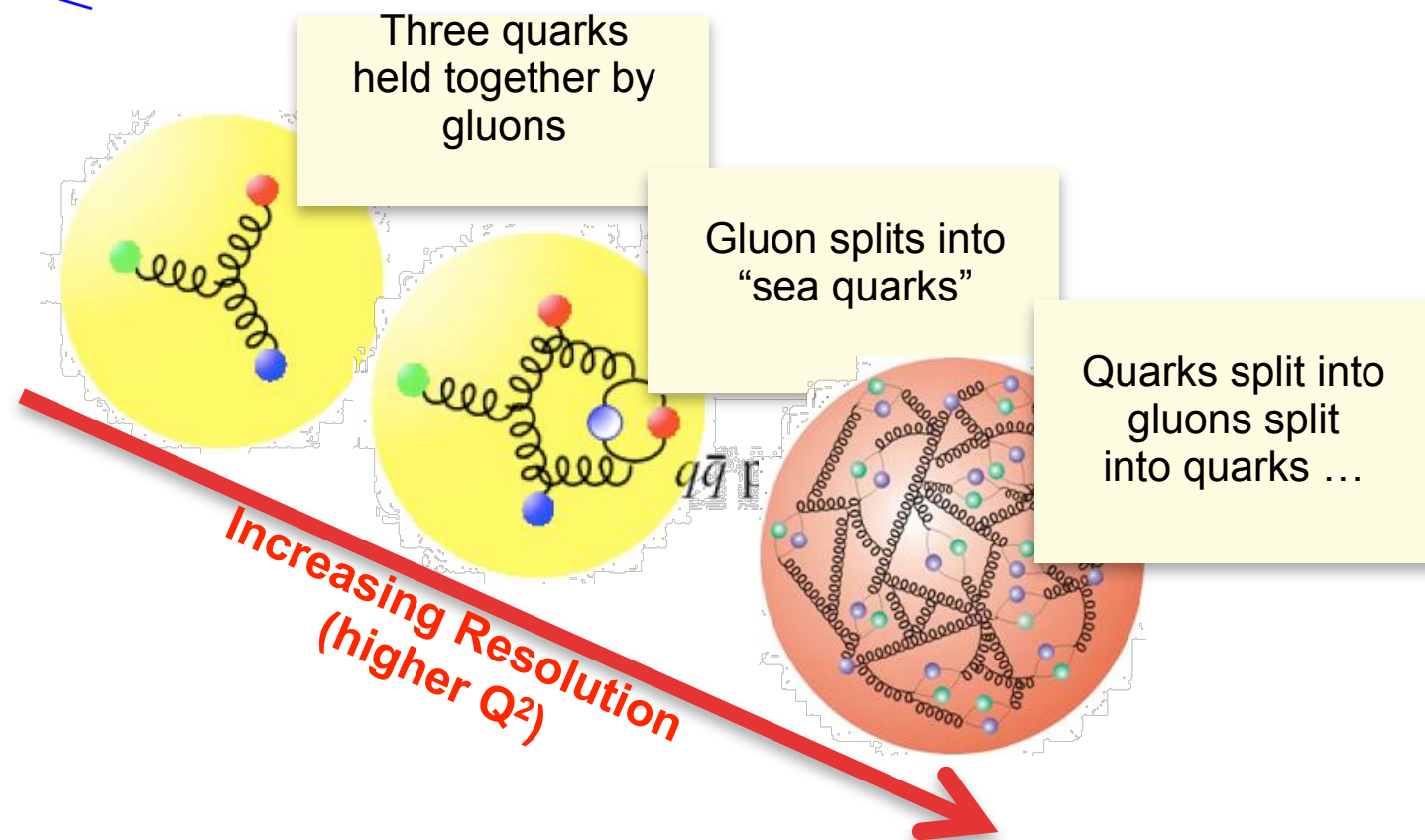
$$Q^2 = 4E_e E'_e \sin^2\left(\frac{\theta'_e}{2}\right)$$

$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of
inelasticity

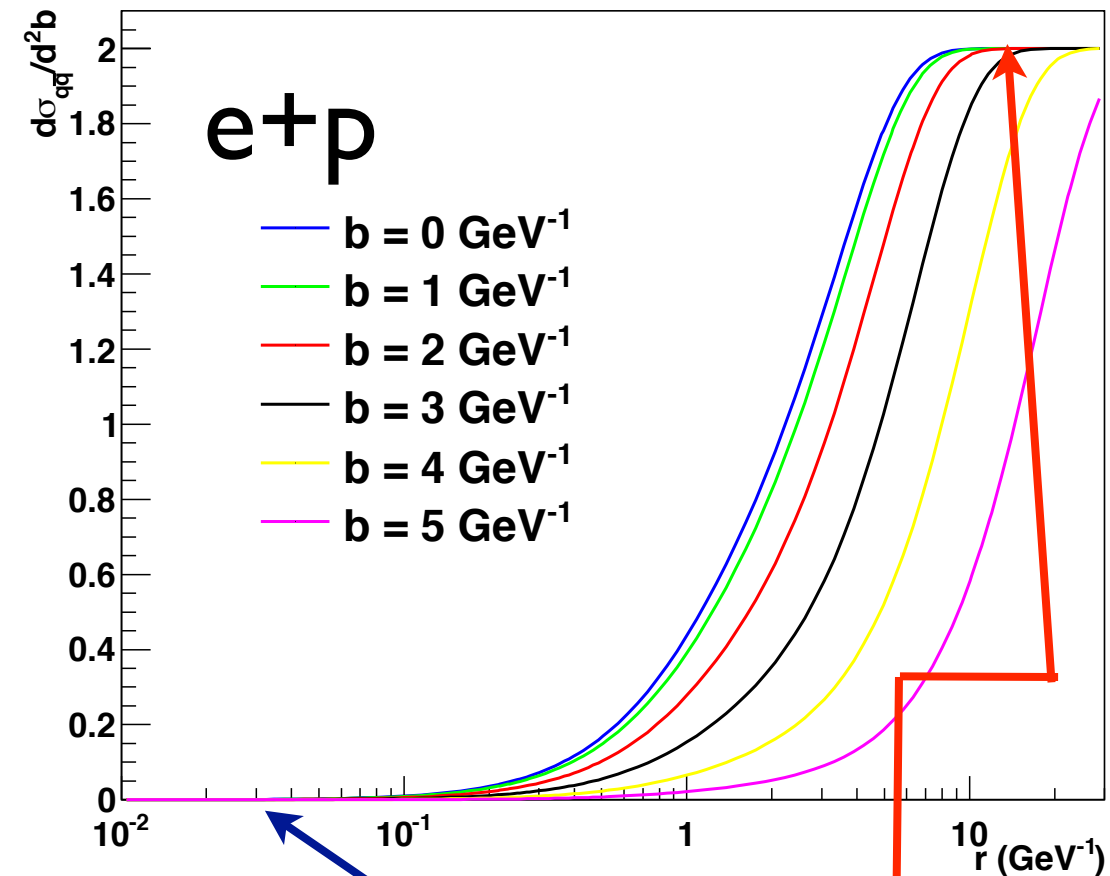
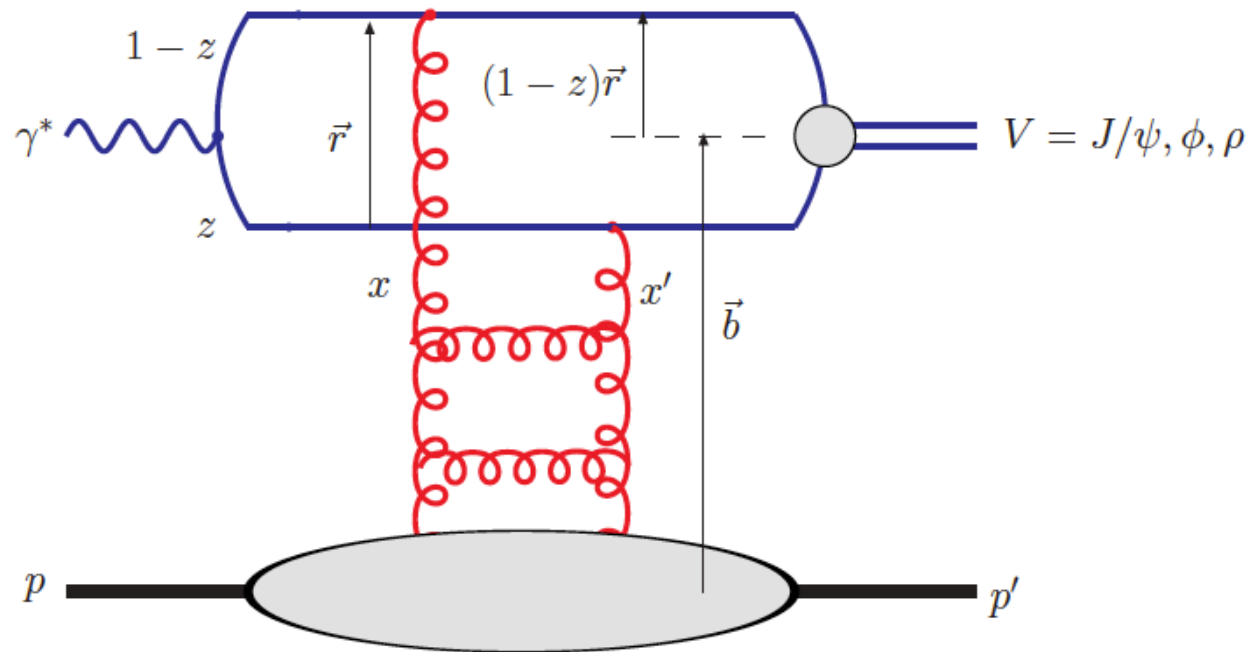
$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of
momentum
fraction of
struck
quark



Getting a “Feel” for Non-Linear QCD

Dipole Model: $\frac{d\sigma_{q\bar{q}}}{d^2b} = 2\mathcal{N}(x, r, b)$



$$\mathcal{N}(x, r, b) = 1 - \exp \left(-r^2 \frac{\pi^2}{2N_c} \alpha_s(\mu^2) x G(x, \mu^2) T(b) \right)$$

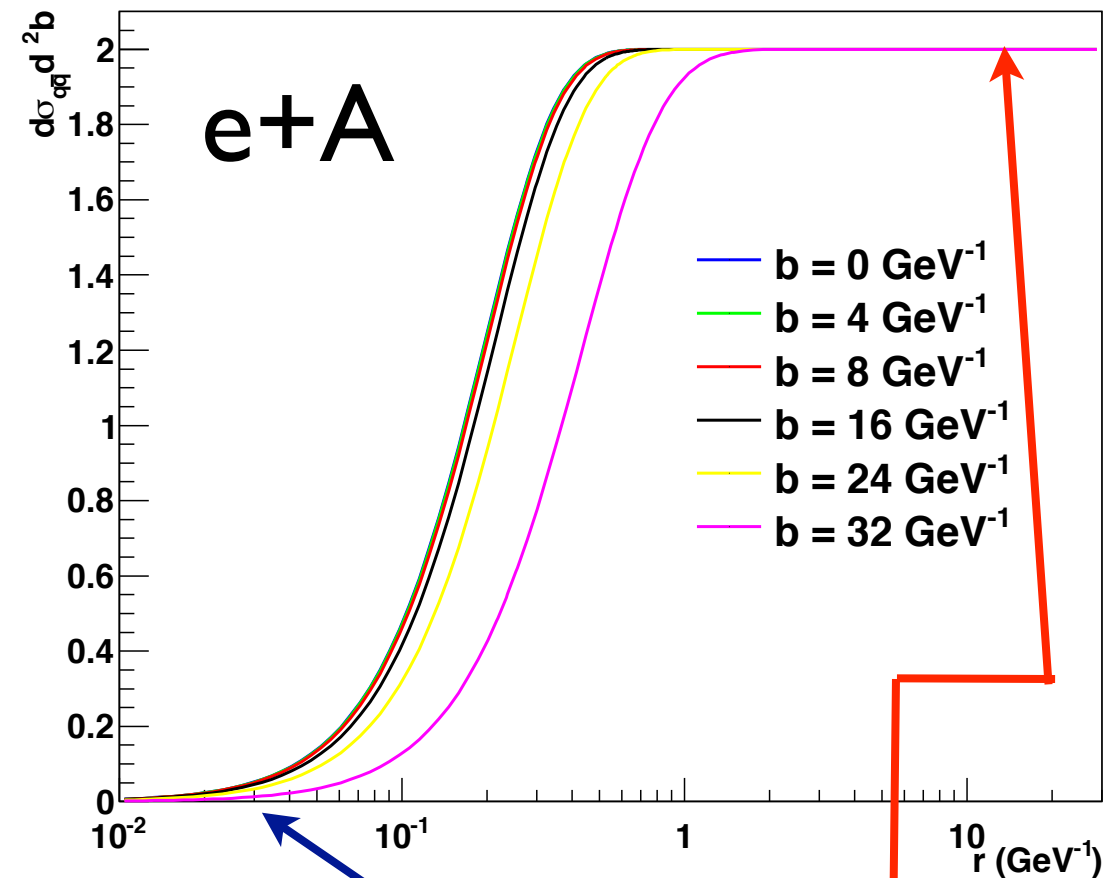
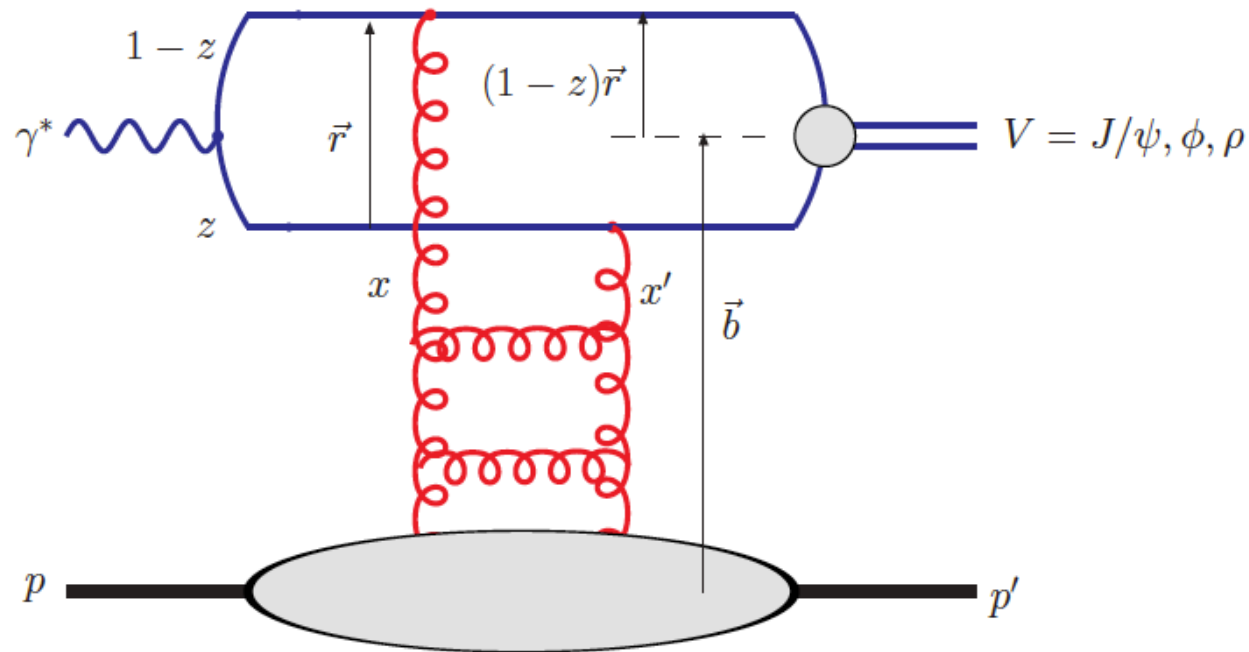
\mathcal{N} = Dipole Scattering Amplitude

0 dilute system, linear QCD

1 saturated, non-linear regime

Getting a “Feel” for Non-Linear QCD

Dipole Model: $\frac{d\sigma_{q\bar{q}}}{d^2b} = 2\mathcal{N}(x, r, b)$



$$\mathcal{N}(x, r, b) = 1 - \exp \left(-r^2 \frac{\pi^2}{2N_c} \alpha_s(\mu^2) x G(x, \mu^2) T(b) \right)$$

\mathcal{N} = Dipole Scattering Amplitude

0 dilute system, linear QCD

1 saturated, non-linear regime



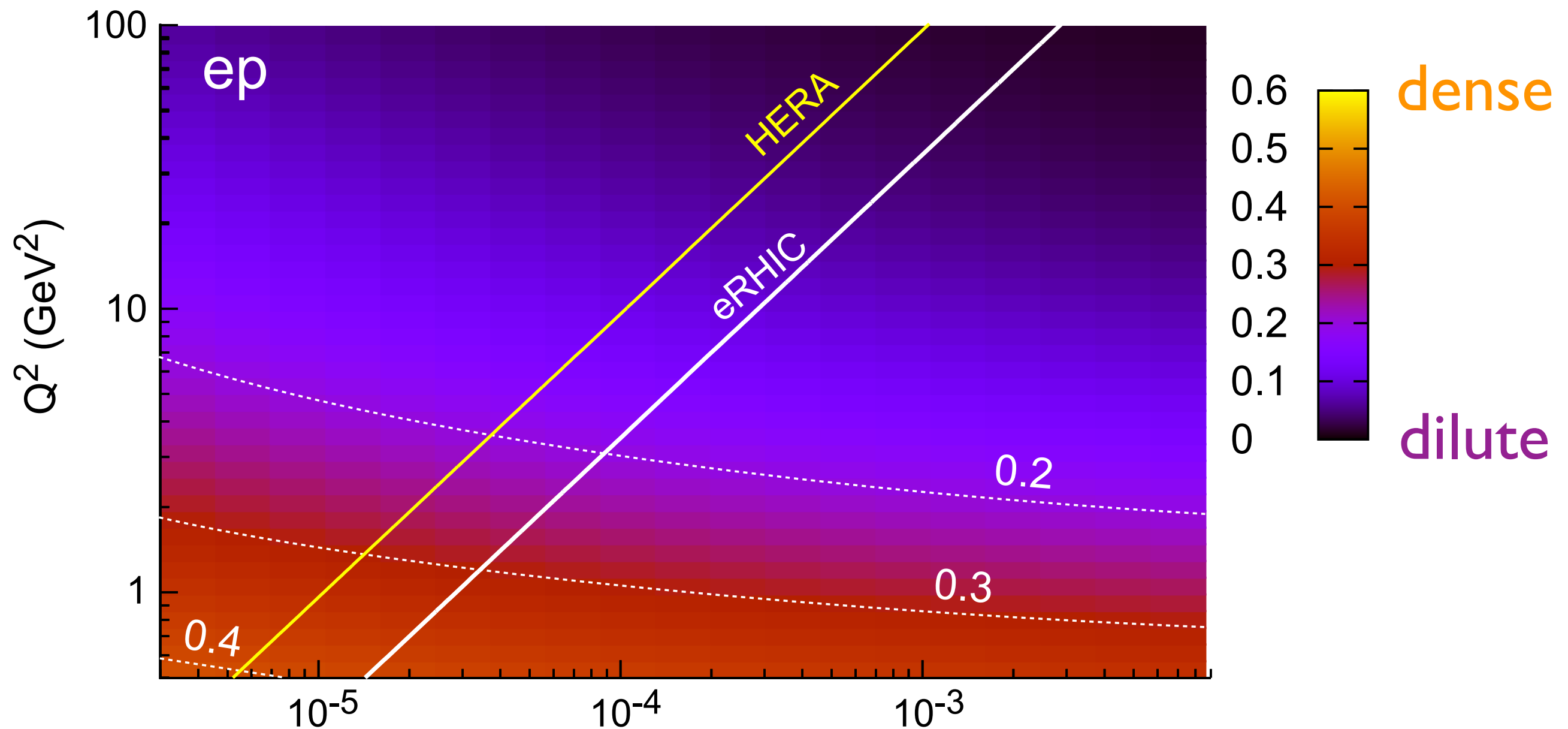
Getting a “Feel” for Non-Linear QCD

To assess typical values of \mathcal{N} calculate average:

$$\langle \mathcal{N} \rangle_{2,L} = \frac{\int d^2b d^2r dz [\psi^* \psi]_{2,L} \mathcal{N}^2}{\int d^2b d^2r dz [\psi^* \psi]_{2,L} \mathcal{N}}$$

$$\langle \mathcal{N} \rangle_2 \rightarrow F_2$$

$$\langle \mathcal{N} \rangle_L \rightarrow F_L$$





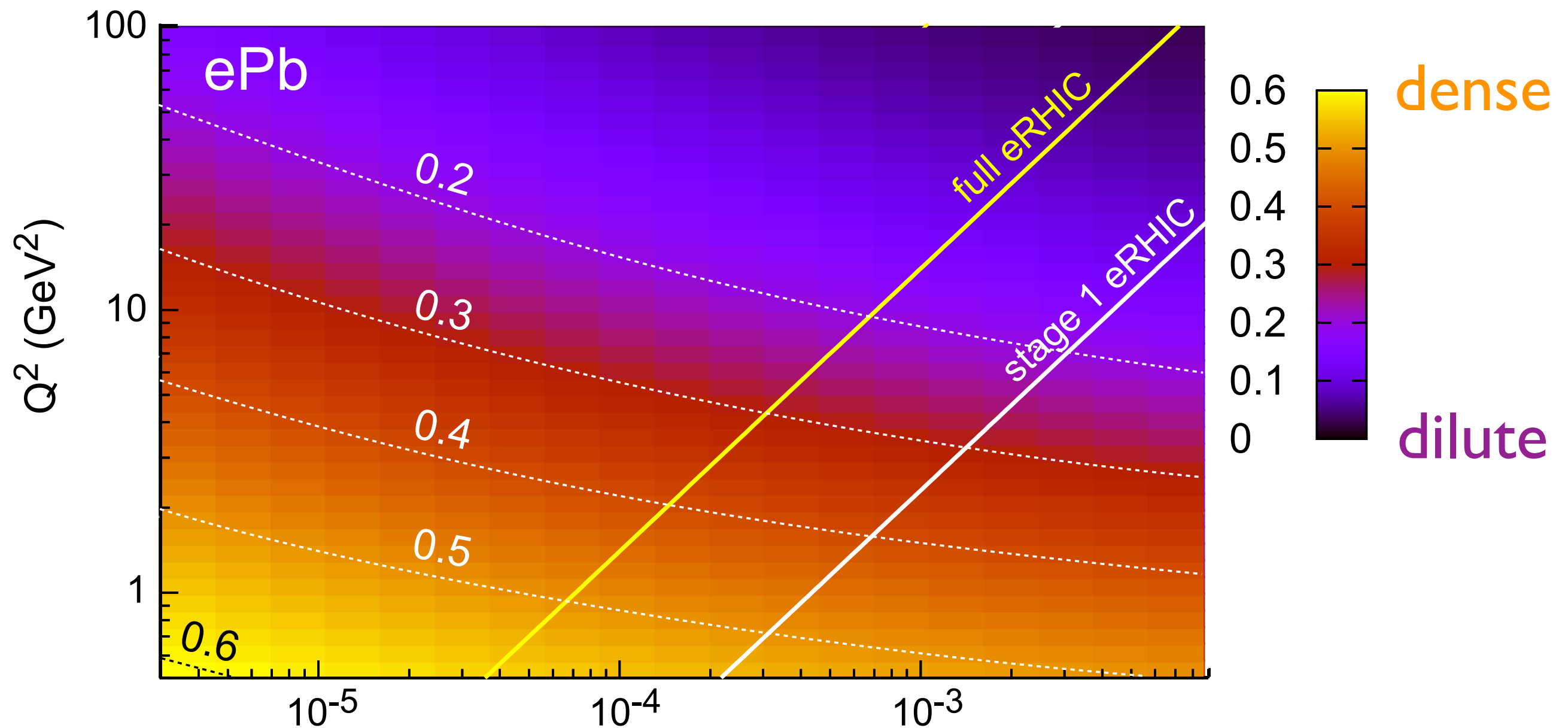
Getting a “Feel” for Non-Linear QCD

To assess typical values of \mathcal{N} calculate average:

$$\langle \mathcal{N} \rangle_{2,L} = \frac{\int d^2b d^2r dz [\psi^* \psi]_{2,L} \mathcal{N}^2}{\int d^2b d^2r dz [\psi^* \psi]_{2,L} \mathcal{N}}$$

$$\langle \mathcal{N} \rangle_2 \rightarrow F_2$$

$$\langle \mathcal{N} \rangle_L \rightarrow F_L$$

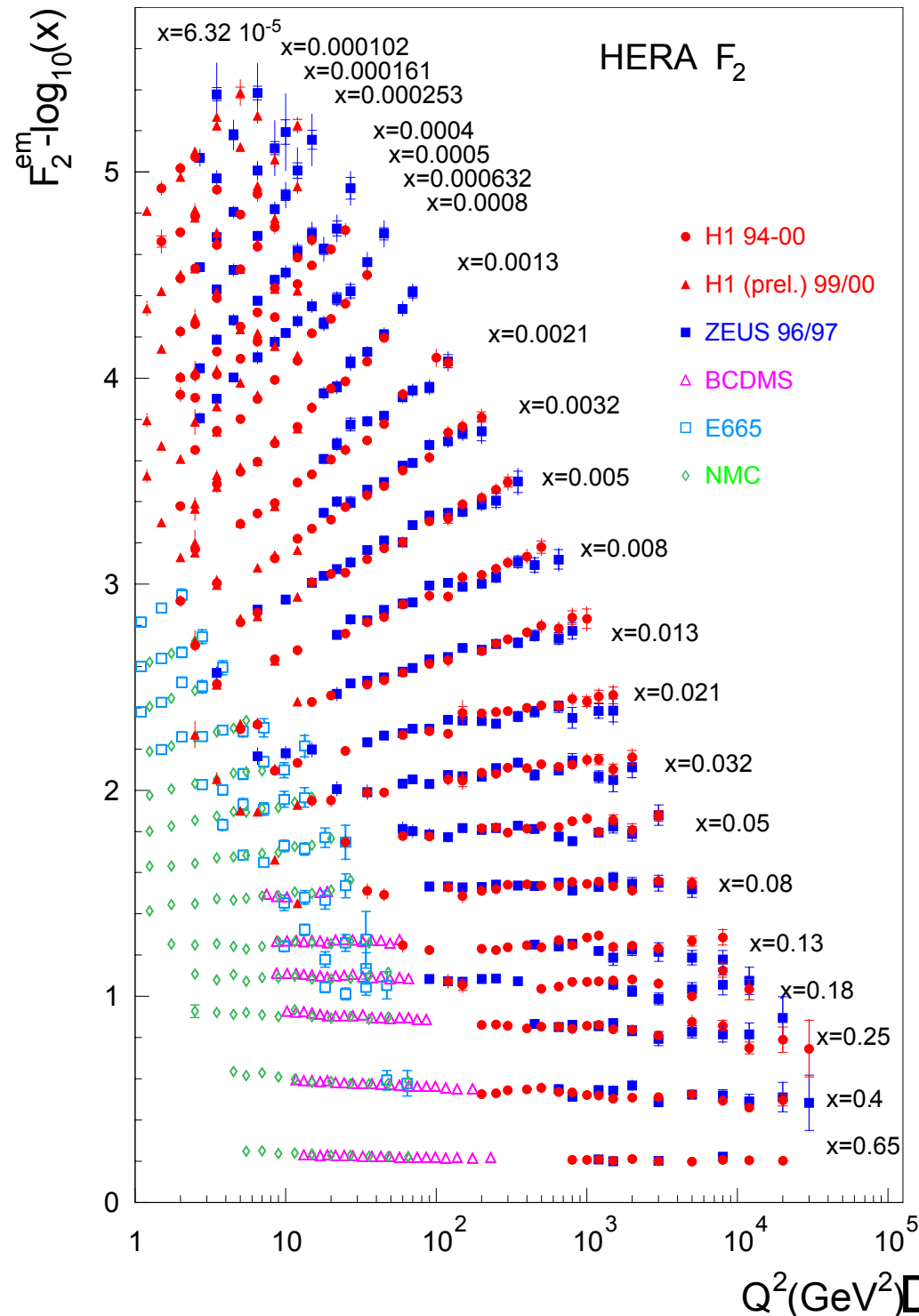


Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

**quark+anti-quark
momentum distributions**

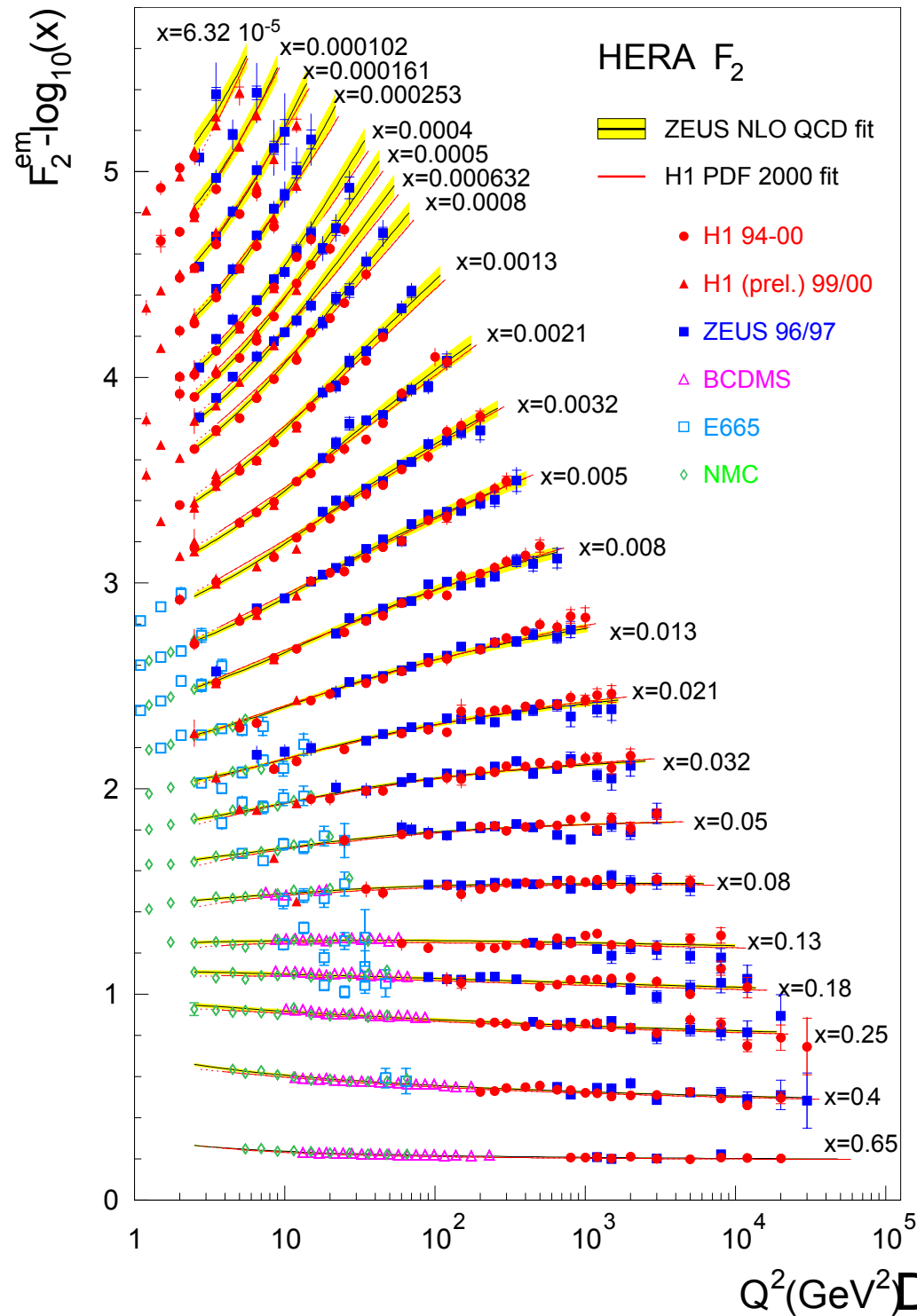
**gluon momentum
distribution**



Measuring the glue via Structure Functions

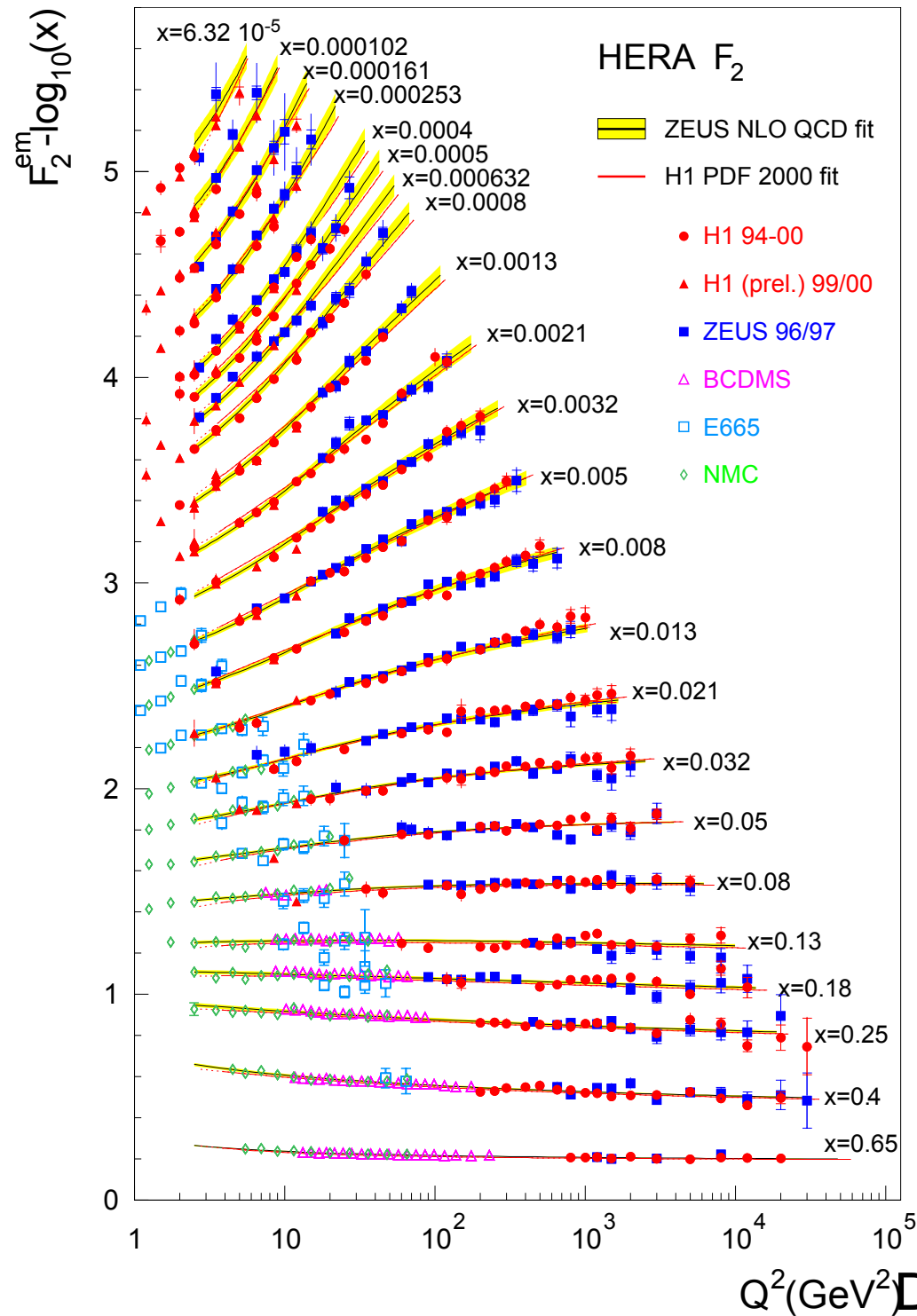
$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP
Evolution $\Rightarrow G(x, Q^2)$

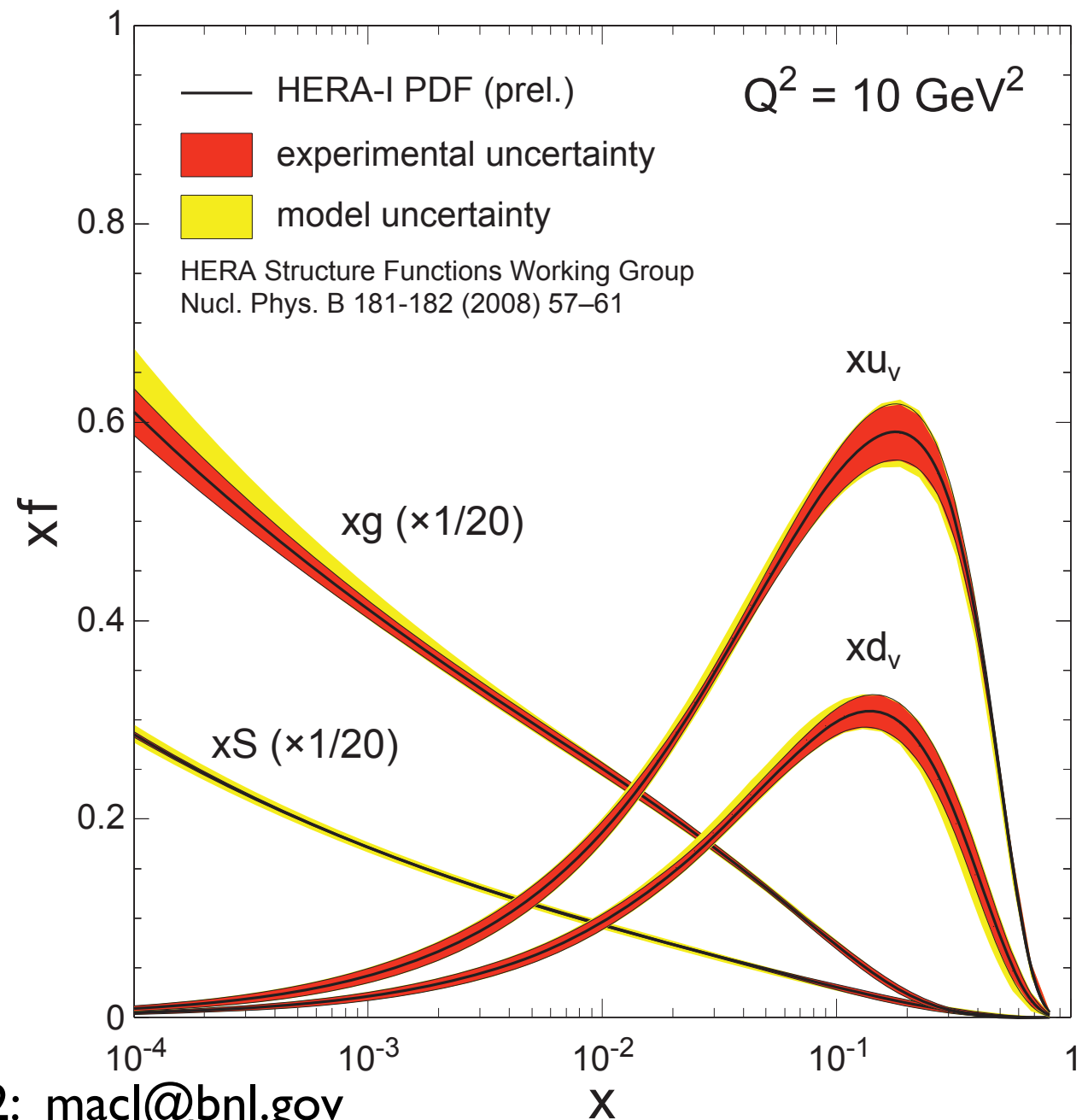


Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

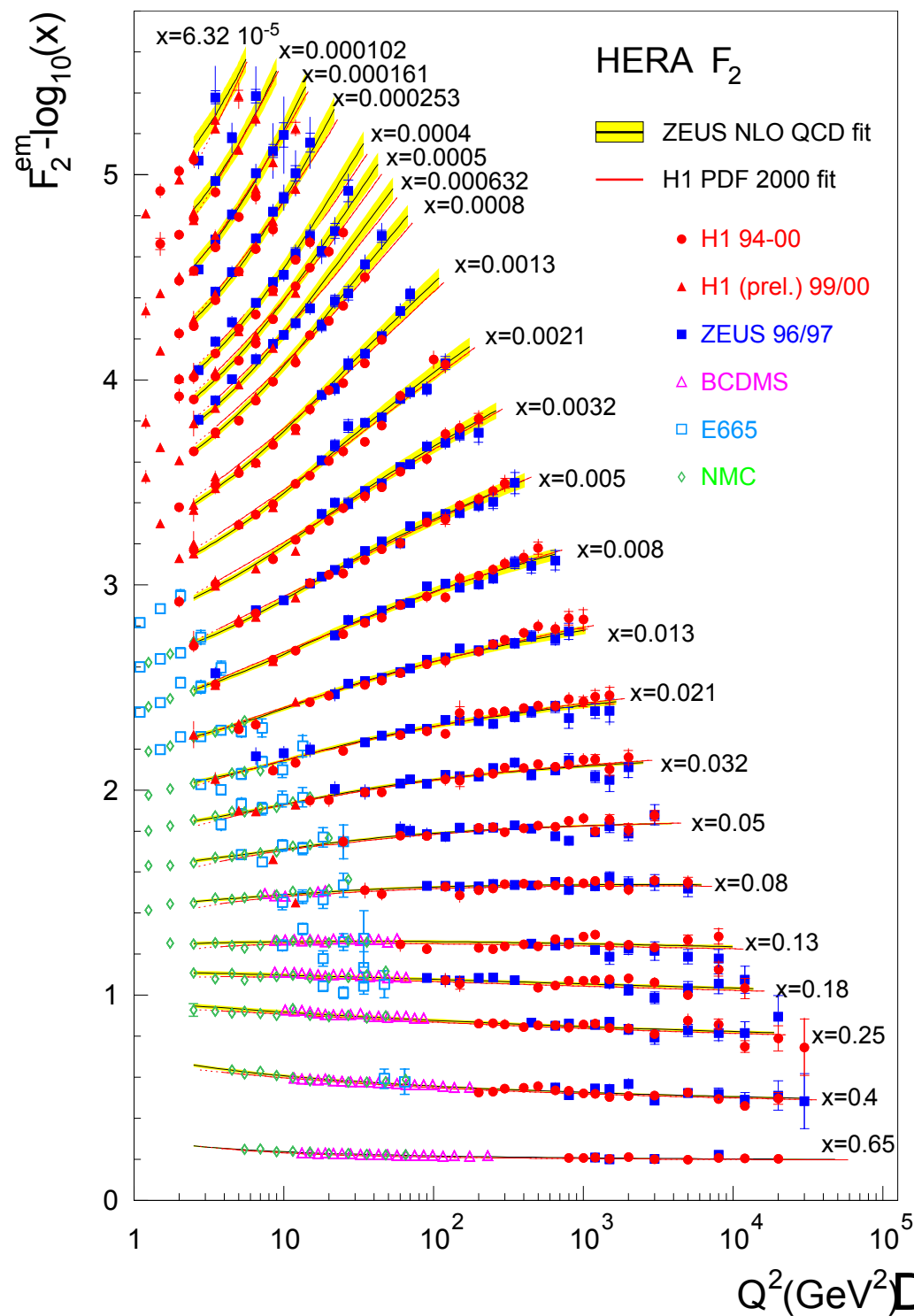


Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP
Evolution $\Rightarrow G(x, Q^2)$

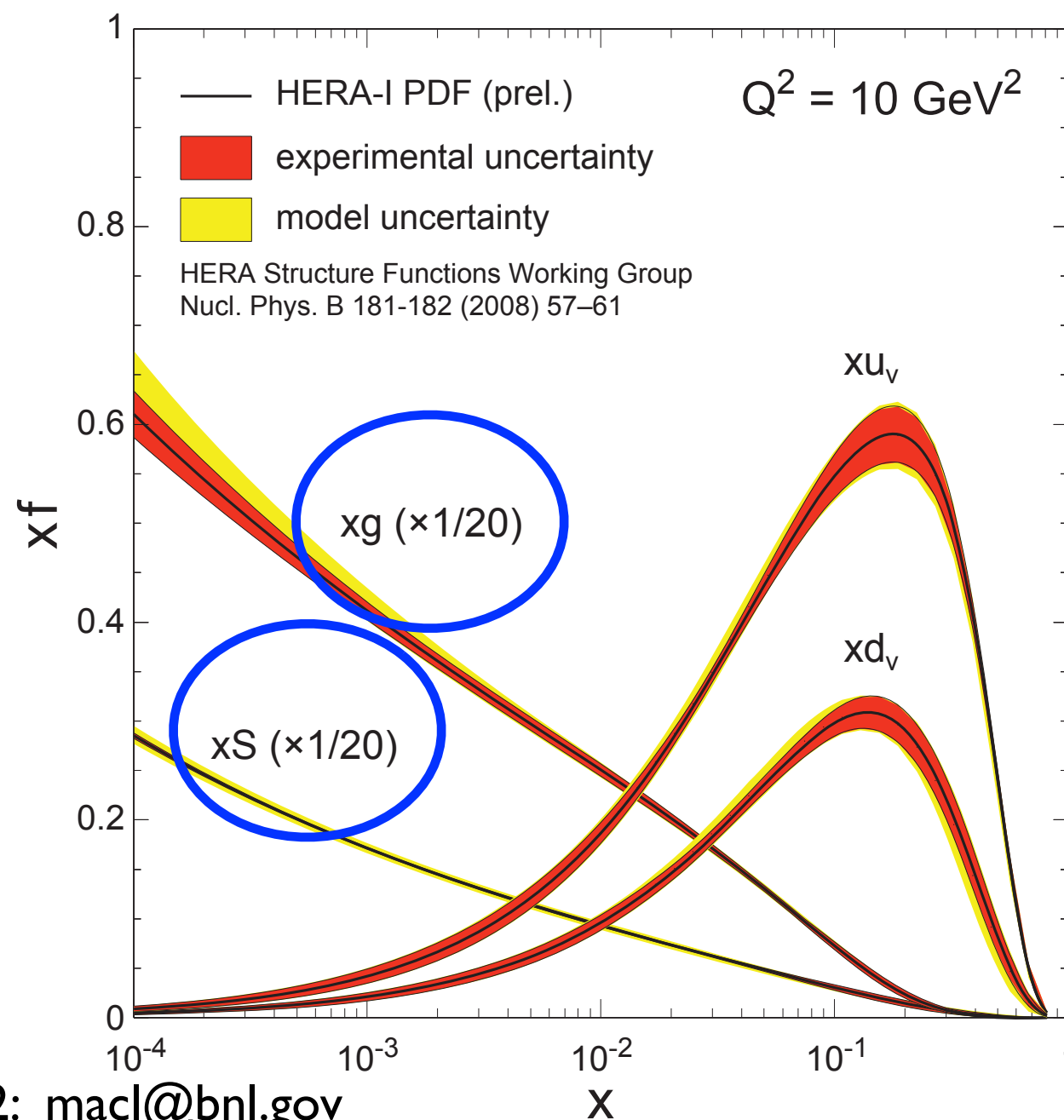


Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

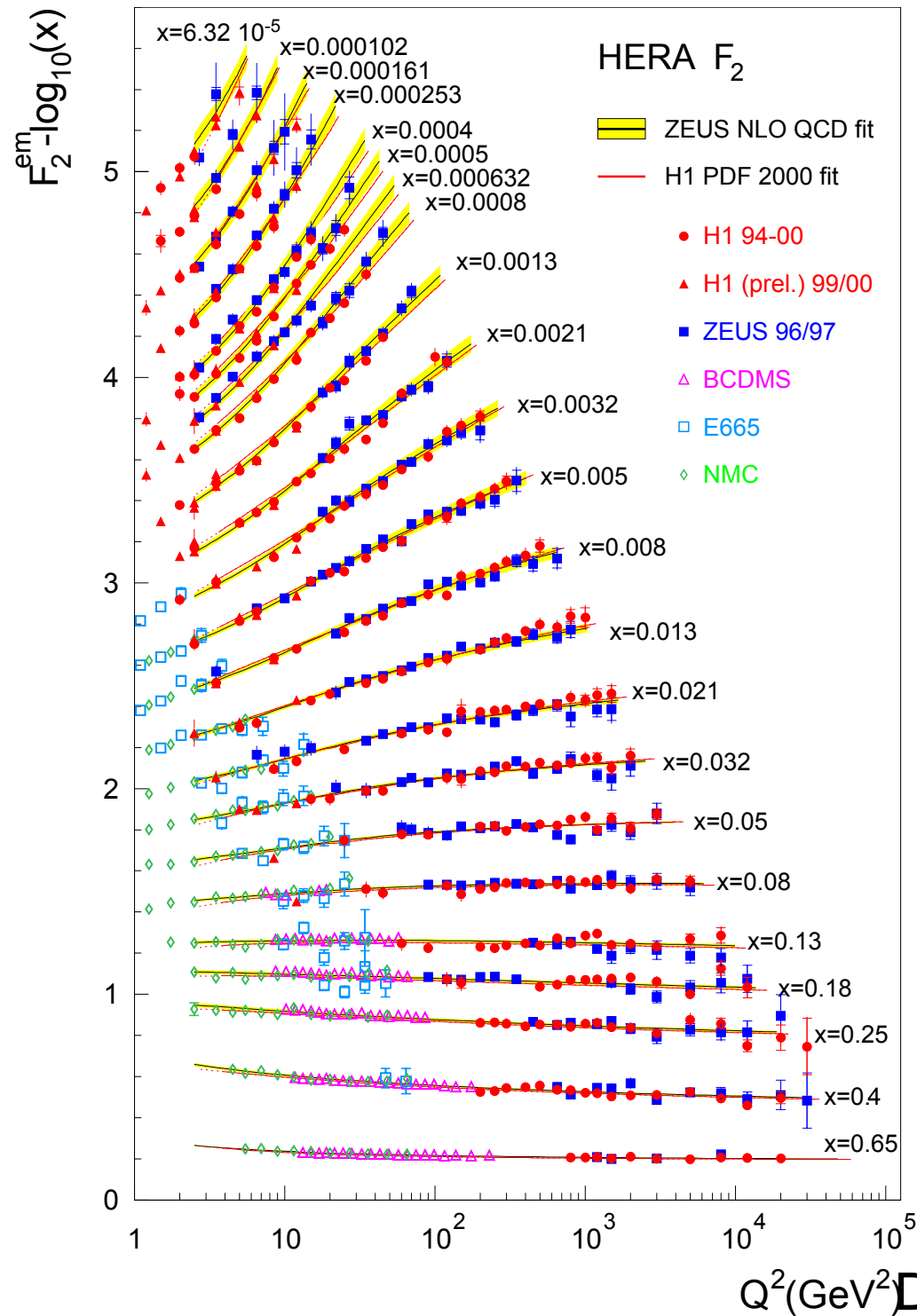


Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP
Evolution $\Rightarrow G(x, Q^2)$

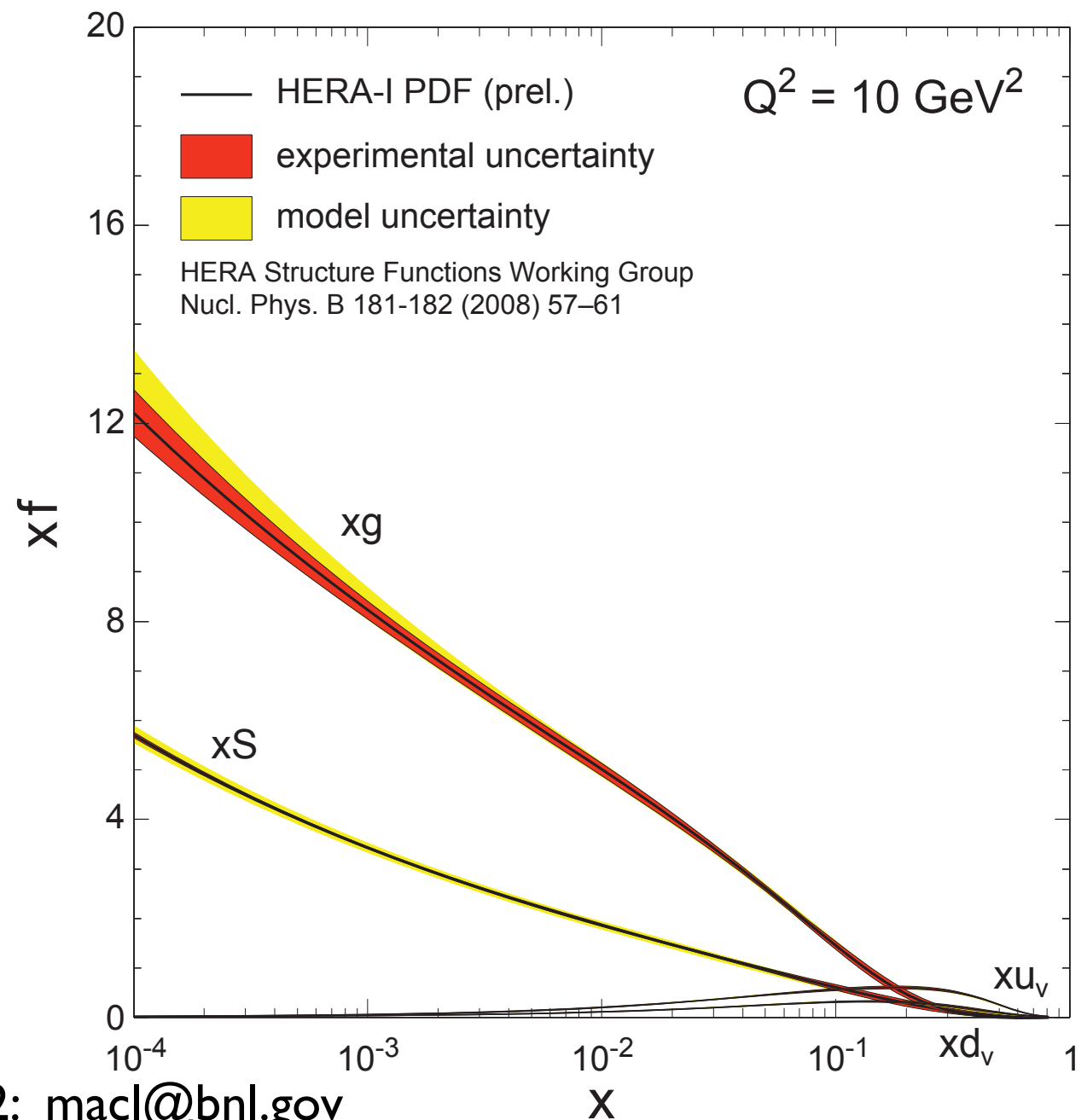


Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$



Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP
Evolution $\Rightarrow G(x, Q^2)$

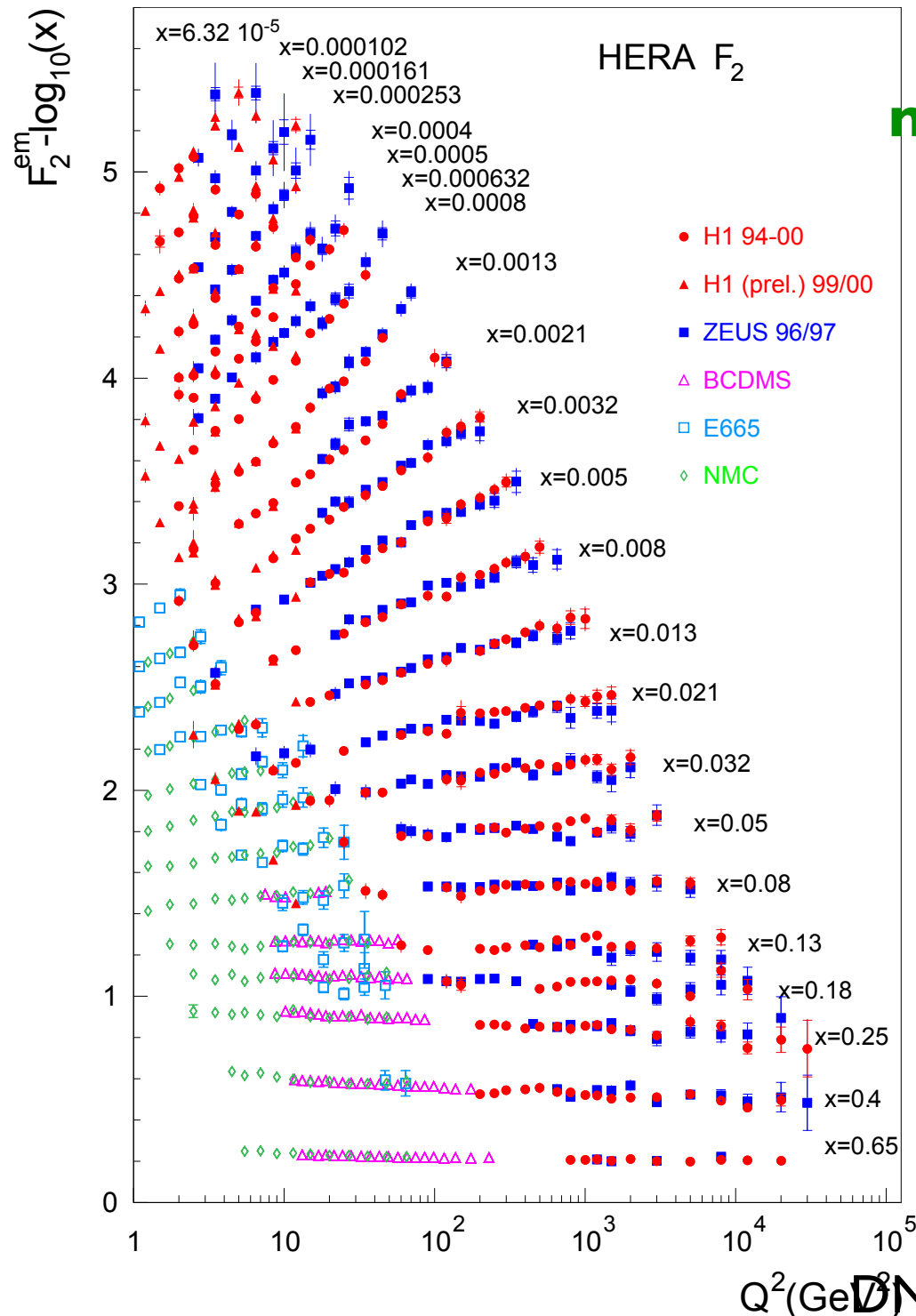


Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$

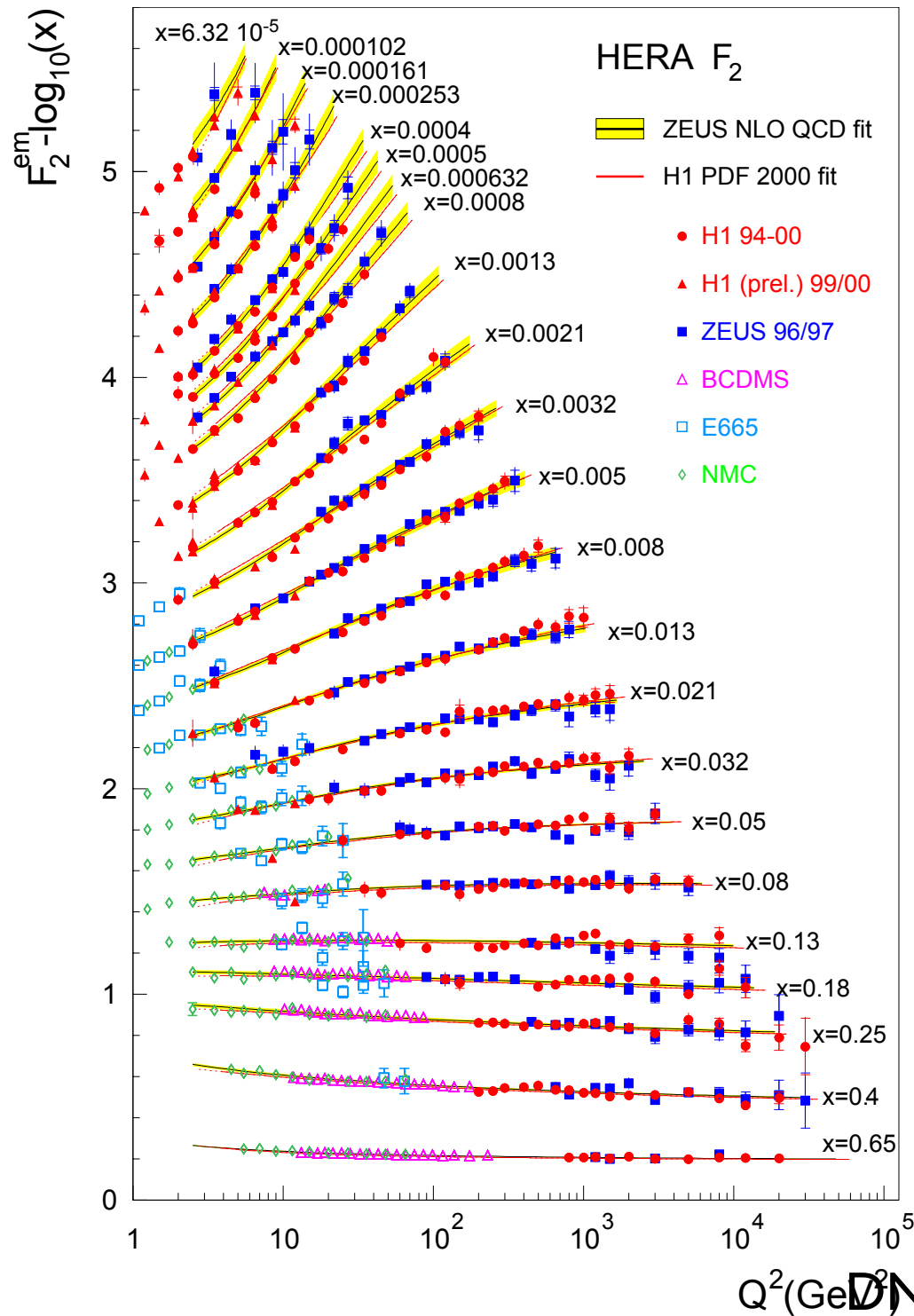
**quark+anti-quark
momentum distributions**

**gluon momentum
distribution**



Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y_+} F_L^A(x, Q^2)$$



Scaling violation: $dF_2/d\ln Q^2$ and linear DGLAP
Evolution $\Rightarrow G(x, Q^2)$

